



Auroral processes in satellite data



© *AURORA EXPERIENCE*

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Overview

- Setting the scene
- Auroral processes in the **upward** current region
- Auroral processes in the **downward** current region
- Dayside and cusp aurora, theta aurora
- Temporal evolution in the auroral zone

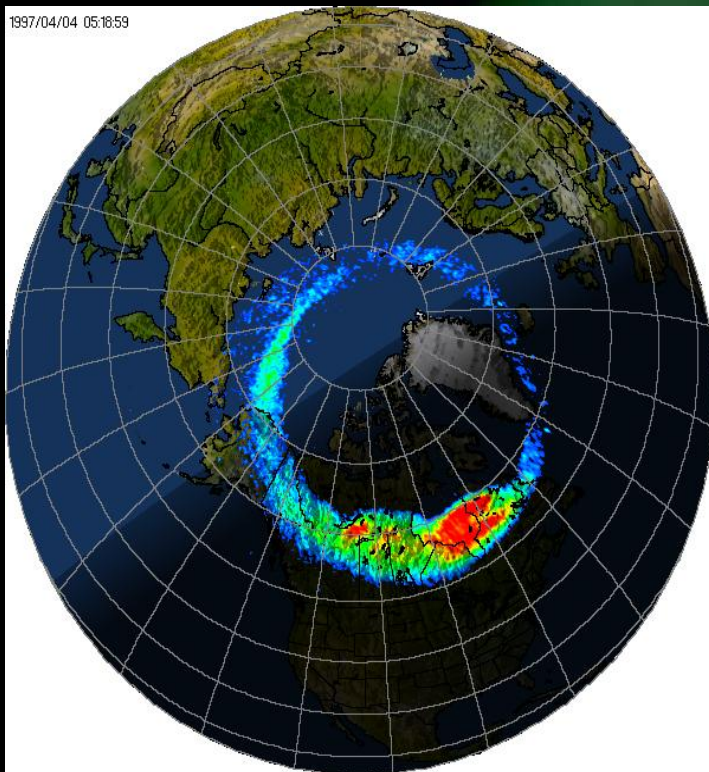
Very useful reference:

Space Science Series of ISSI
Auroral Plasma Physics

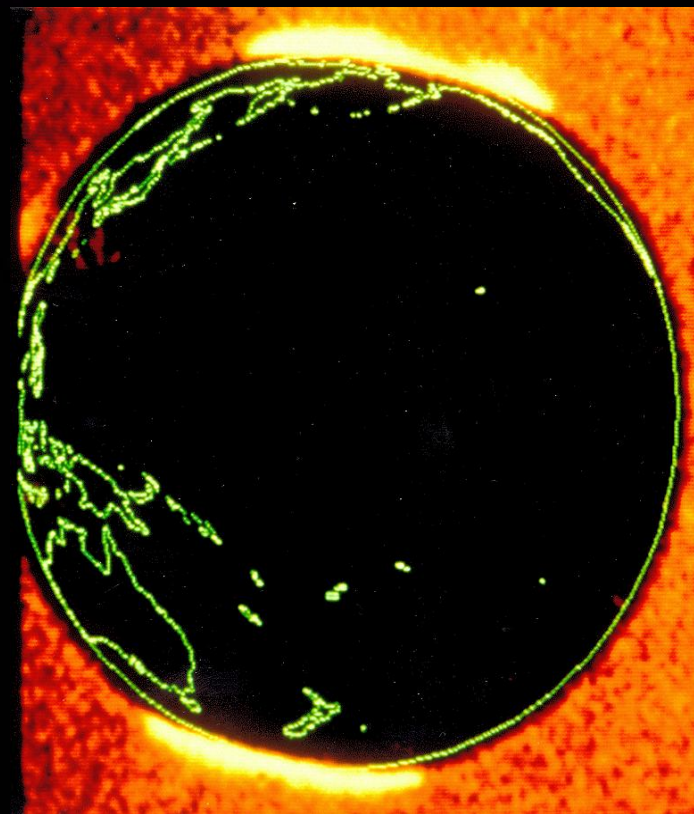
Götz Paschmann, Stein Haaland, Rudolf Treumann

(Space Science Reviews, vol 103, 1-4, 2002)

