



Auroral processes in satellite data



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Overview

- Setting the scene
- Auroral processes in the **upward** current region
- Auroral processes in the **downward** current region
- Dayside and cusp aurora, theta aurora
- Temporal evolution in the auroral zone

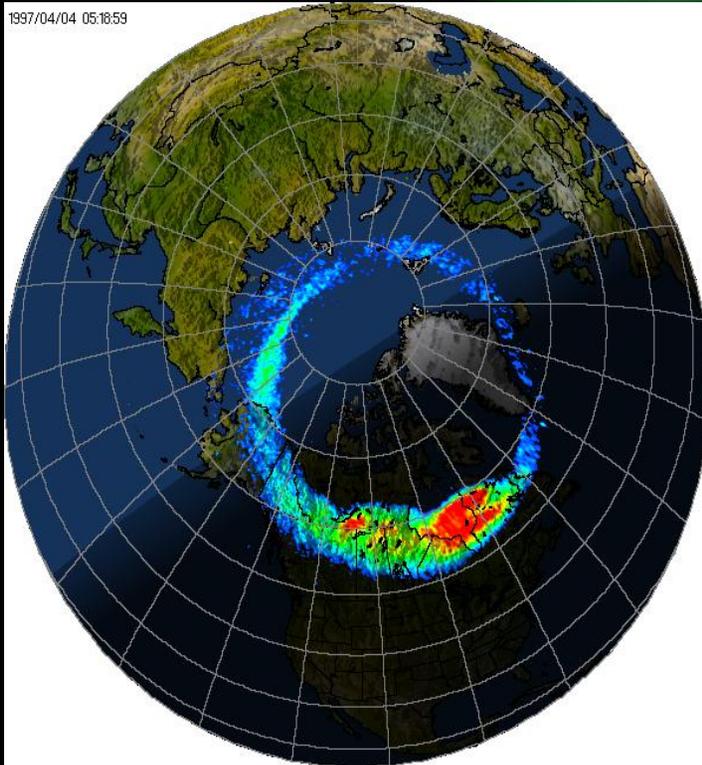
Very useful reference:

Space Science Series of ISSI
Auroral Plasma Physics

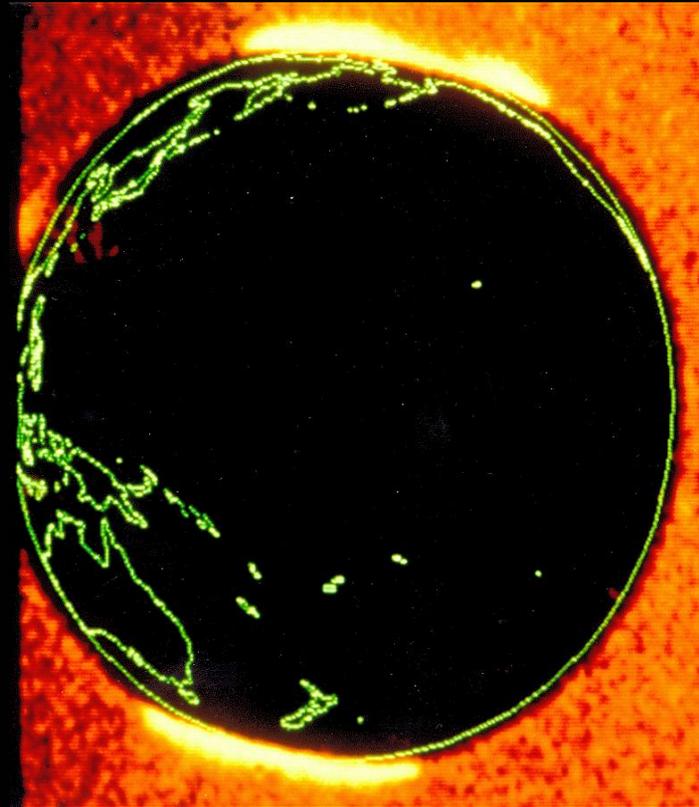
Götz Paschmann, Stein Haaland, Rudolf Treumann

(Space Science Reviews, vol 103, 1-4, 2002)

Auroral ovals

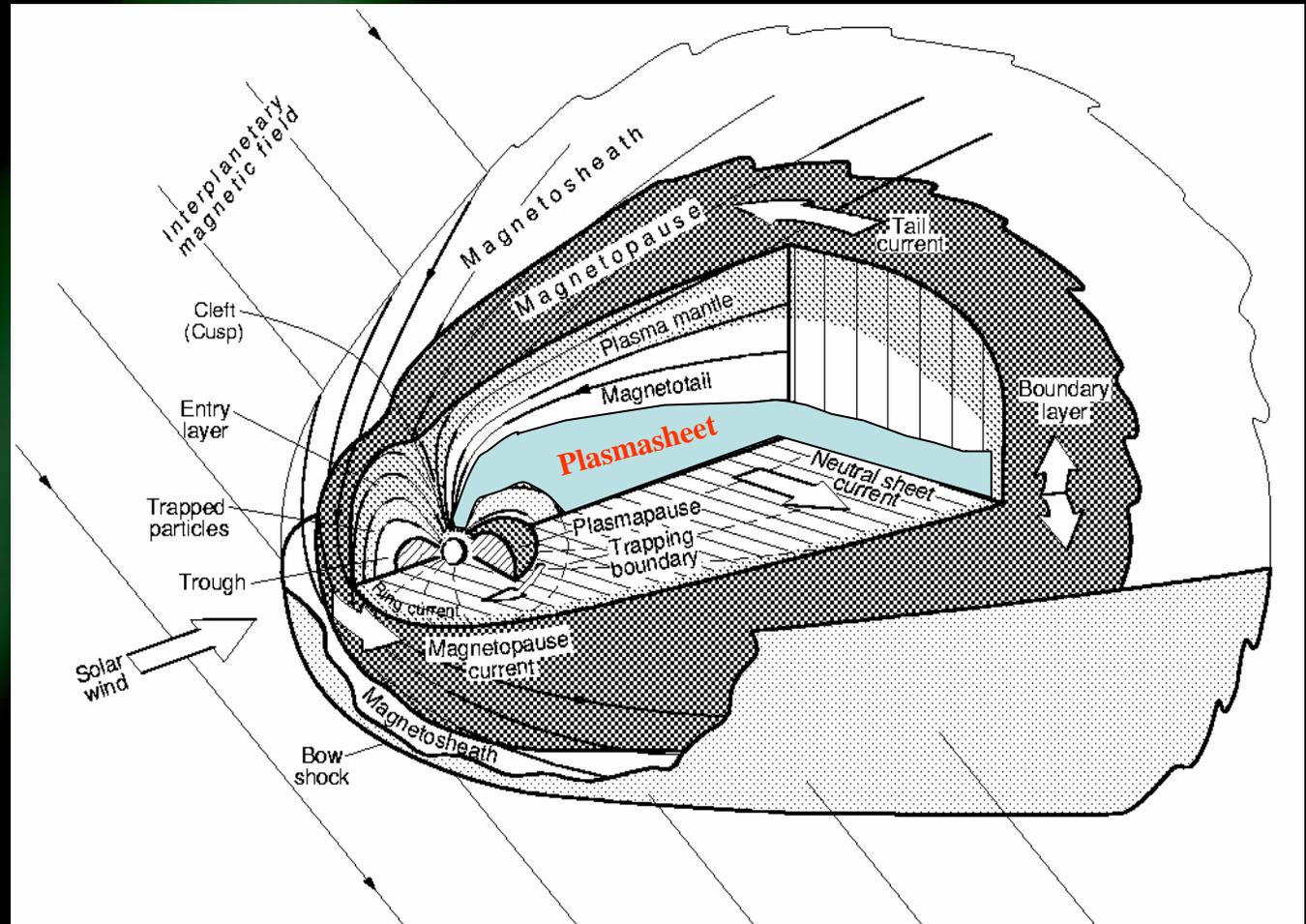


Polar

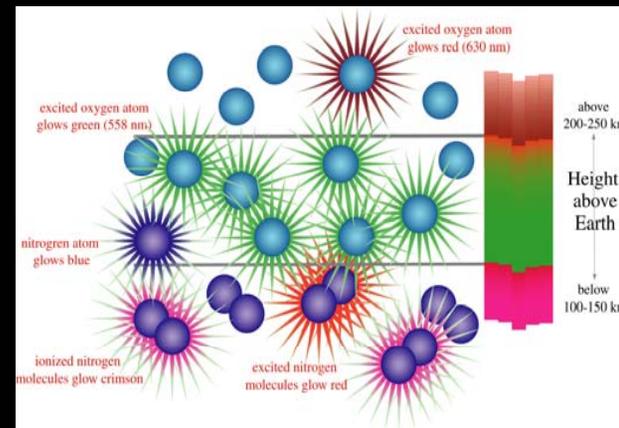
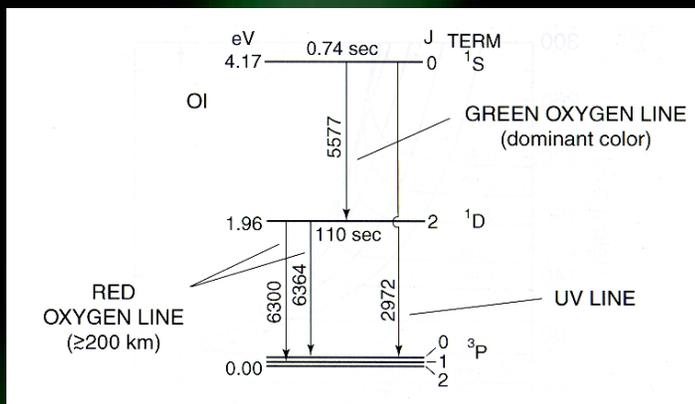
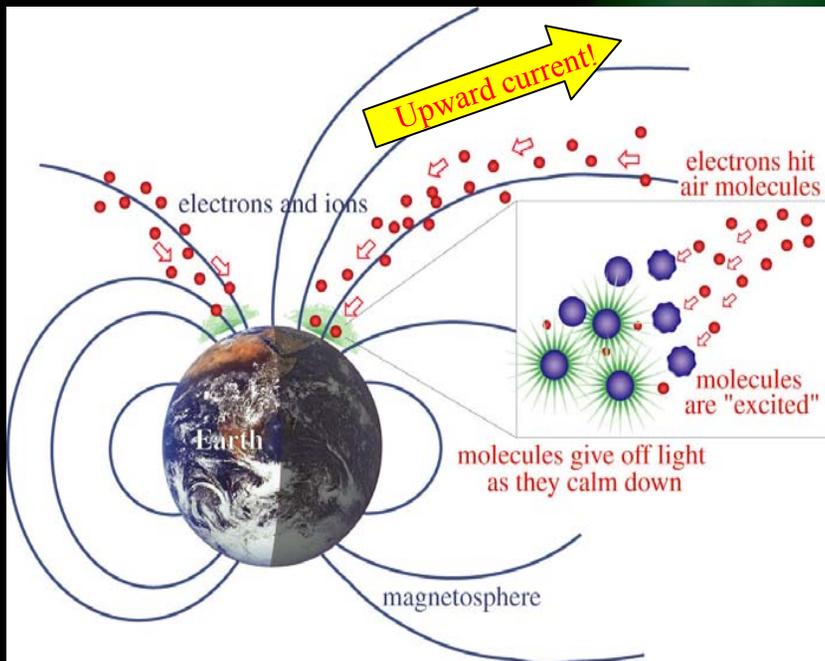


Dynamics Explorer

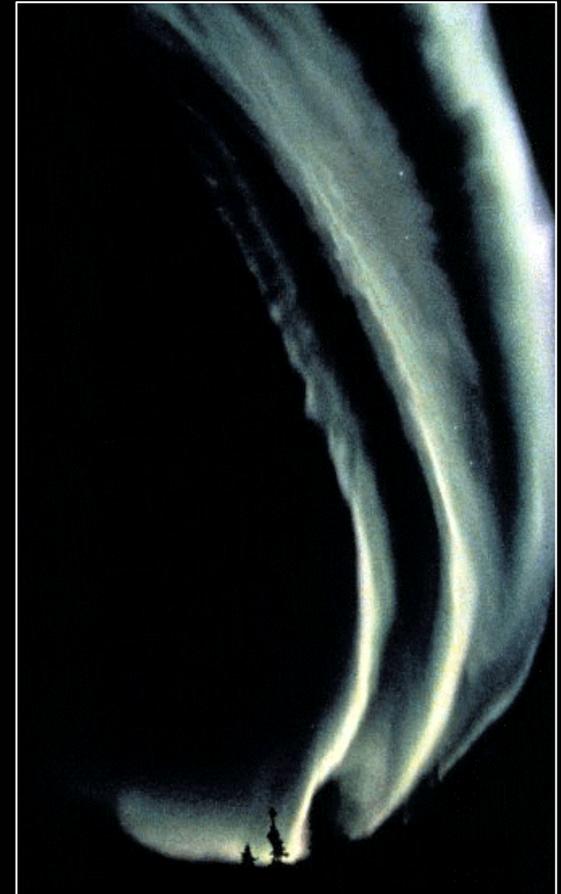
The auroral oval is the projection of the plasmasheet onto the atmosphere