Auroral ovals

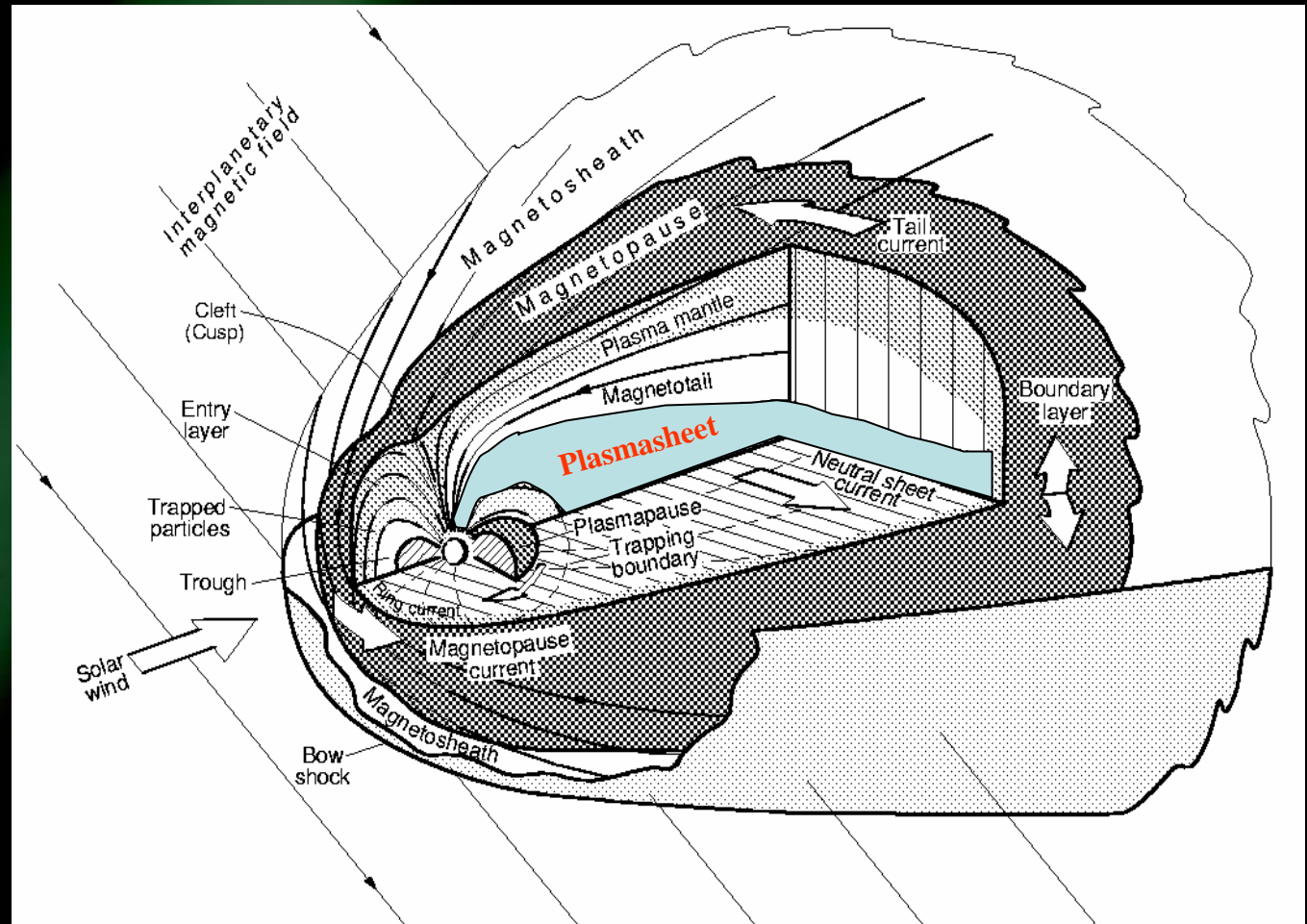


Polar

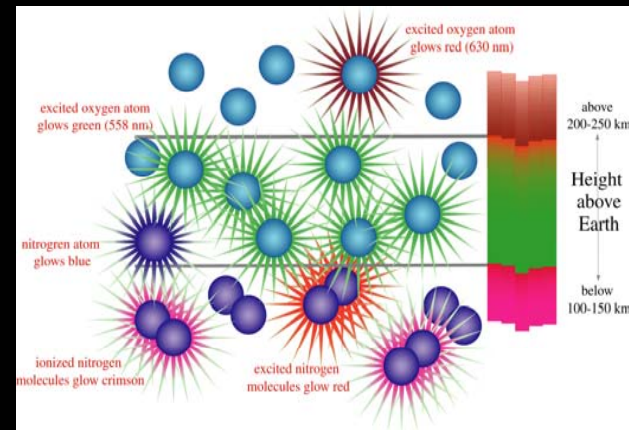
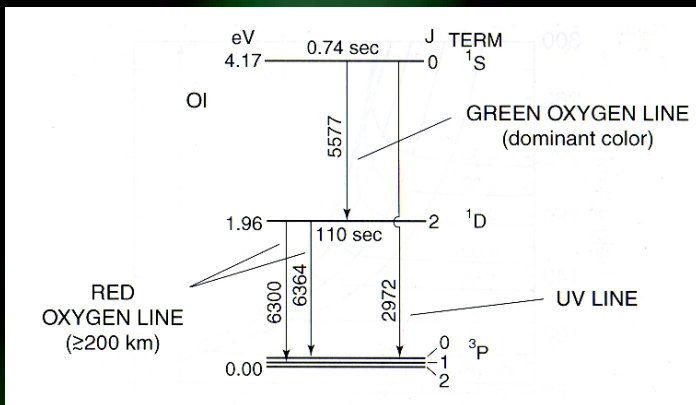
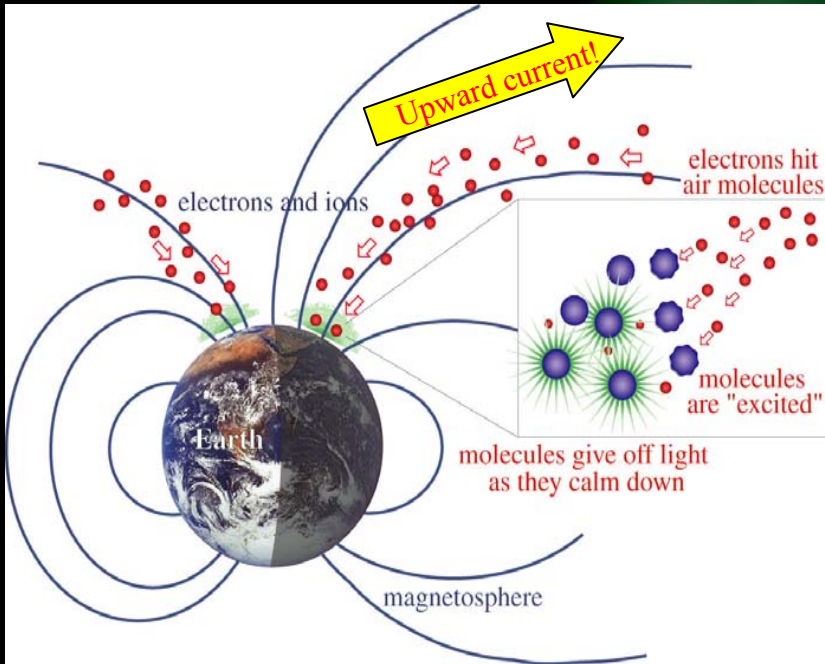


Dynamics Explorer

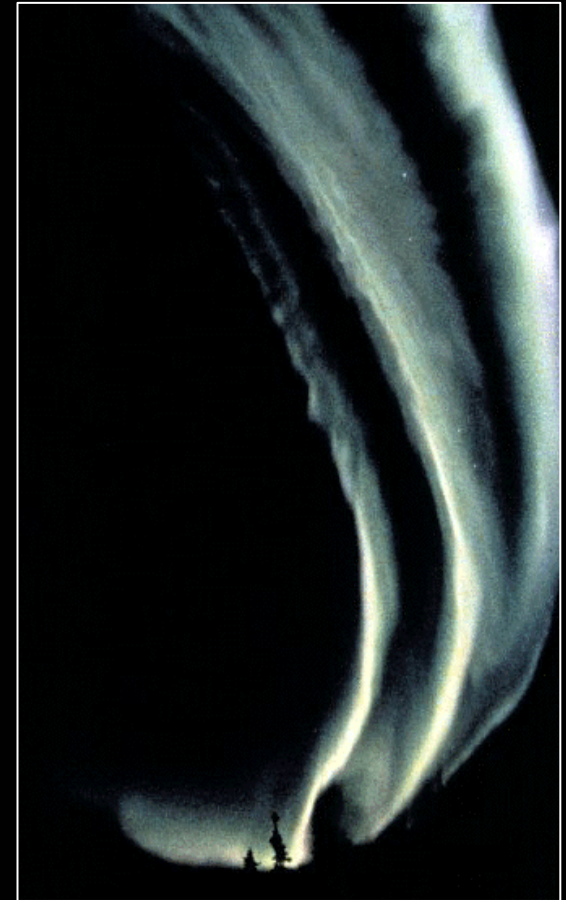
The auroral oval is the projection of the plasmasheet onto the atmosphere



Auroral emissions

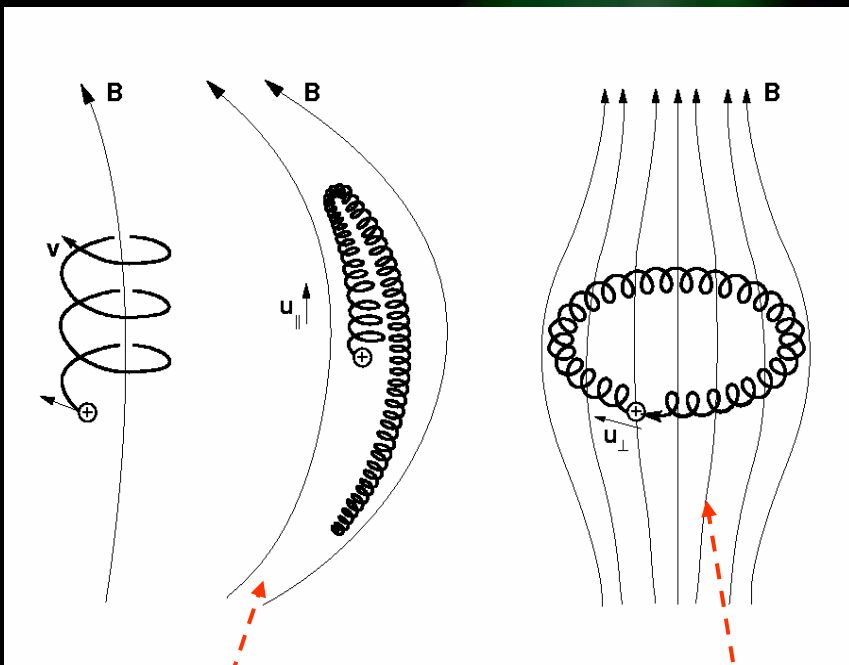


Homogenous auroral arcs



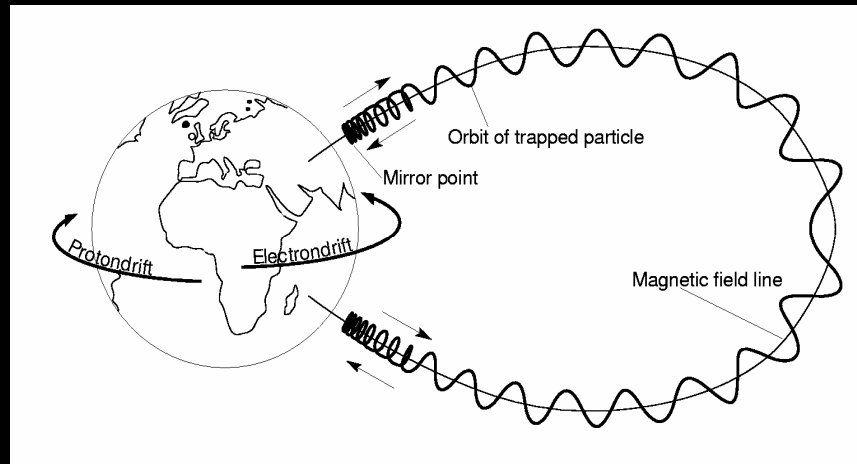
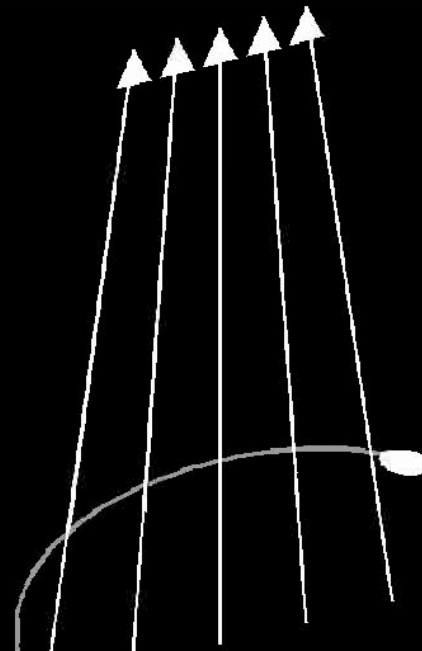
Particle motion in the geomagnetic field

gyration longitudinal oscillation azimuthal drift

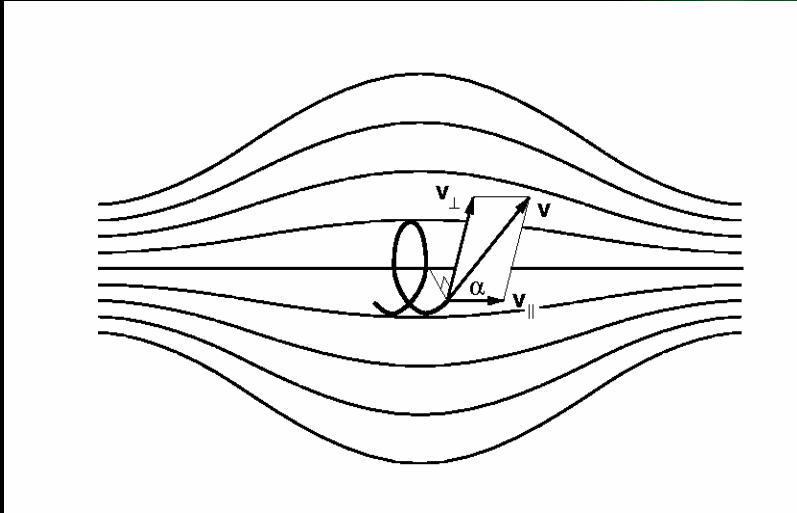


Magnetic mirror

grad B drift



Magnetic mirror



The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

$mv^2/2$ constant ➔

$$\frac{\sin^2 \alpha}{B} = \text{const}$$

particle turns when $\alpha = 90^\circ$ ➔

$$B_{\text{turn}} = B / \sin^2 \alpha$$

If maximal B-field is B_{max} a particle with pitch angle α can only be turned around if

$$B_{\text{turn}} = B / \sin^2 \alpha \leq B_{\text{max}} \quad \text{➔}$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B / B_{\text{max}}}$$

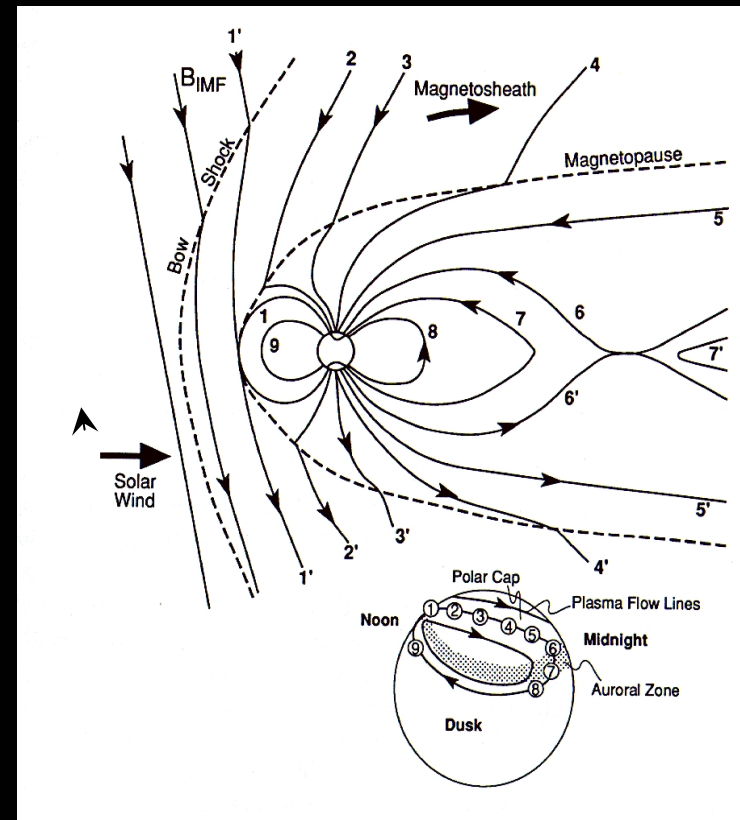
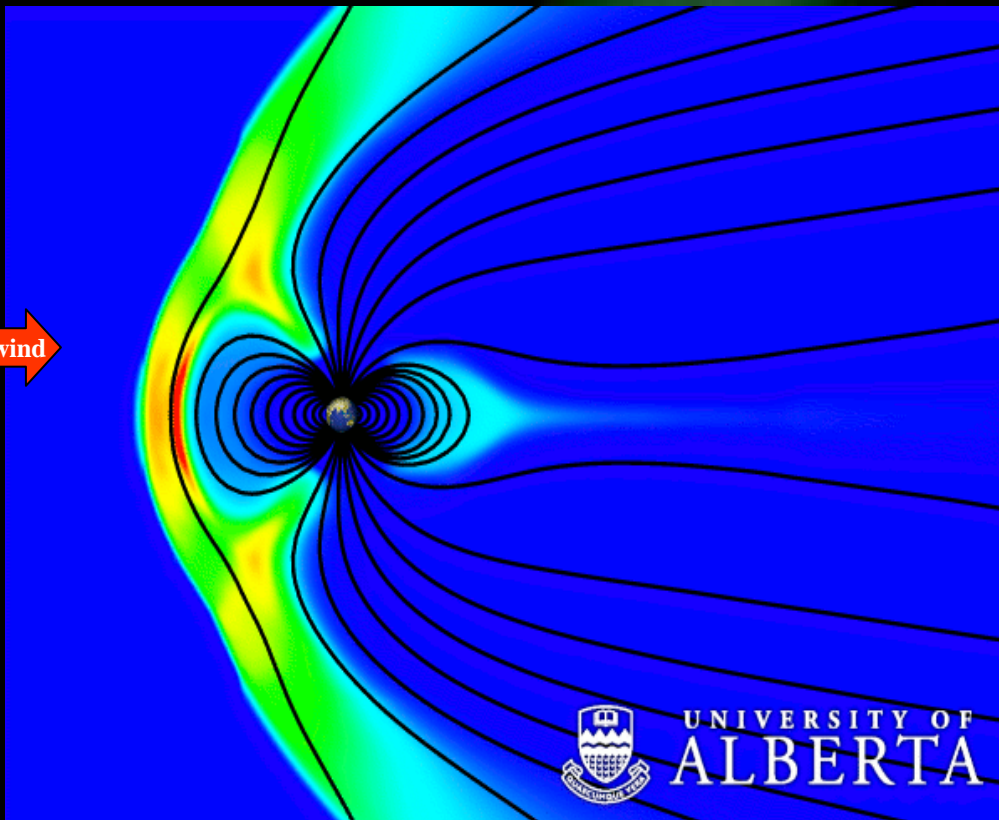
Particles in
loss cone :