Auroral emissions

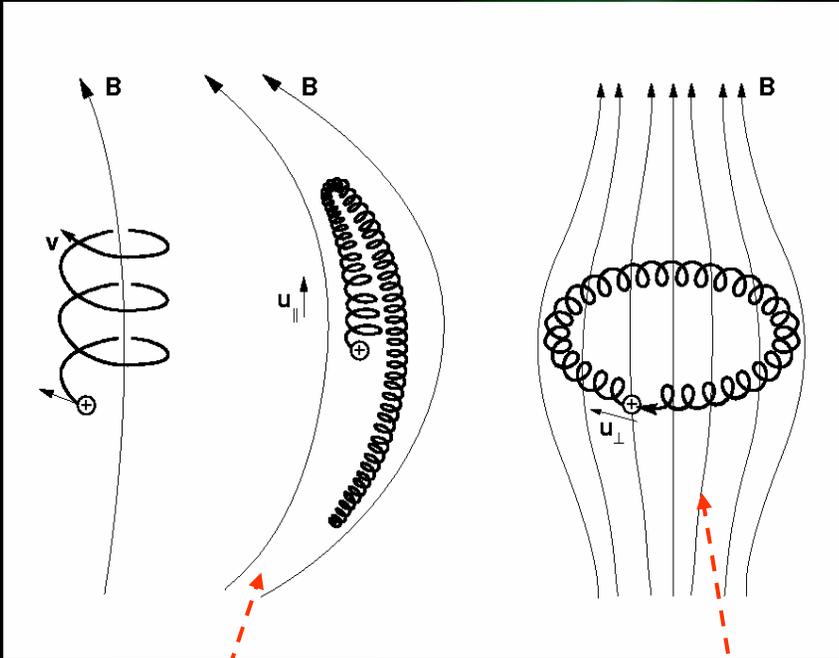


Homogenous auroral arcs



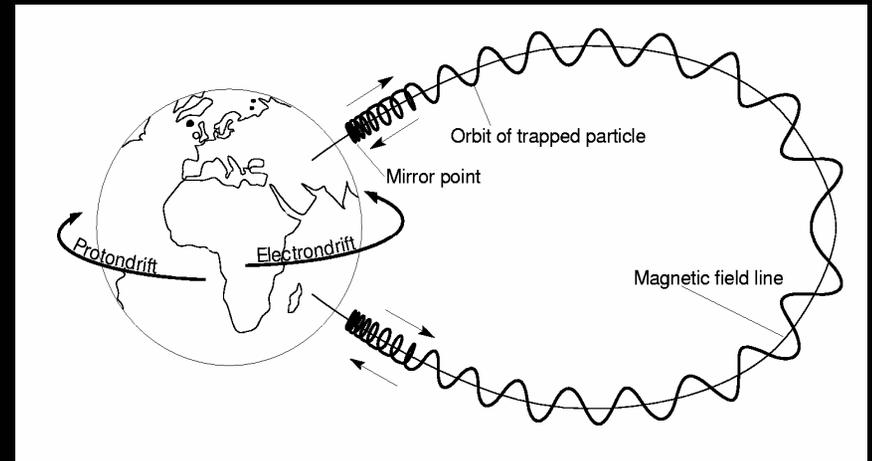
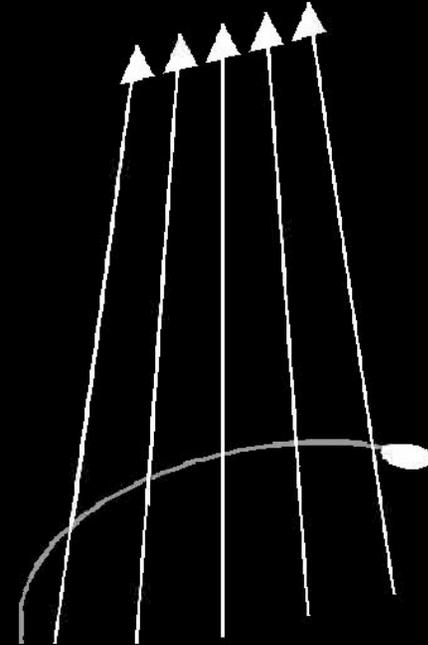
Particle motion in the geomagnetic field

gyration longitudinal oscillation azimuthal drift

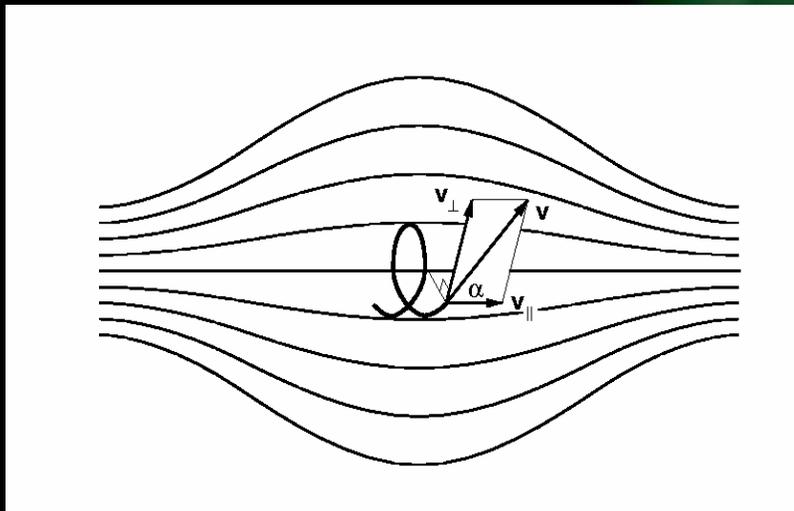


Magnetic mirror

grad B drift



Magnetic mirror



$mv^2/2$ constant ➔

$$\frac{\sin^2 \alpha}{B} = \text{const}$$

particle turns when $\alpha = 90^\circ$ ➔

$$B_{\text{turn}} = B / \sin^2 \alpha$$

If maximal B-field is B_{max} a particle with pitch angle α can only be turned around if

$$B_{\text{turn}} = B / \sin^2 \alpha \leq B_{\text{max}} \quad \text{➔}$$

The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B / B_{\text{max}}}$$

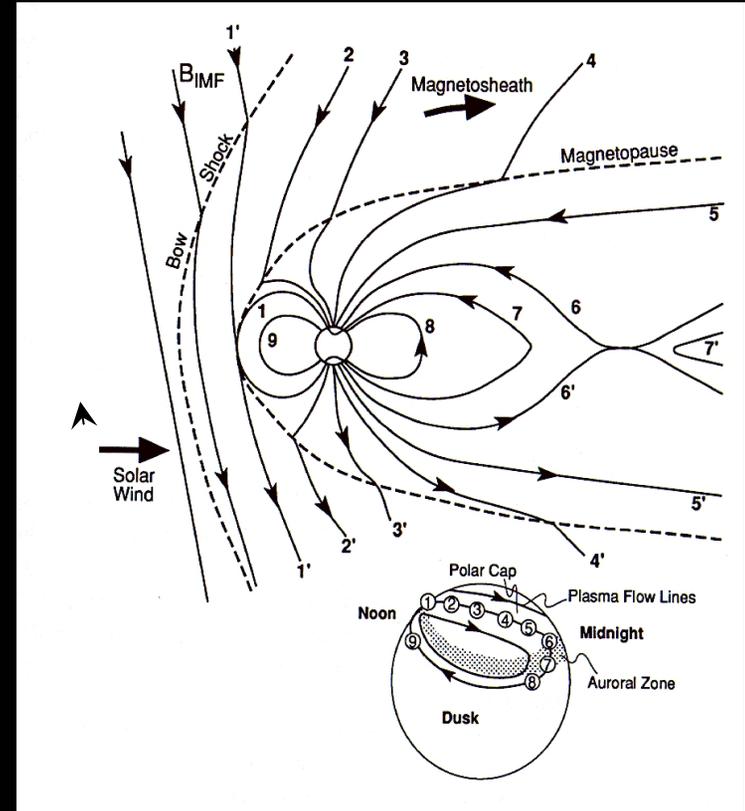
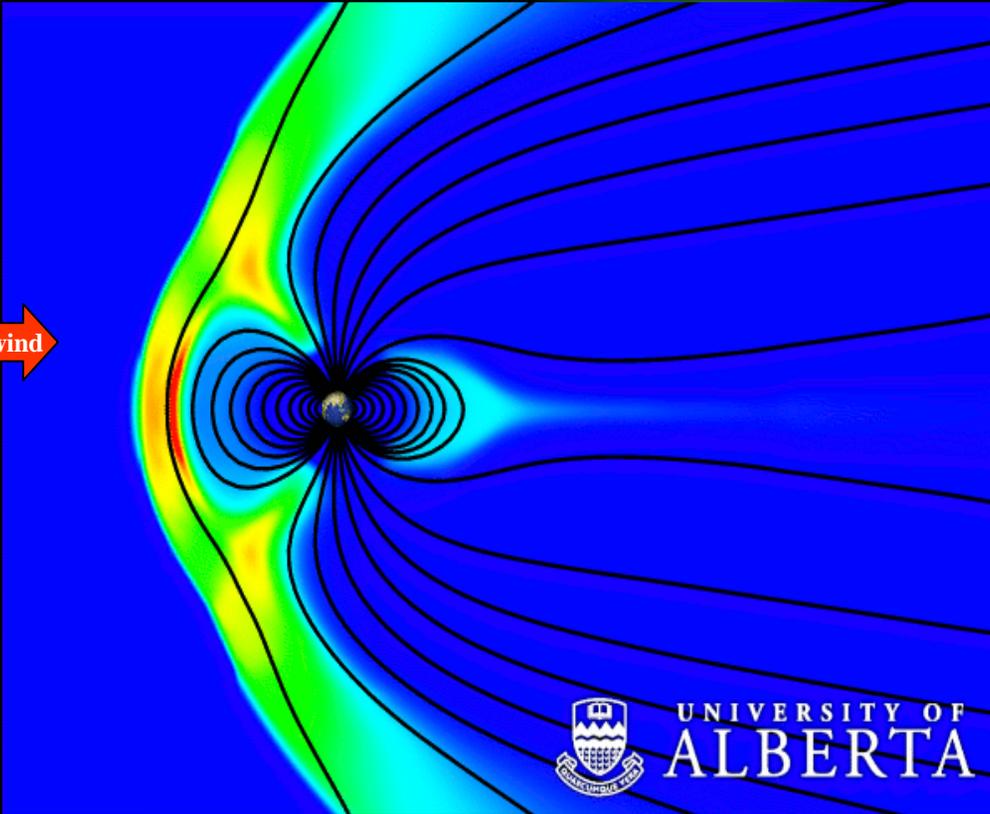
Particles in
loss cone :

$$\alpha < \alpha_{lc}$$

Magnetospheric convection

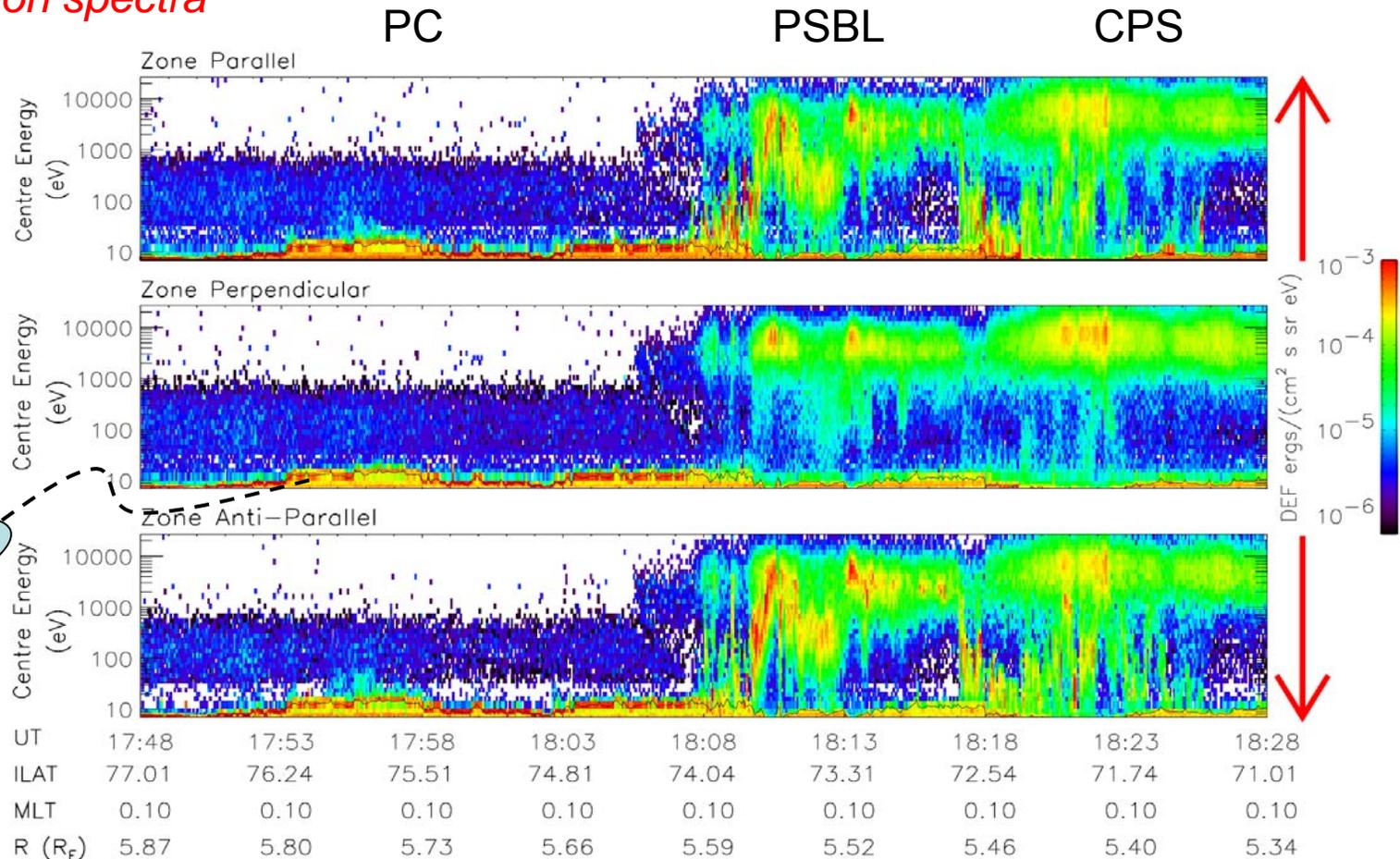


Solar wind



A typical auroral pass - CLUSTER

Electron spectra



Auroral scales

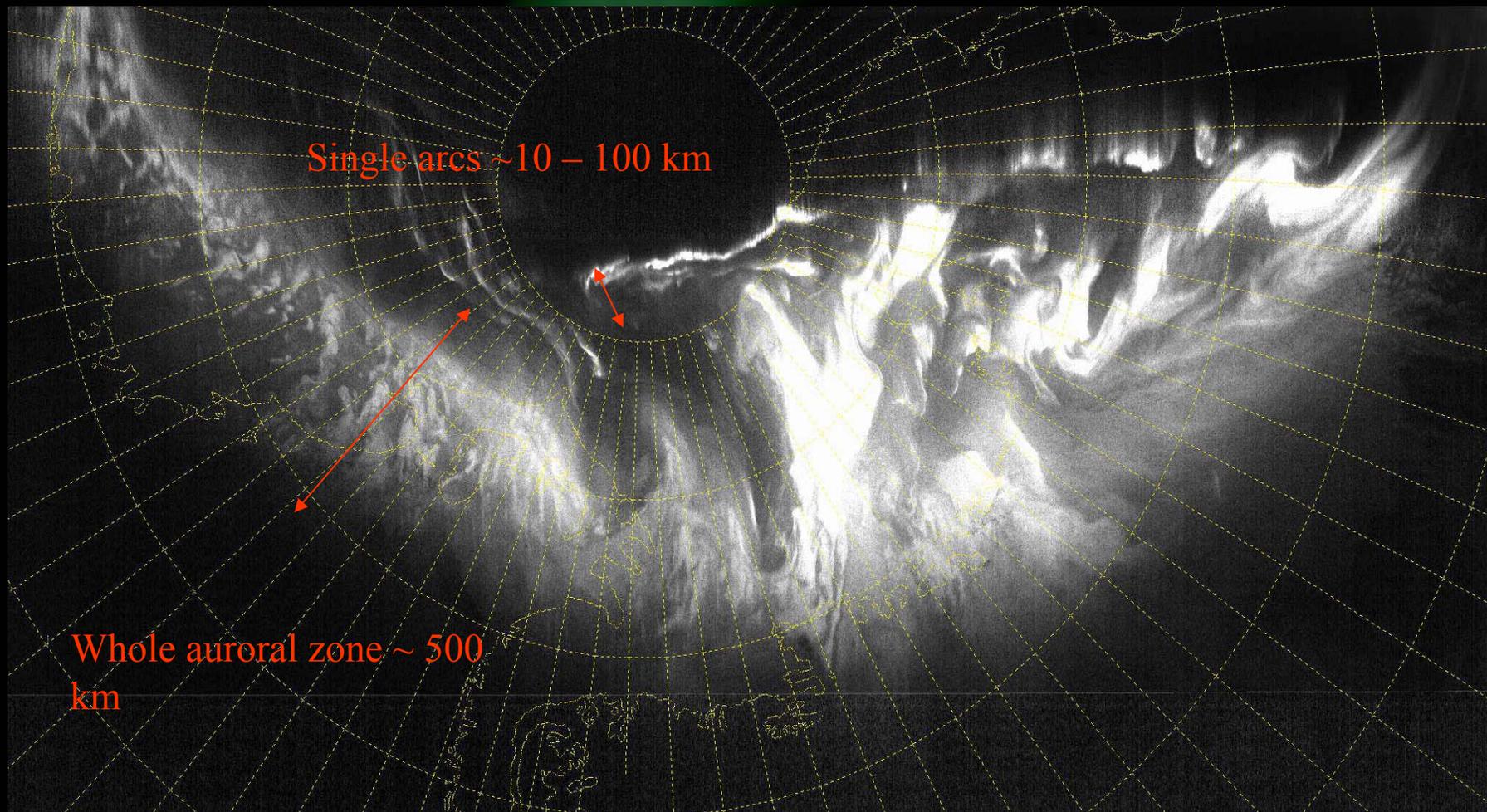
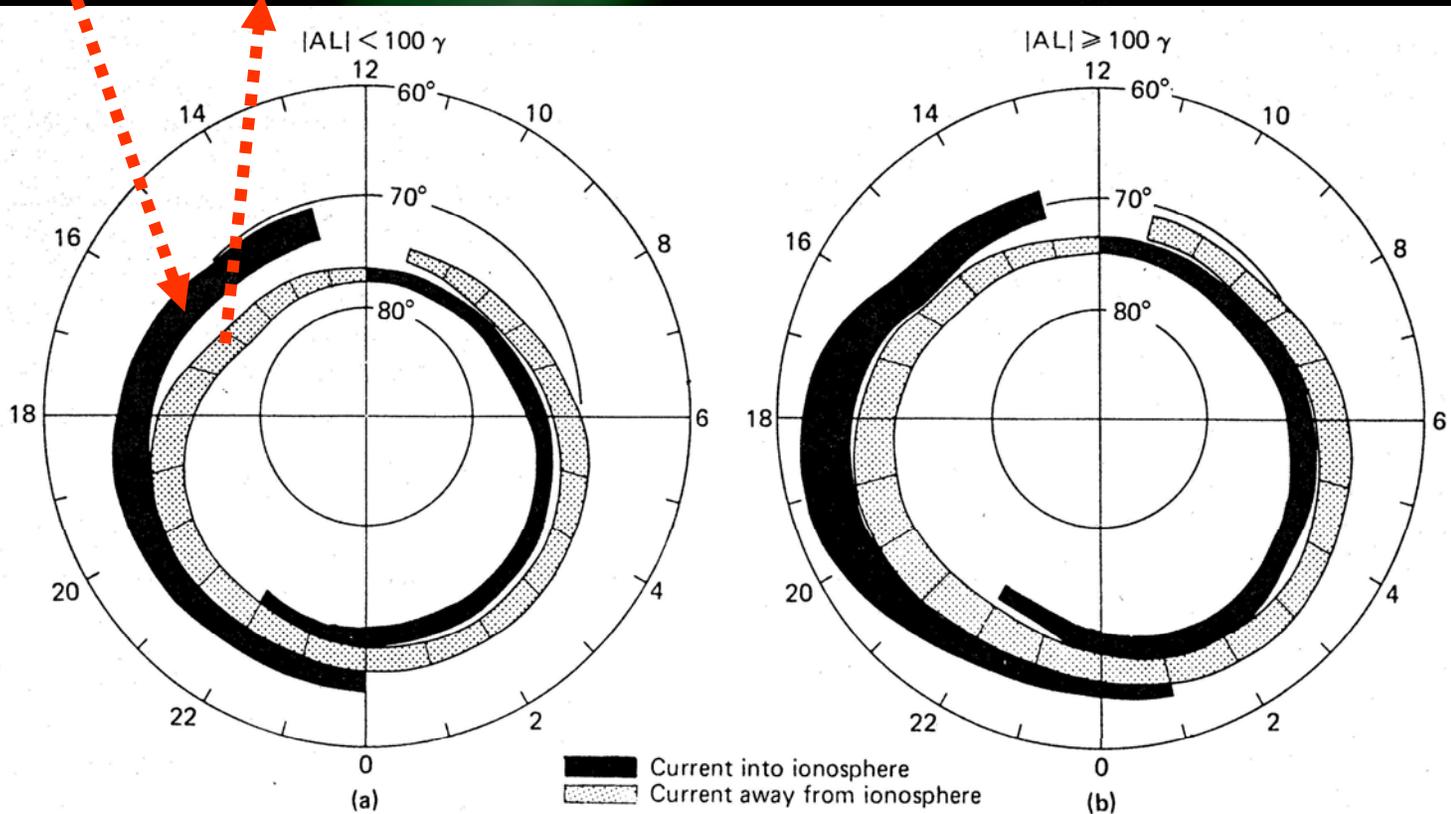


Photo from DMSP satellite

Birkeland currents in the auroral oval

Low geomagnetic activity

High geomagnetic activity



A typical auroral pass - FAST

ΔB_{EW}

↓ dE-flux, e^-

PA, e^-

↑ dE-flux, e^-

n-flux, e^-

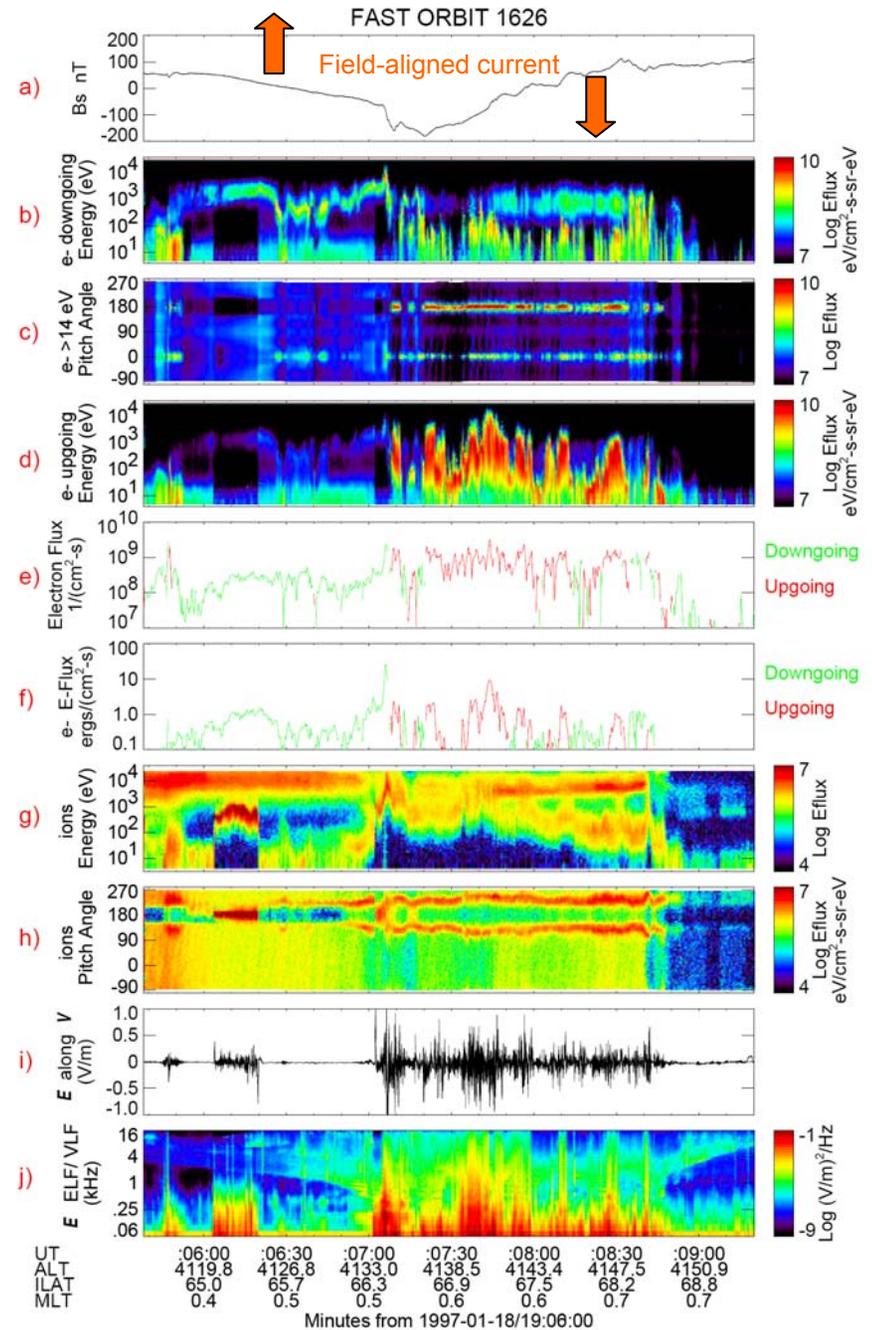
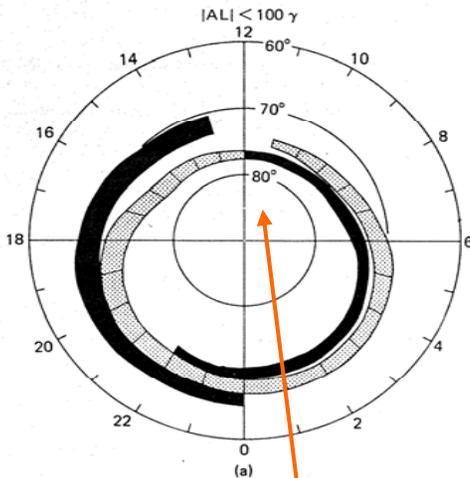
E-flux, e^-