$$\alpha < \alpha_{lc}$$

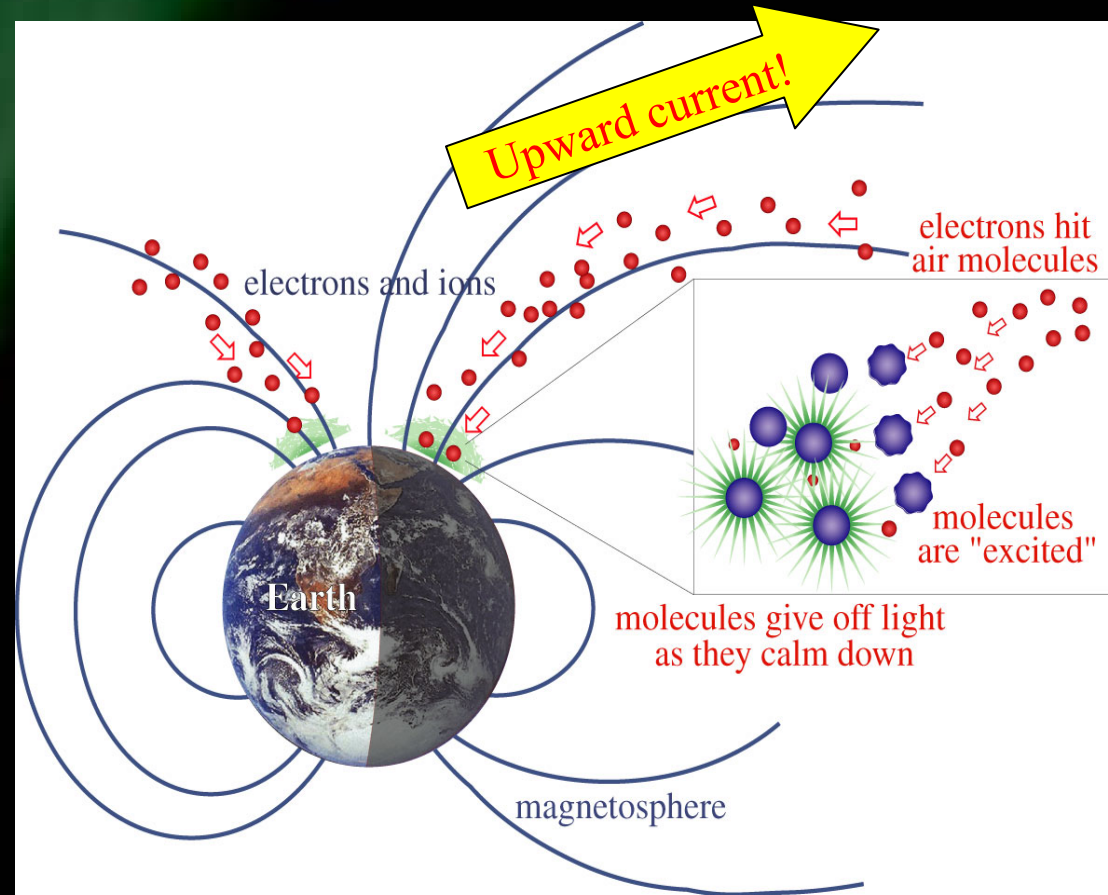
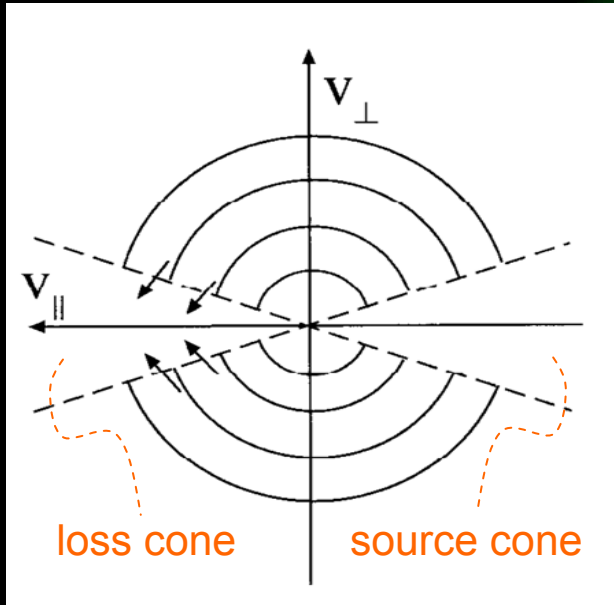
Magnetospheric convection