dE-flux, i^+

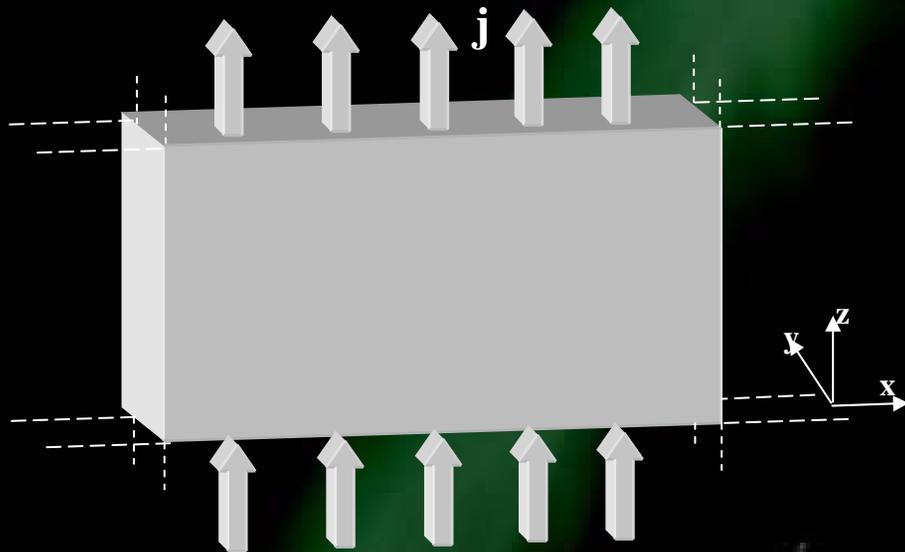
PA, i^+

E-perp (DC)

E (AC)



Current sheet approximation and Ampère's law



$$\left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z}, \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}, \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) = \mu_0 (j_x, j_y, j_z)$$

But $\frac{\partial}{\partial x} = 0$ and $\frac{\partial}{\partial z} = 0$

$$\left(\frac{\partial B_z}{\partial y}, 0, -\frac{\partial B_x}{\partial y} \right) = \mu_0 (0, 0, j_z)$$

Ampère's law (no time dependence):

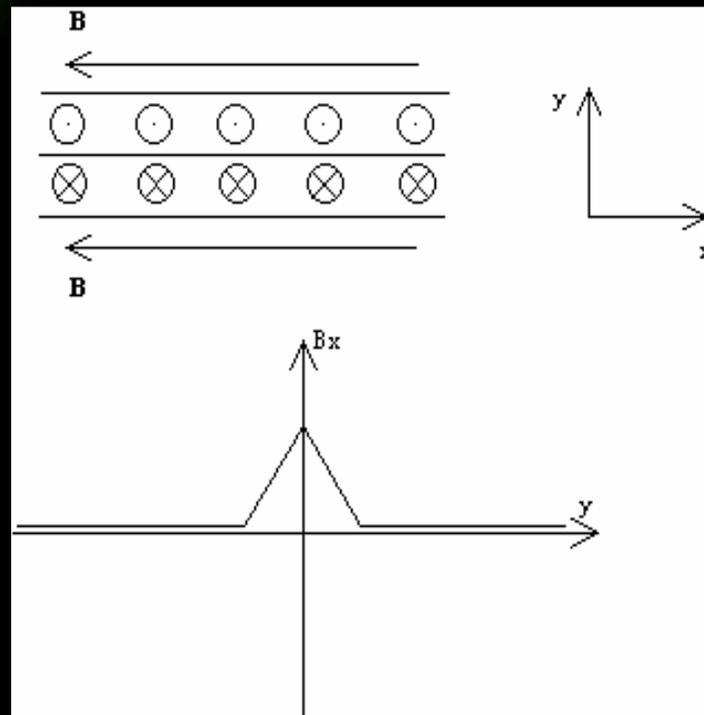
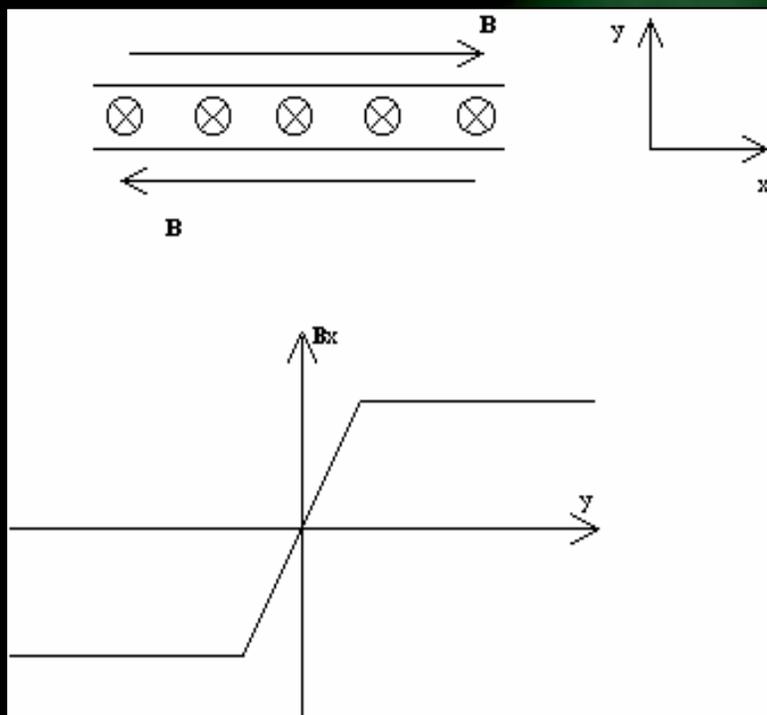
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