Solar wind

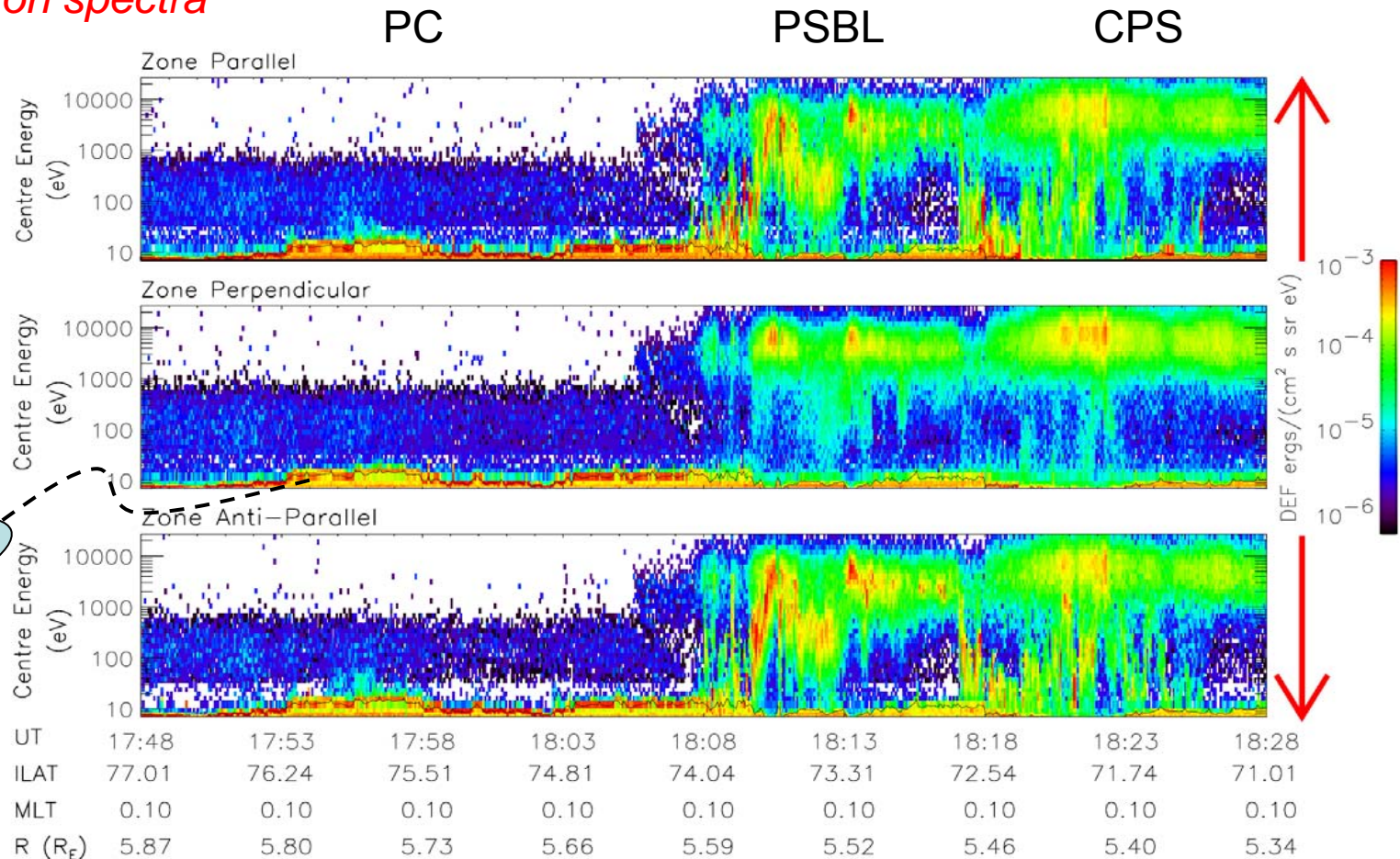


Atmospheric collisions - emissions



A typical auroral pass - CLUSTER

Electron spectra



Auroral scales

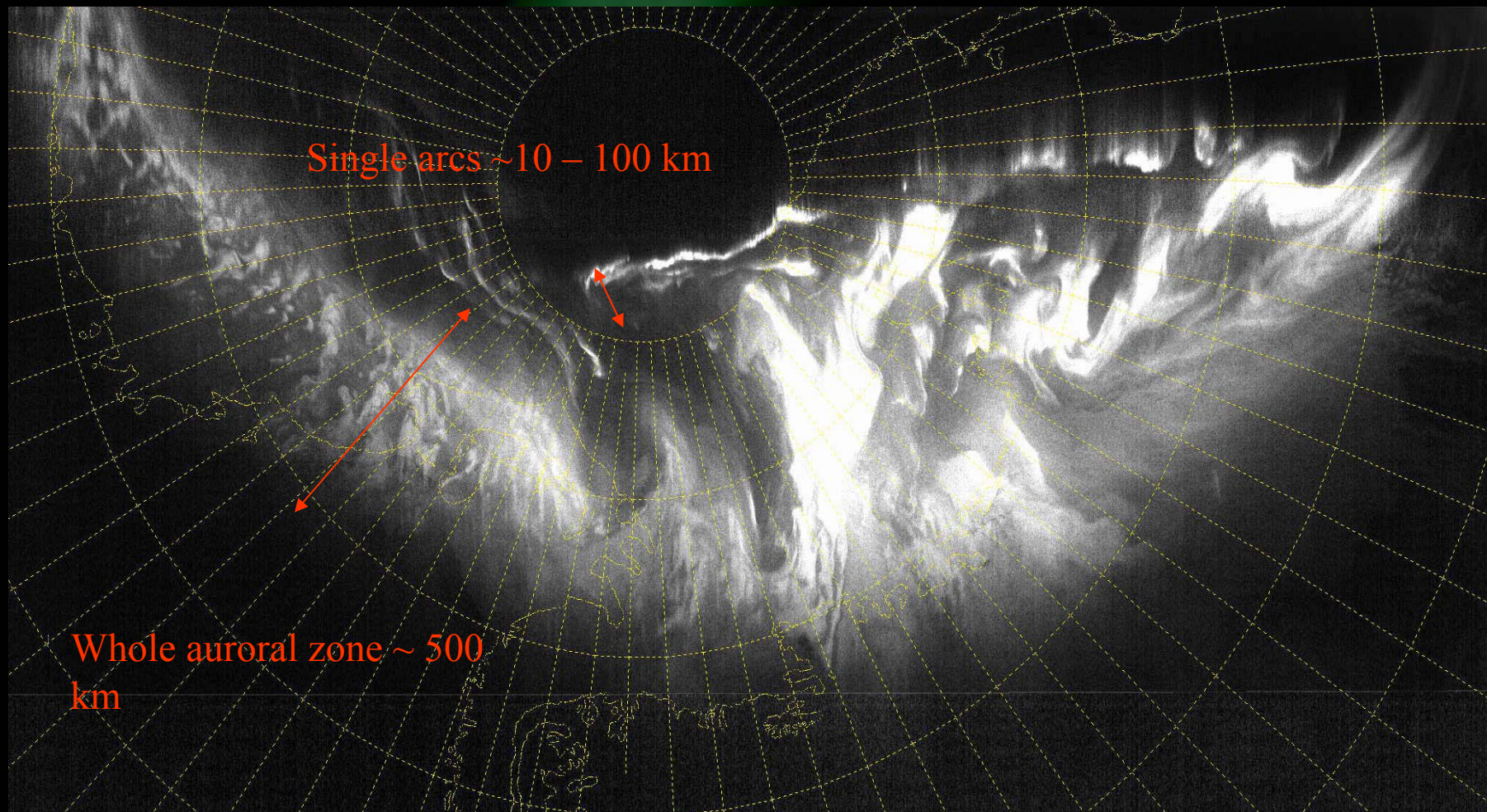
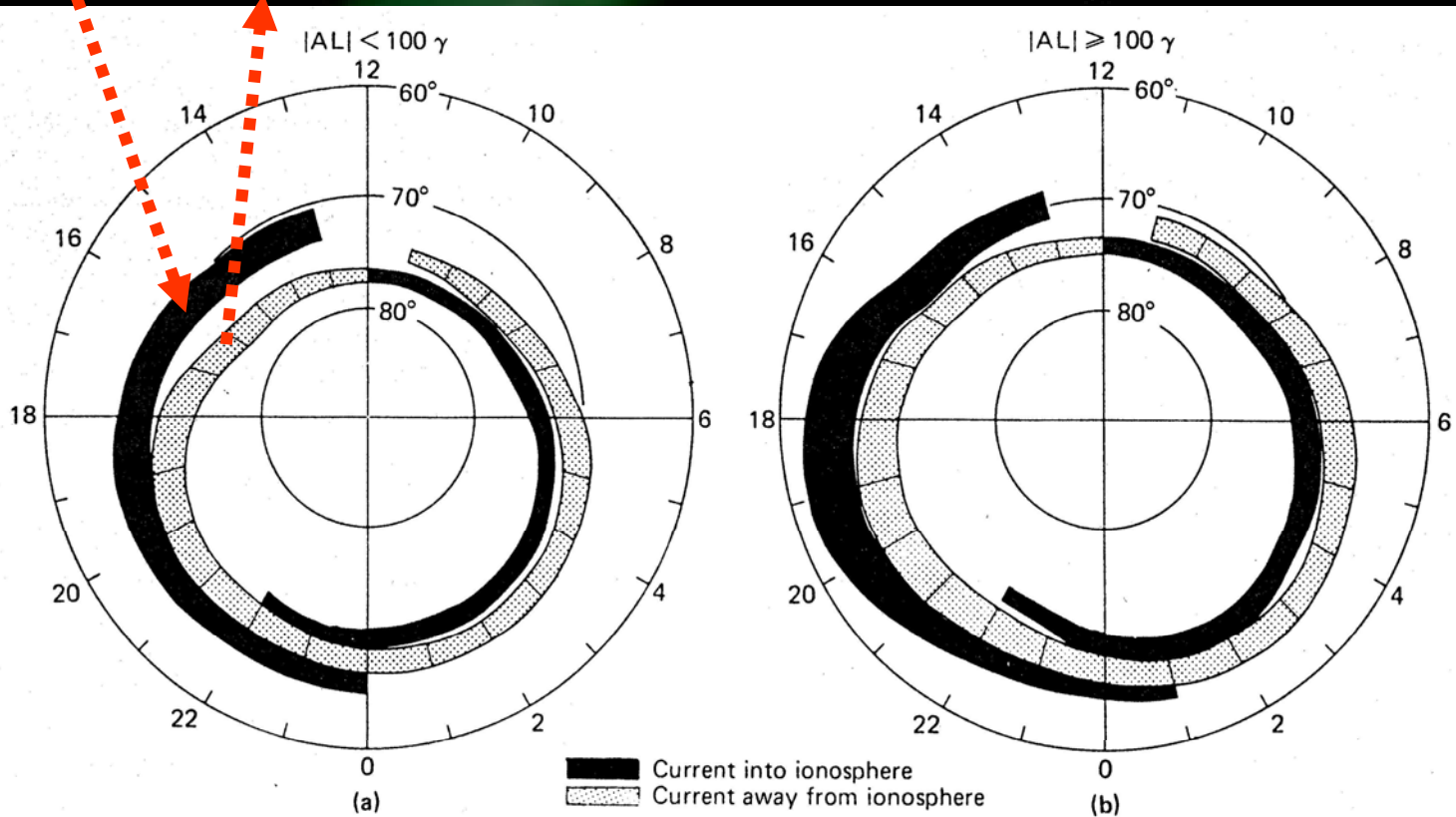


Photo from DMSP satellite

Birkeland currents in the auroral oval

Low geomagnetic activity

High geomagnetic activity



A typical auroral pass - FAST

ΔB_{EW}

↓ dE-flux, e^-

PA, e^-

↑ dE-flux, e^-

n-flux, e^-

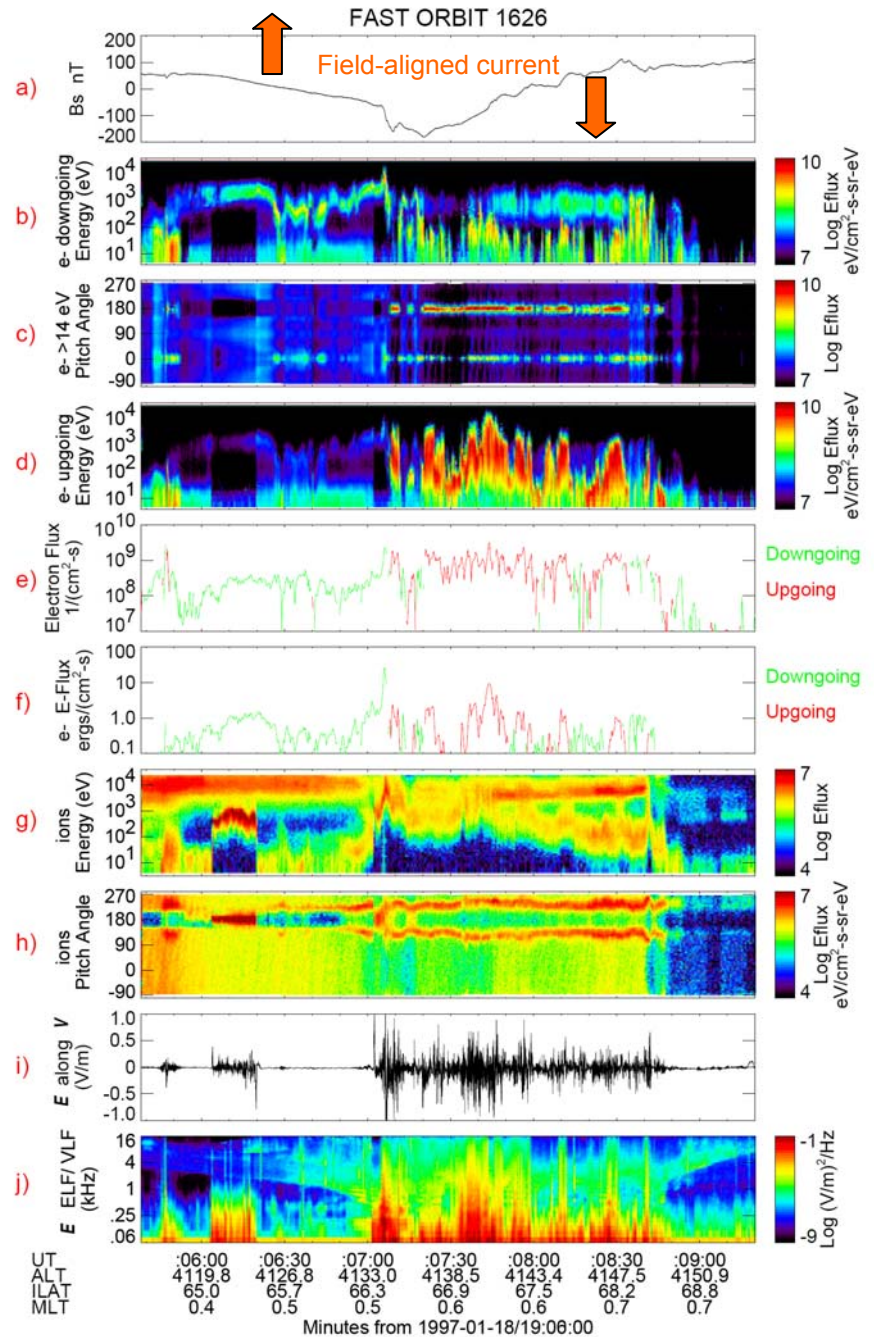
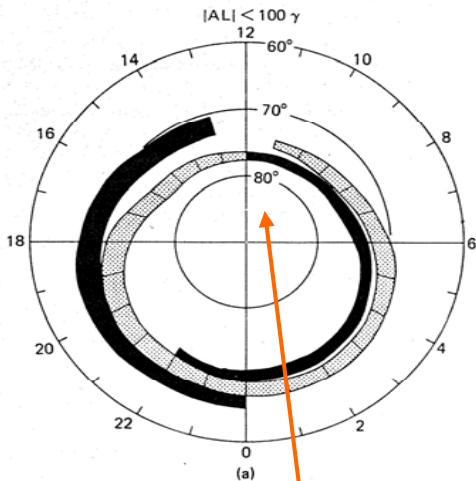
E-flux, e^-

dE-flux, i^+

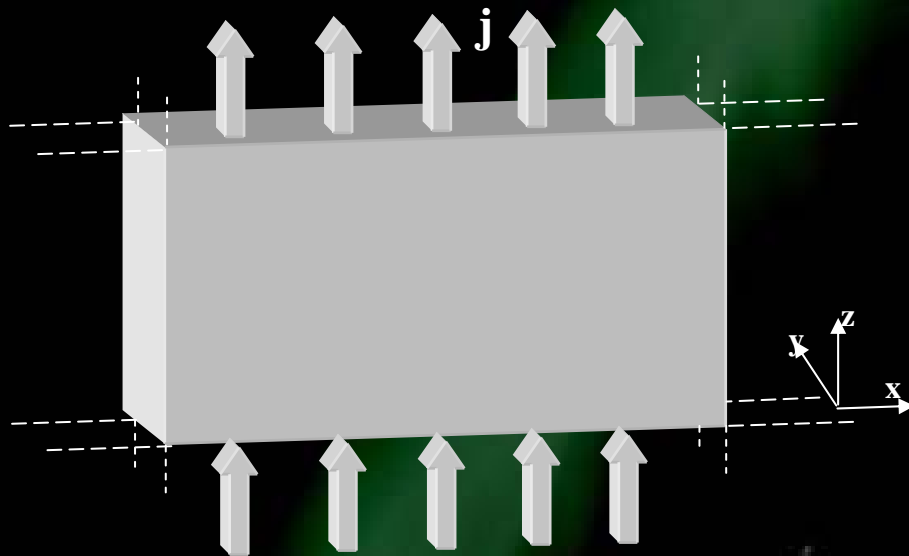
PA, i^+

E-perp (DC)