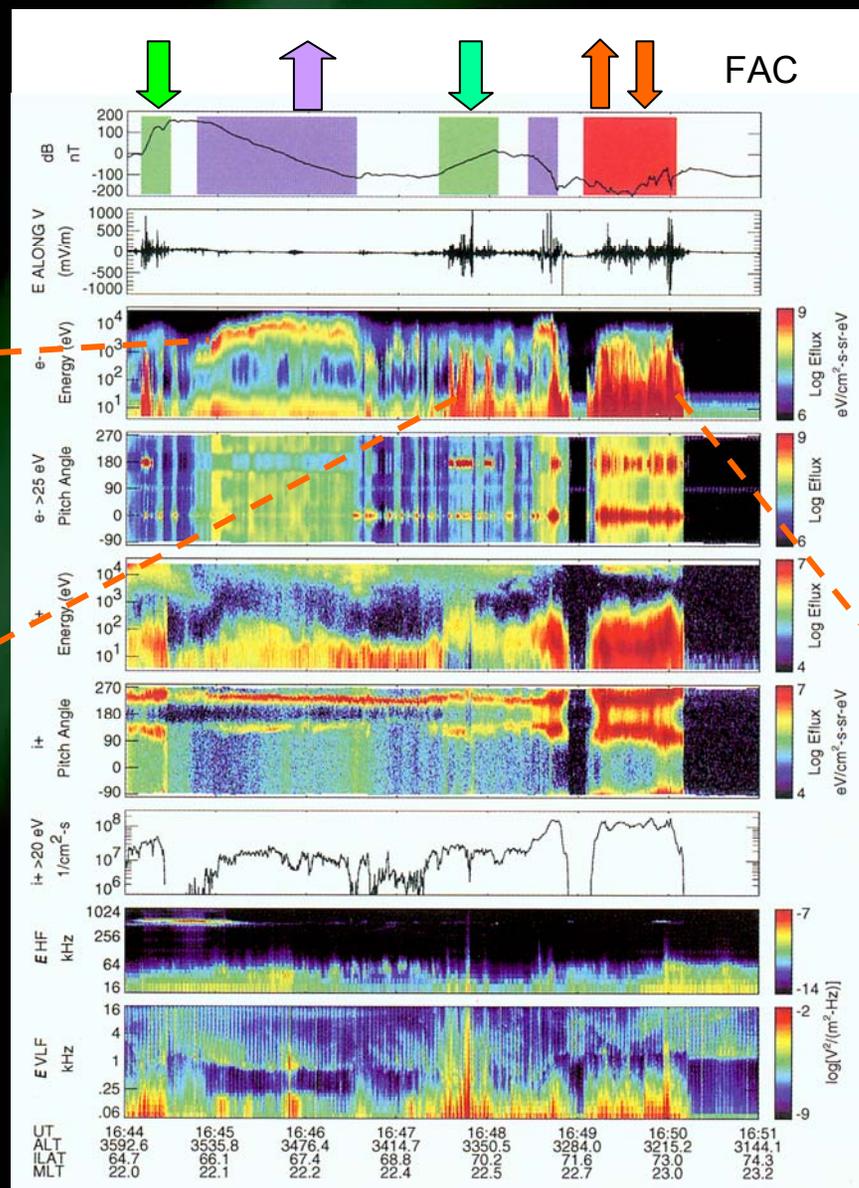
Current sheet

Determination of current density by magnetic field measurement



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

Upward and downward current regions



Inverted V

Upgoing electron beams

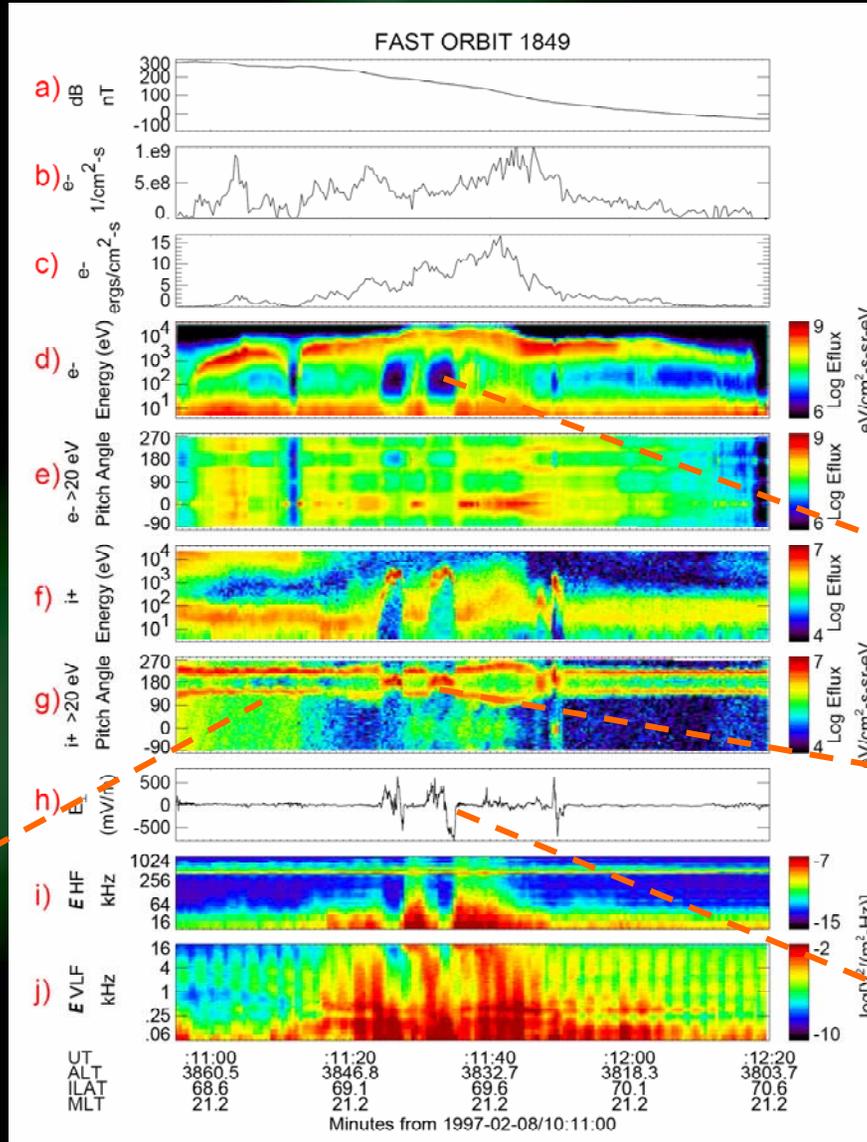
180° = ↑

0° = ↓

Alfvénic aurora

Upward current region

Inverted V arc



180° = ↑

0° = ↓

Ion conic

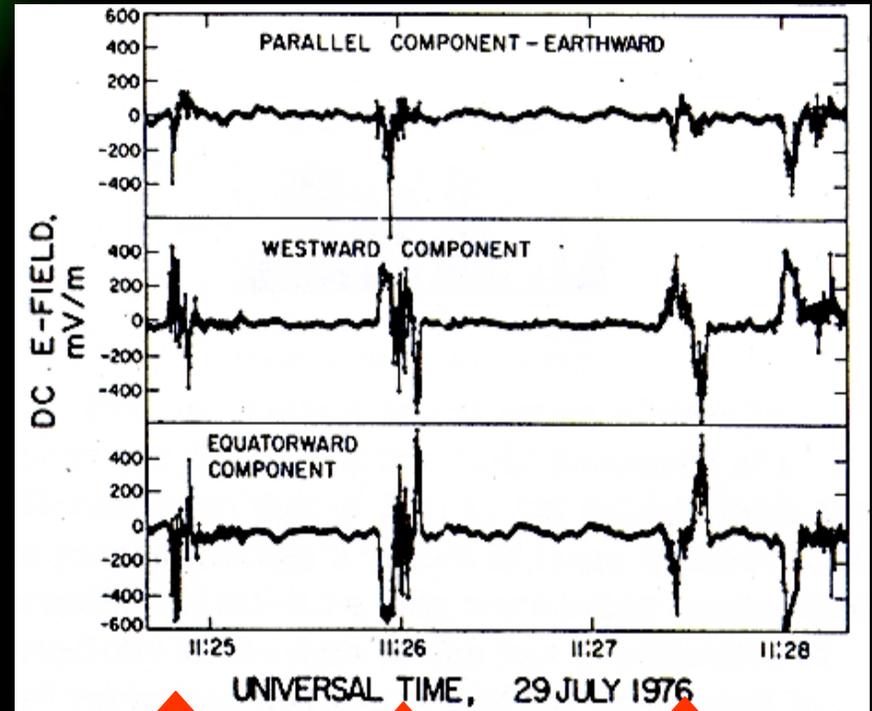
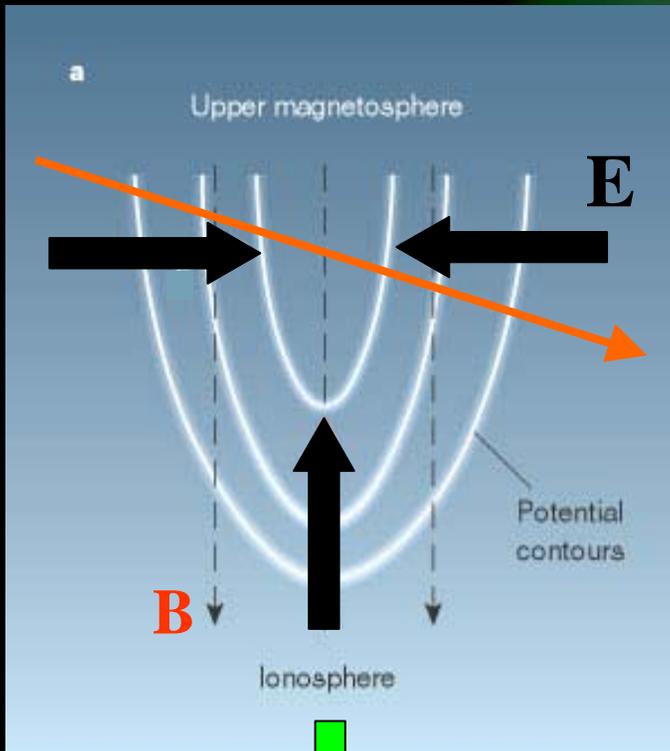
Visual limit $\sim 10^{-3} \text{ Jm}^{-2}\text{s}^{-1}$
 $= 1 \text{ erg cm}^{-2}\text{s}^{-1}$

Auroral density cavity

Ion beam

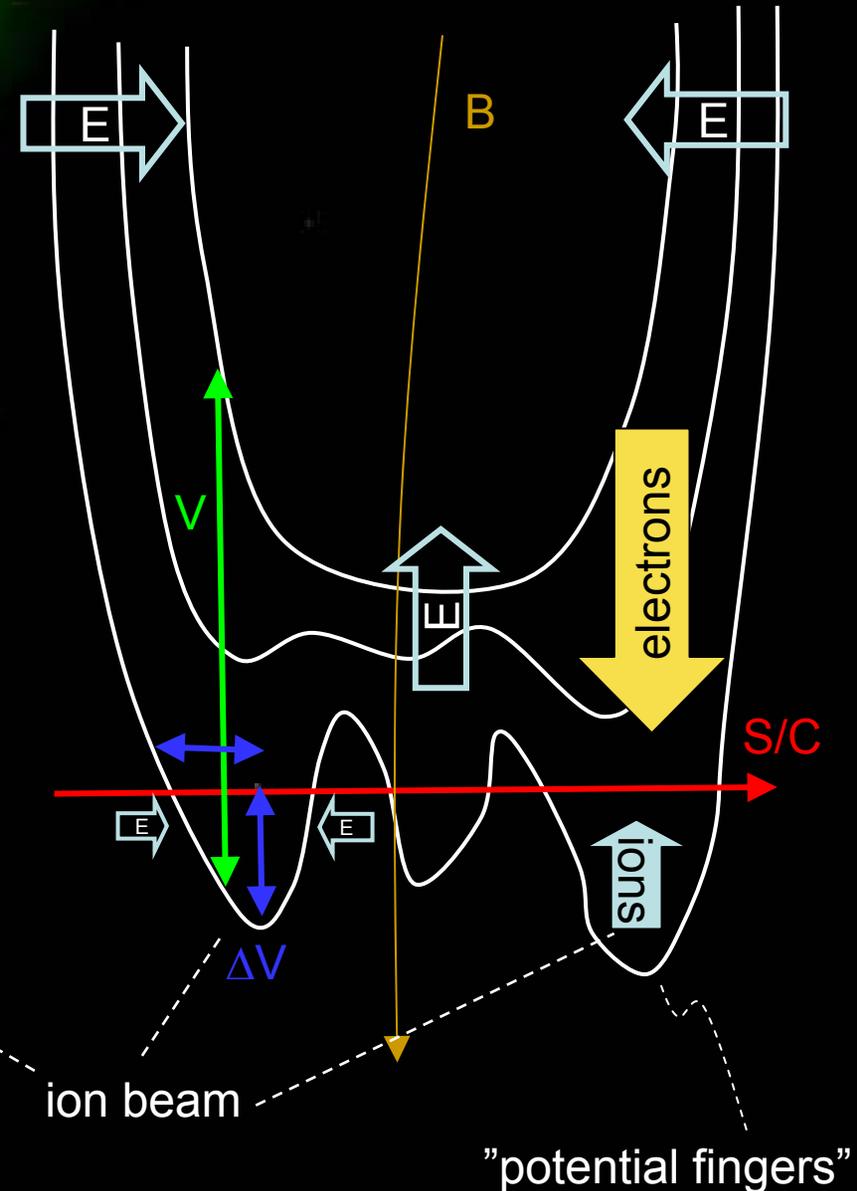
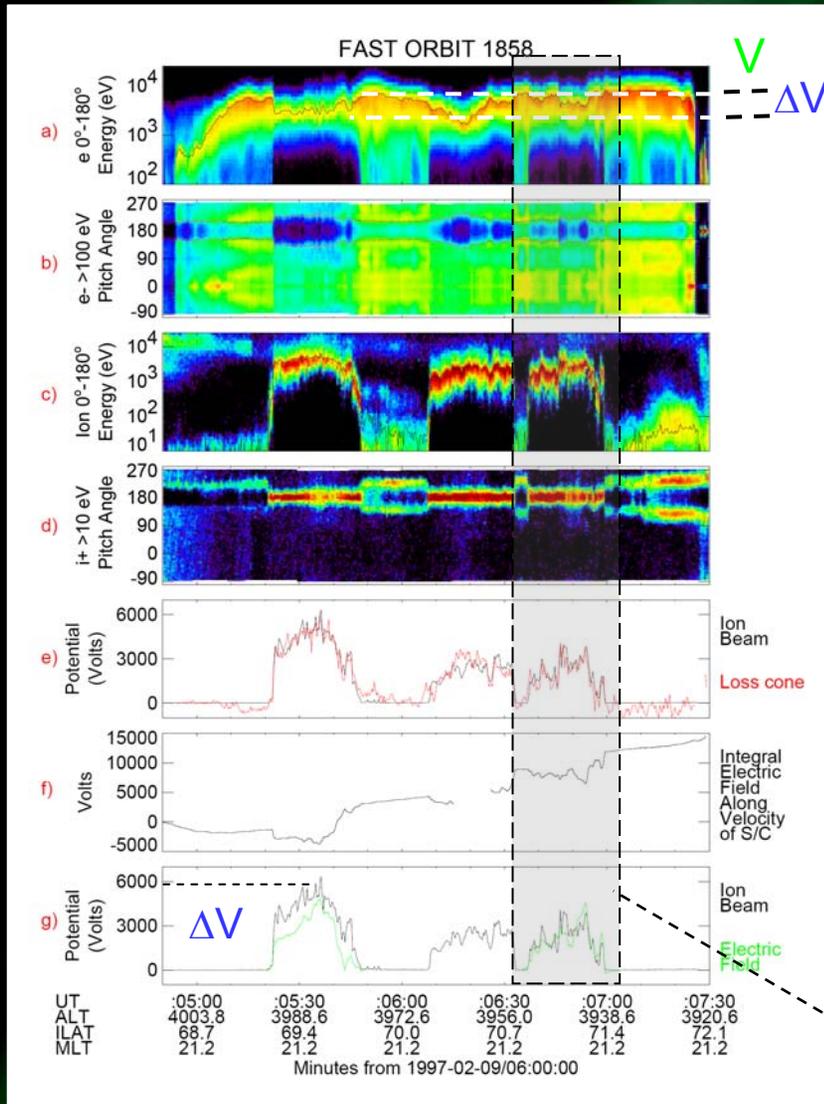
Bipolar electric field structures

Satellite signatures of U potential

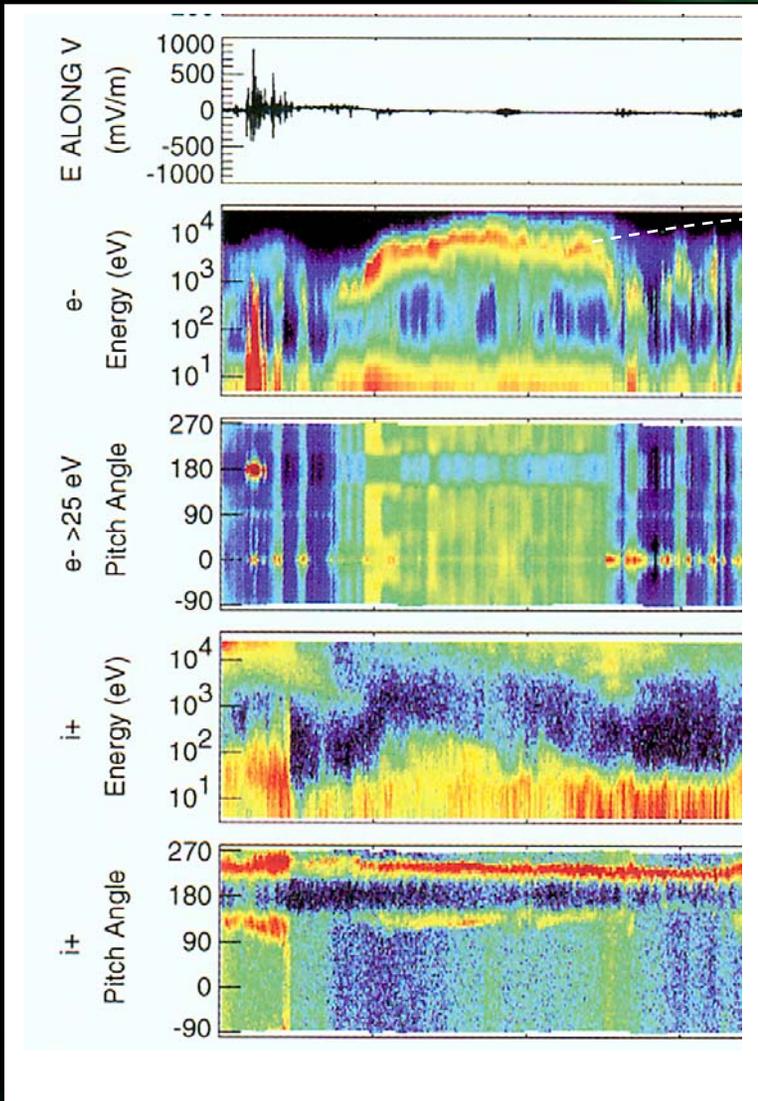


Measurements made by the ISEE satellite (Mozer et al., 1977)