E (AC)



Current sheet approximation and Ampère's law



$$\left(\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z}, \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}, \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) = \mu_0 (j_x, j_y, j_z)$$

But $\frac{\partial}{\partial x} = 0$ and $\frac{\partial}{\partial z} = 0$

$$\left(\frac{\partial B_z}{\partial y}, 0, -\frac{\partial B_x}{\partial y} \right) = \mu_0 (0, 0, j_z)$$

Ampère's law (no time dependence):

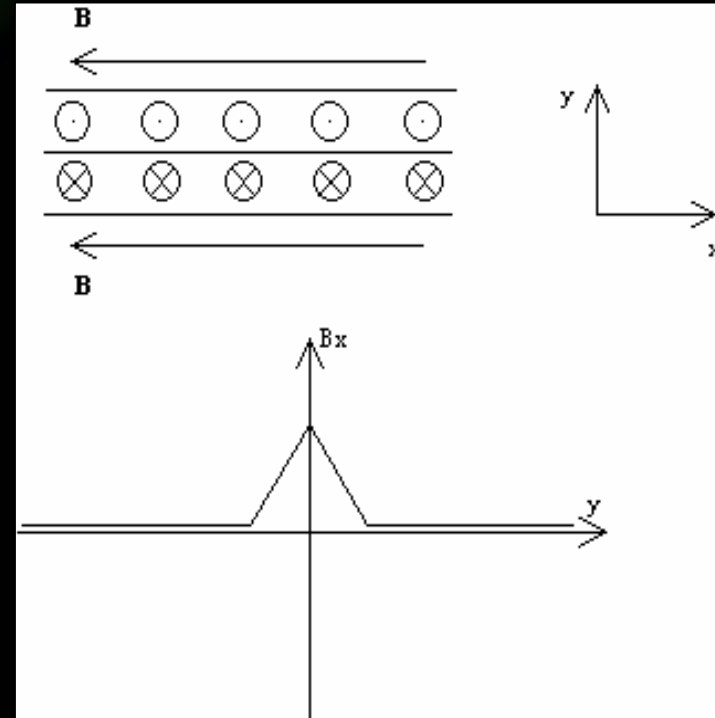
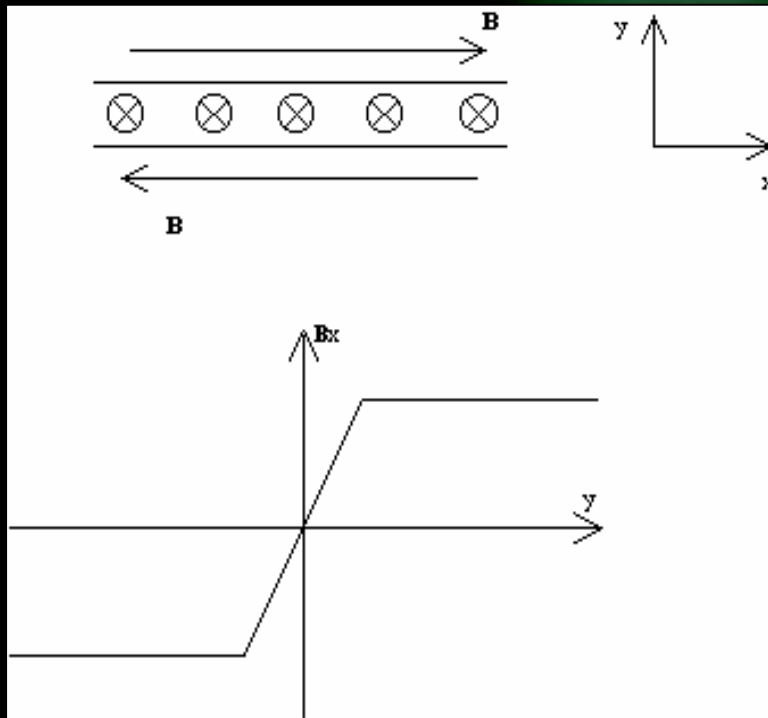
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