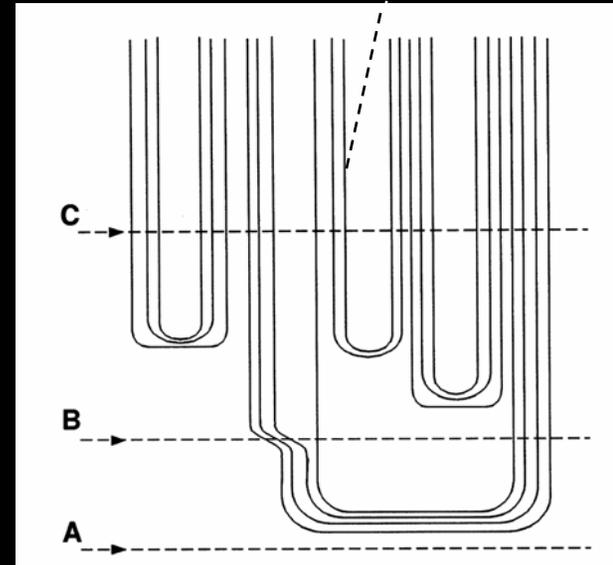
Acceleration potential structure I



Acceleration potential structure II

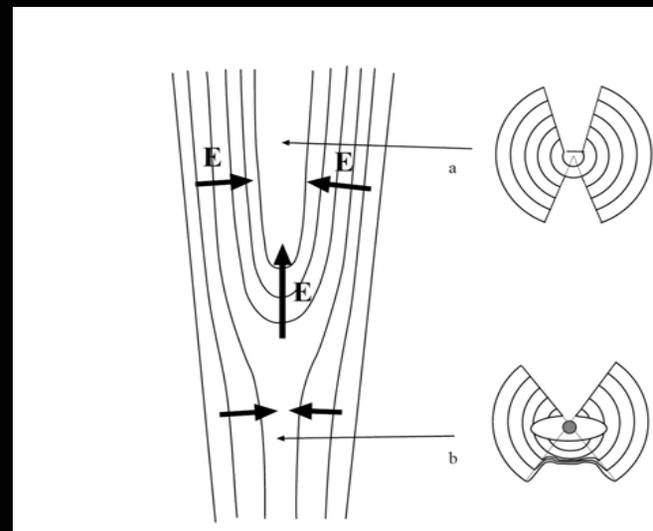
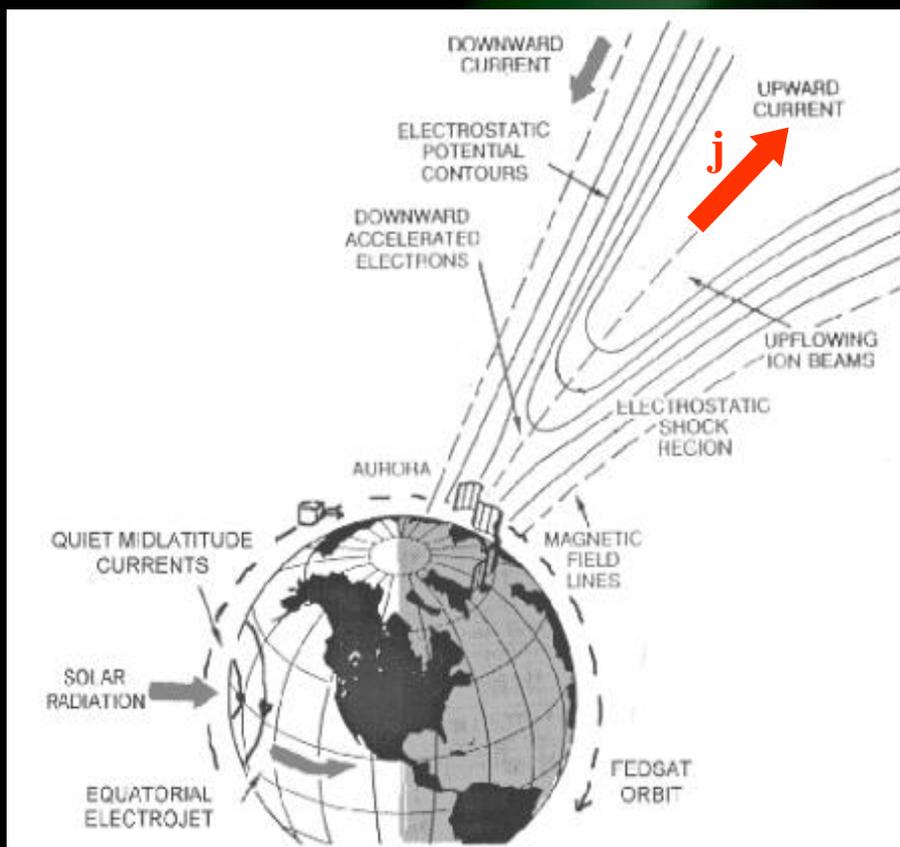


Localized regions of higher energy electrons without associated ion beams

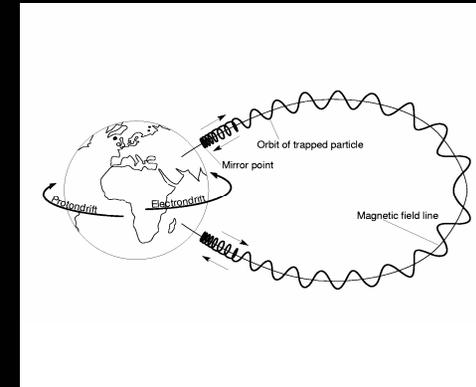
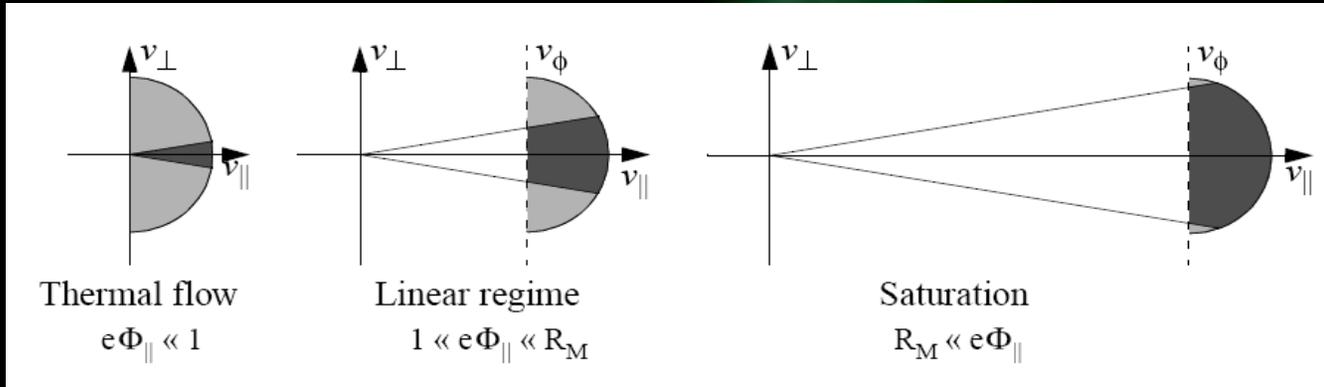


Why particle acceleration?

- The magnetosphere often acts as a current generator
- Electrons are accelerated downwards by upward E-field.
- This increases the pitch-angle of the electrons, and more electrons can reach the ionosphere, where the current can be closed.



Auroral currents – Knight relation



Fraction of particles in the loss cone:

$$f = \frac{\pi\theta_{lc}}{2\pi} \sim \frac{B_{ms}}{B_{ion}} \sim 2 \times 10^{-4}$$

Thermal current:

$$j_{\parallel,ms} = n_0 e v_{th} f$$

$$j_{\parallel,ion} = n_0 e v_{th} f \frac{B_{ms}}{B_{ion}} =$$

$$= n_0 e v_{th} \frac{B_{ms}}{B_{ion}} \frac{B_{ion}}{B_{ms}} = n_0 e v_{th} \approx$$

$$[n_e = 0.1 \text{ cm}^{-3}, T_e = 1 \text{ keV}] \approx$$

$$\sim 1 \mu\text{A/m}^2$$

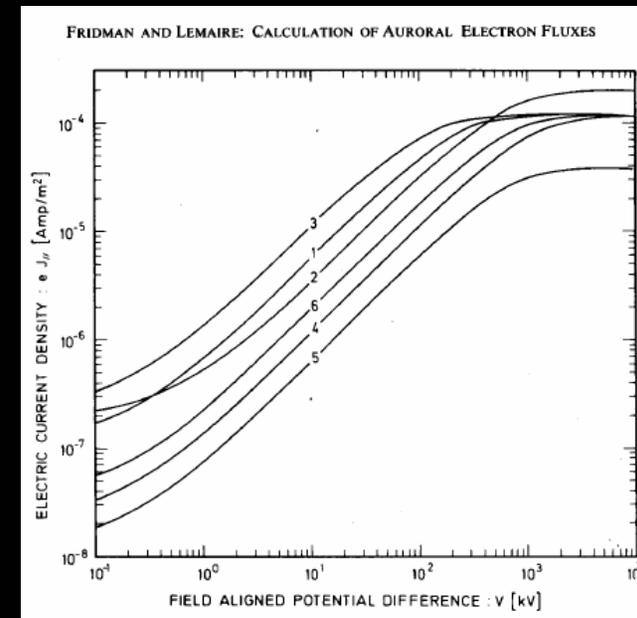
Apply a parallel potential drop:

$$j_{\parallel,ion} = n_0 e v_{th} \frac{B_{ion}}{B_{ms}} \left[1 - \frac{e^{-xe\Phi_{\parallel}/T_e}}{1+x} \right]$$

Linear regime :

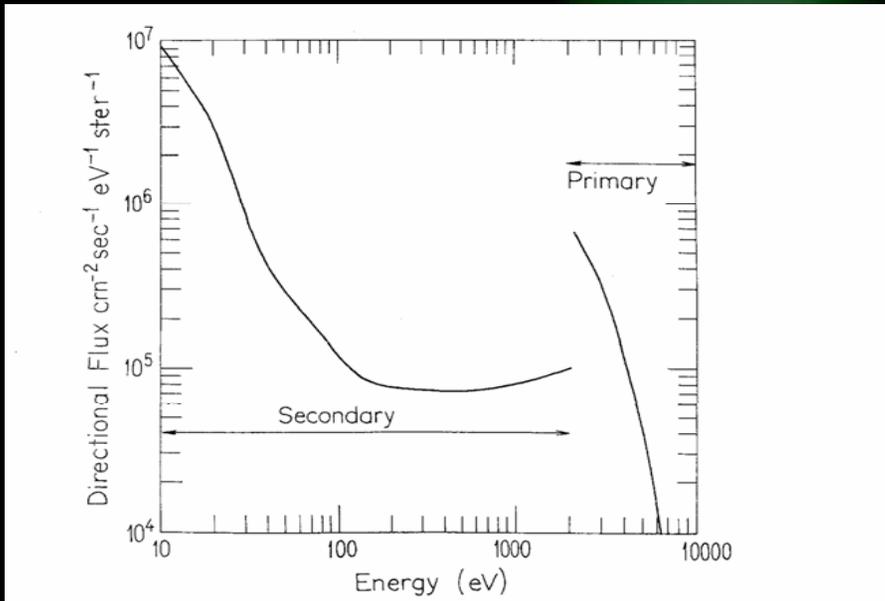
$$j_{\parallel,ion} \approx n_0 e v_{th} \frac{e\Phi_{\parallel}}{k_B T_e} = K \Phi_{\parallel}$$

$$K = \frac{e^2 n_0}{\sqrt{2\pi m_e k_B T_e}} \sim 10^{-9} \text{ S/m}^2$$

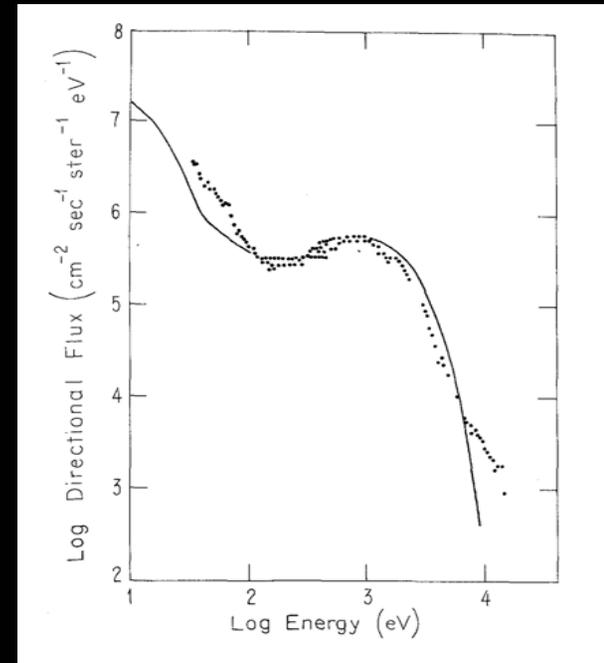


$$v_{th} = \sqrt{T_e / 2\pi m_e} \quad x = \frac{1}{B_i / B_0 - 1}$$

Particle distributions associated with inverted V's

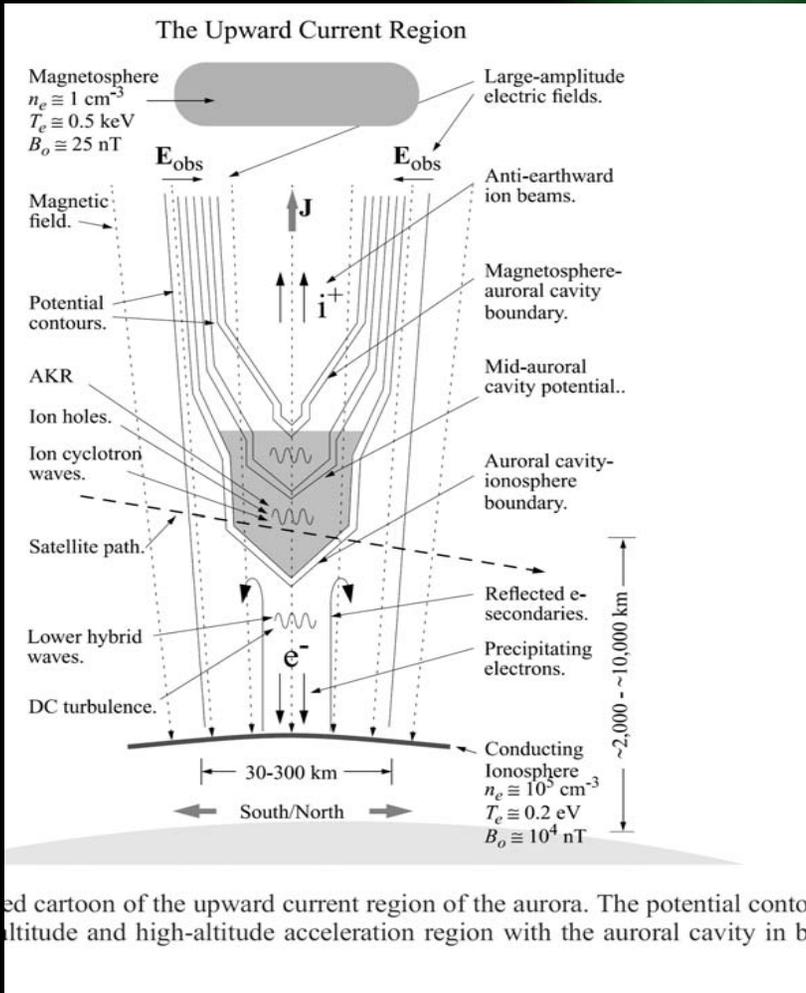


Model of cold beam producing secondaries (Evans, 1974)

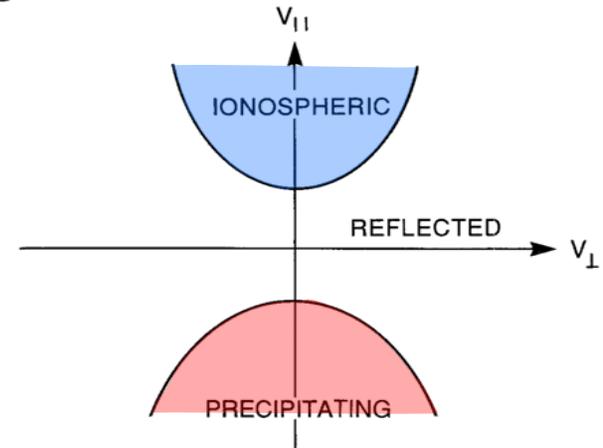


Model of hot electron beam and secondaries (Evans, 1974). Data from Franck and Ackerson, 1971

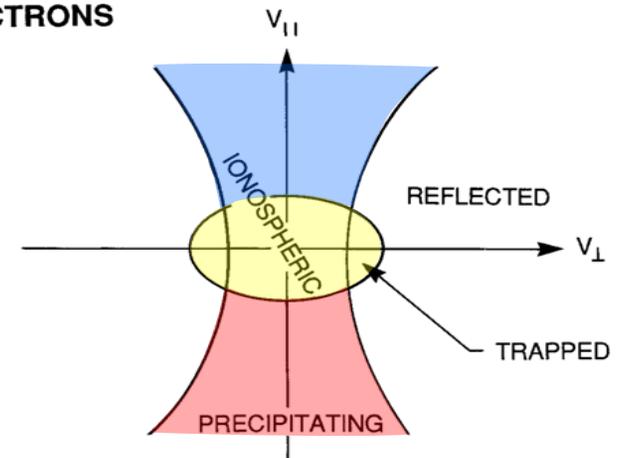
Auroral cavity and trapped populations



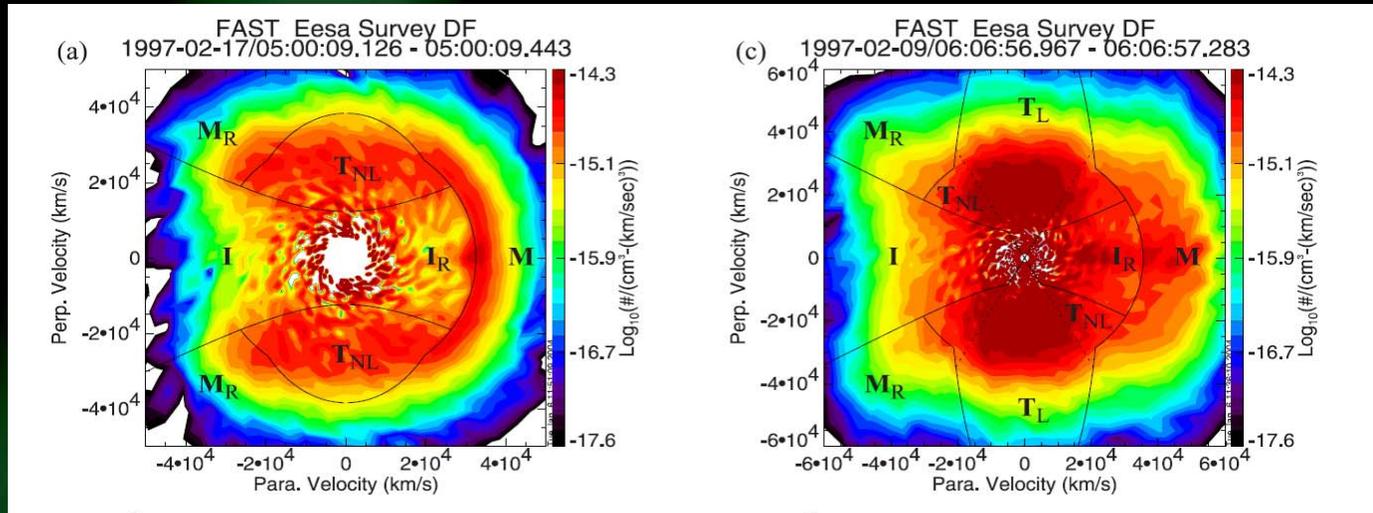
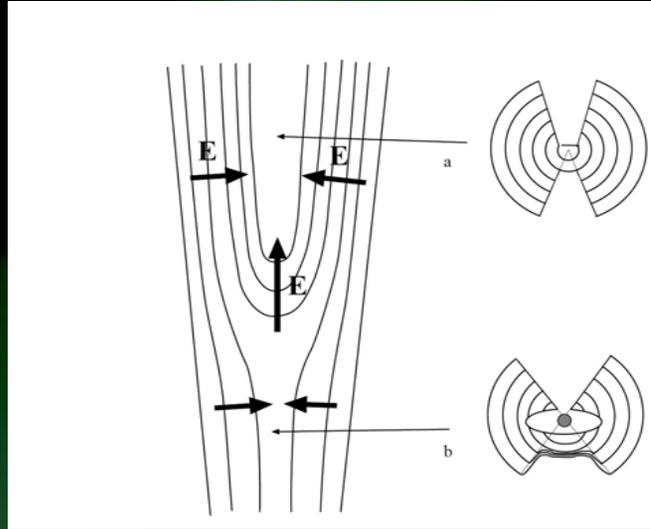
(a) IONS



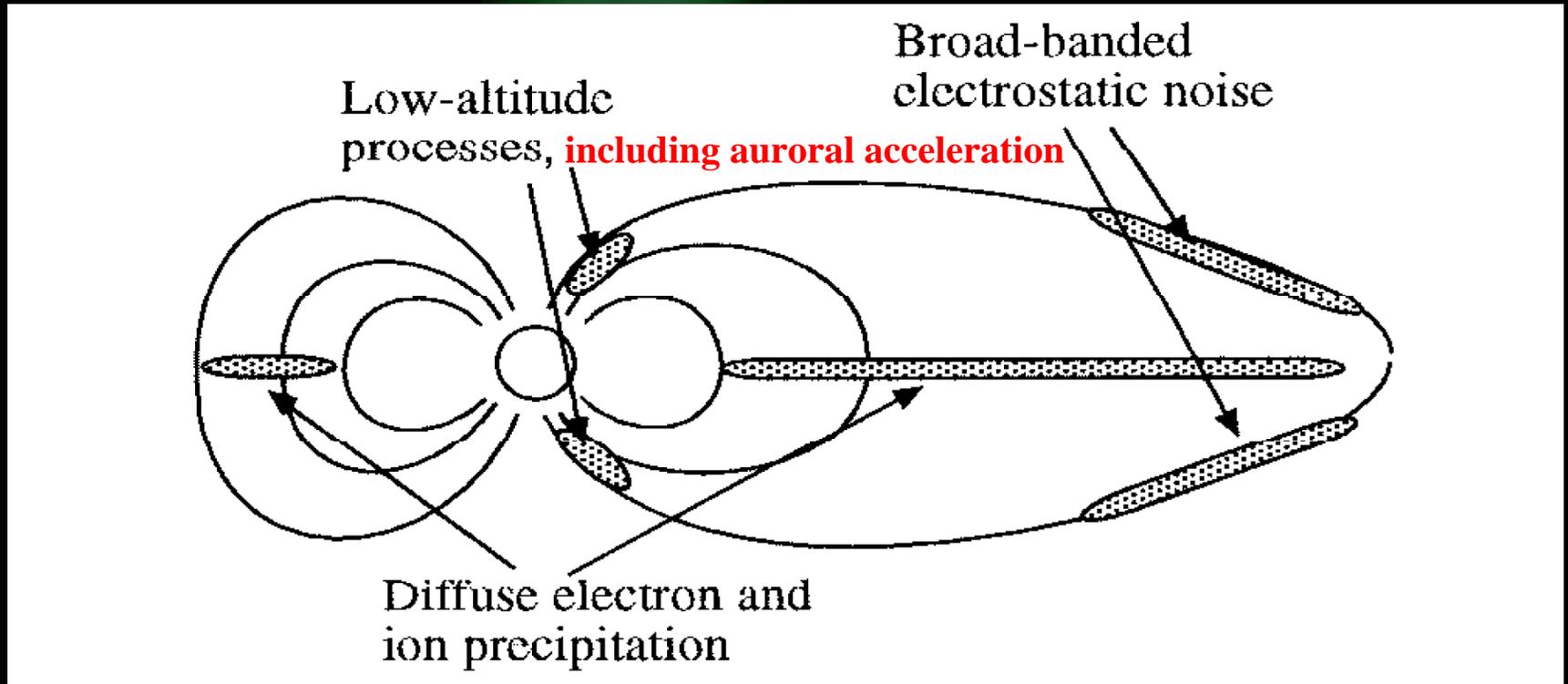
(b) ELECTRONS



Accelerated Maxwellian



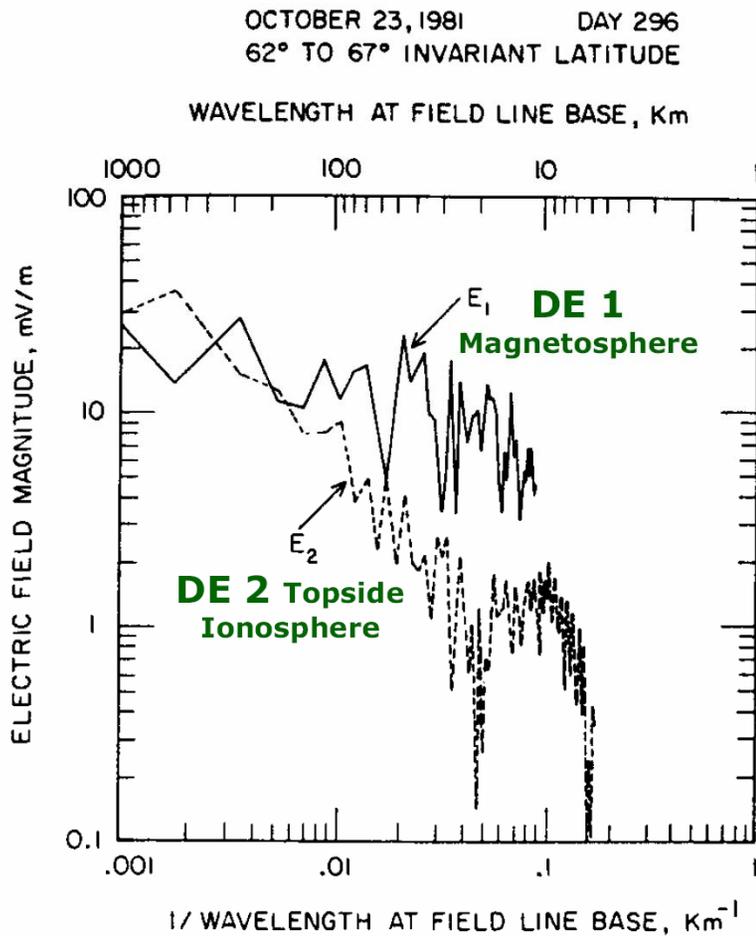
Acceleration regions



Koskinen

Auroral acceleration region typically situated at altitude of 1-3 R_E

Mapping of auroral electric fields



Large scales

Small scales

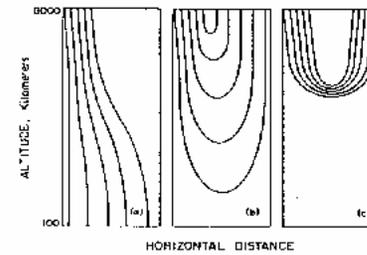


Figura 8. Contururi echipotentiale posibile in zona de accelerare (v. text). Din Hudson and Mozer (1978)

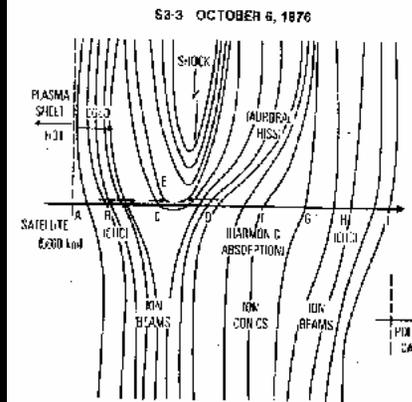
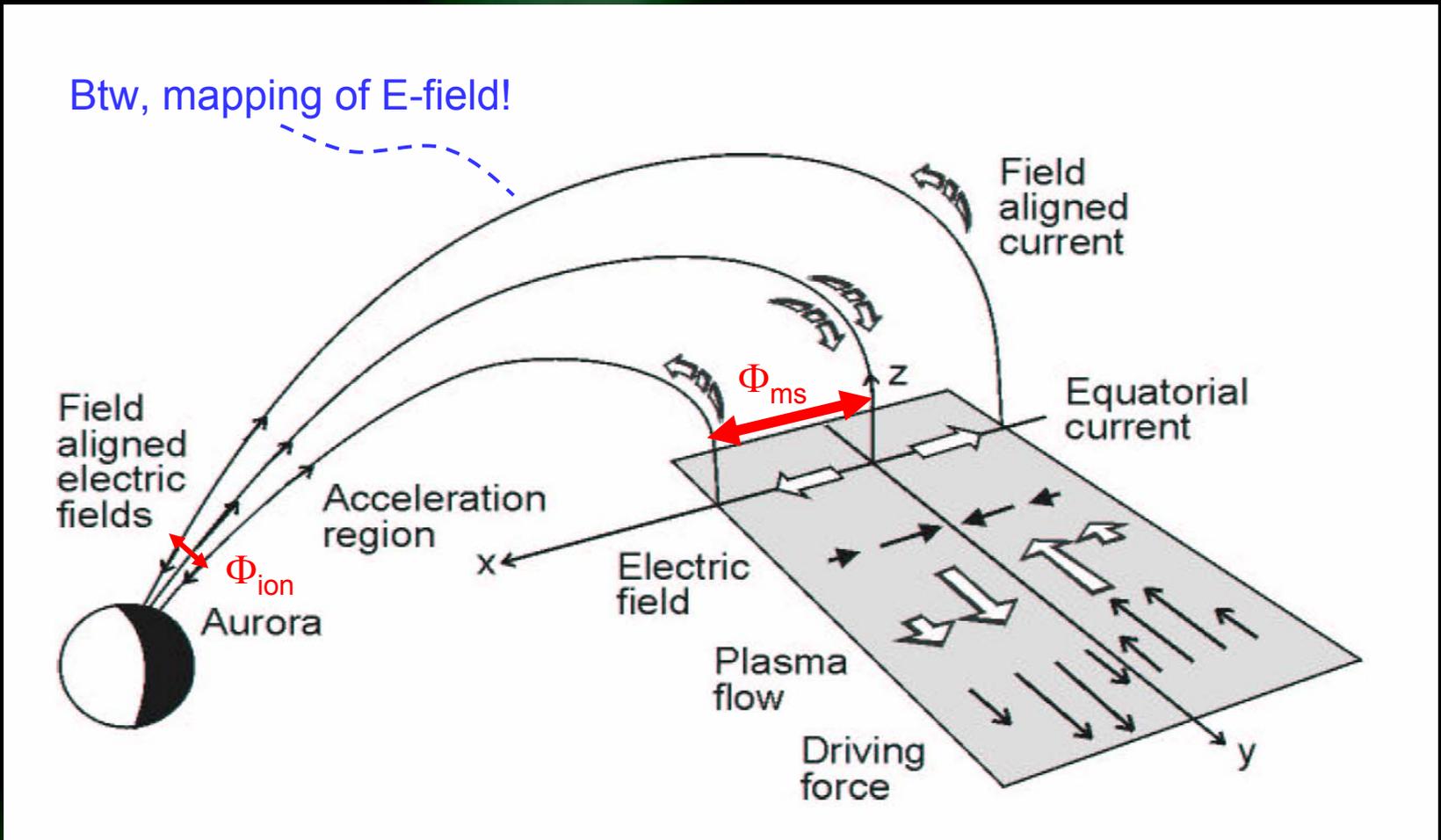


Figura 9. Reconstructia calitativa a contururilor echipotentiale din date 53-3. Figura ilustreaza 1 minut de observatii. Din Mizera et al. (1982)

Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (Weimer et al., 1985)

Static, medium-scale MI-coupling

MI-coupling critical scale size II



Static, medium-scale MI-coupling

MI-coupling critical scale size III

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp} + \mathbf{E}_{\perp} \cdot \nabla_{\perp} \Sigma_P + (\mathbf{E}_{\perp} \times \nabla_{\perp} \Sigma_H) \cdot \hat{\mathbf{b}}$$

$$K(\Phi_{ms} - \Phi_{ion}) \approx \Sigma_P \frac{\Phi_{ion}}{L^2}$$

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp}$$

$$\Phi_{ion} = \left(1 + \frac{\Sigma_P}{KL^2} \right)^{-1} \Phi_{ms}$$

$$j_{\parallel} = K \Delta \Phi_{\parallel} = K(\Phi_{ms} - \Phi_{ion})$$

$$\text{When } L \gg \sqrt{\frac{\Sigma_P}{K}} : \Phi_{ion} \approx \Phi_{ms}$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp ion}$$

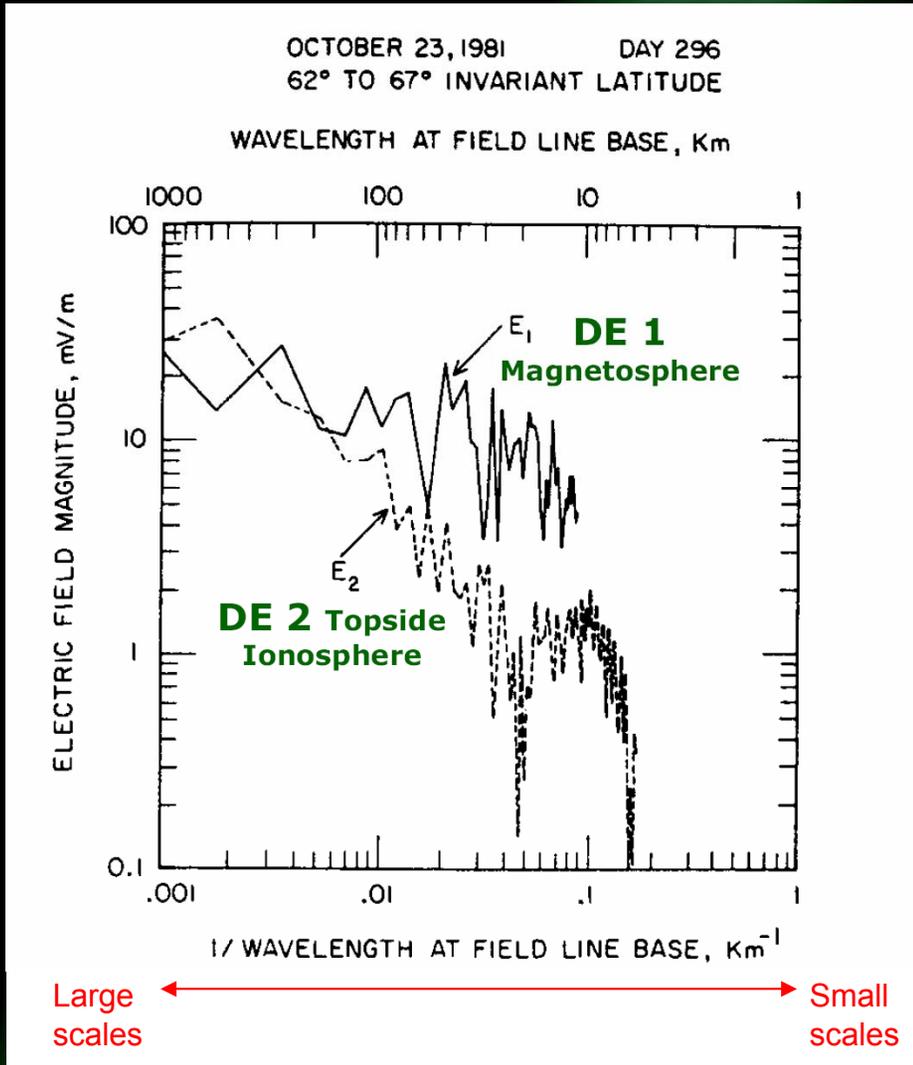
$$\text{When } L \ll \sqrt{\frac{\Sigma_P}{K}} : \Phi_{ion} \ll \Phi_{ms}$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp}^2 \Phi_{ion}$$

The last case means $|\Phi_{ms}| \approx |\Delta \Phi_{\parallel}|$

Static, medium-scale MI-coupling

MI-coupling critical scale size



Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (*Weimer et al., 1985*)

1/WAVELENGTH AT FIELD LINE BASE, Km⁻¹

Fig. 4. Electric field spectrums from day 296 (October 23) of 1981. The spectrums are obtained from a Fourier transform of the electric field data between 62° and 67° invariant latitude. The solid line shows the spectrum of the electric field measured by DE 1. The solid line shows the spectrum of the electric field measured by DE 2. The ordinate values are obtained from the square root of the "spectral power density." The actual units are mV m⁻¹ km^{1/2}. *Weimer et al, 1985*