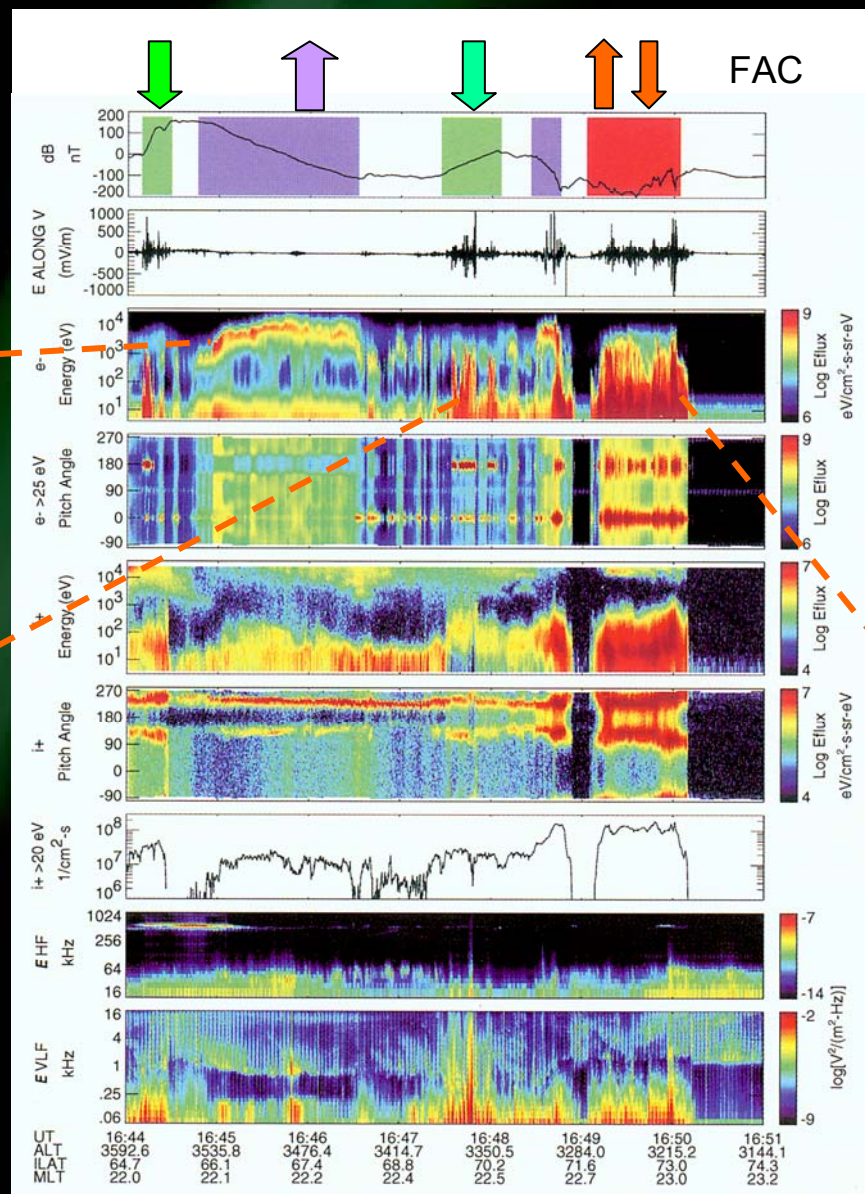
Current sheet

Determination of current density by magnetic field measurement



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

Upward and downward current regions



Inverted V

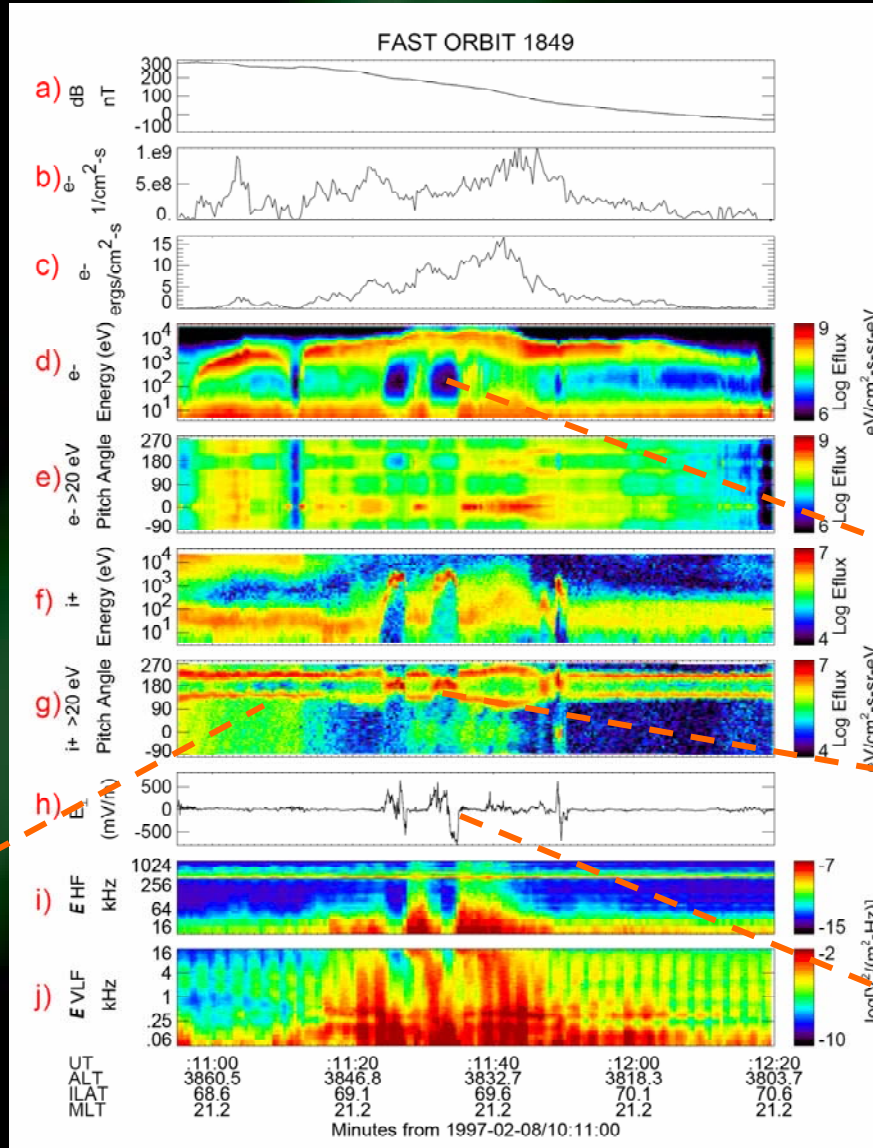
Upgoing electron beams

180° = ↑

0° = ↓

Alfvénic aurora

Upward current region Inverted V arc



180° = ↑

0° = ↓

Ion conic

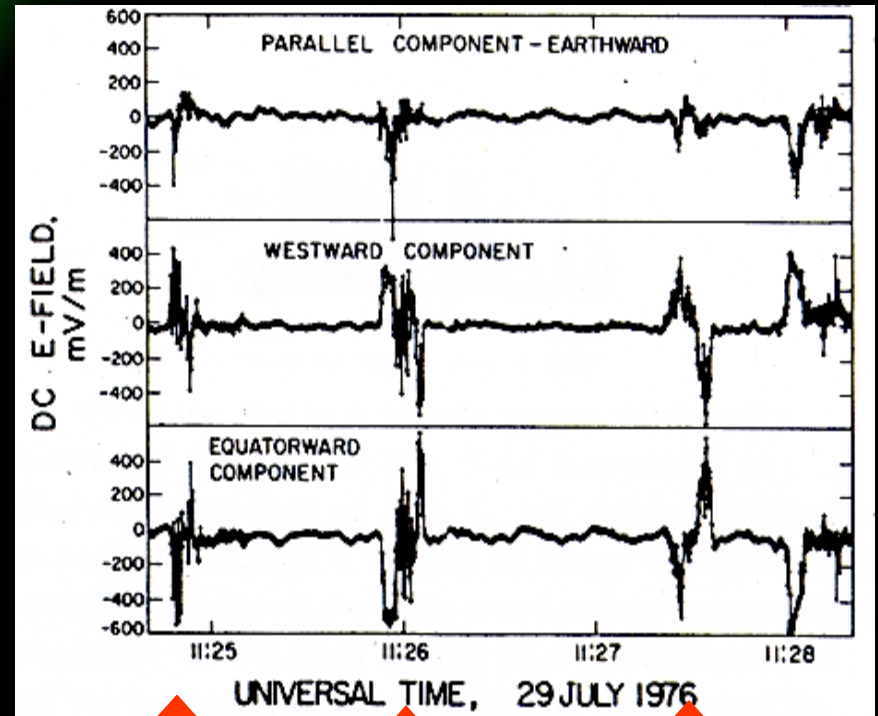
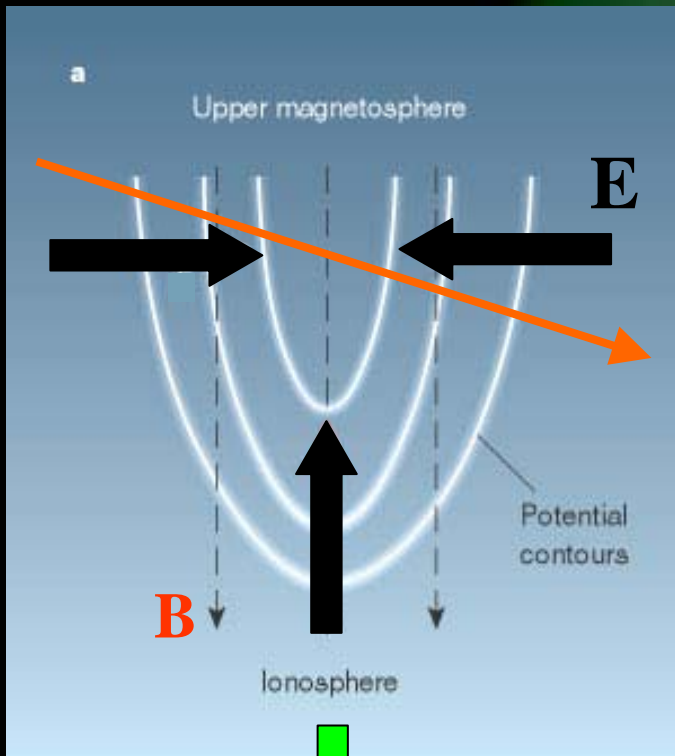
Visual limit $\sim 10^{-3} \text{ Jm}^{-2}\text{s}^{-1}$
 $= 1 \text{ erg cm}^{-2}\text{s}^{-1}$

Auroral density cavity

Ion beam

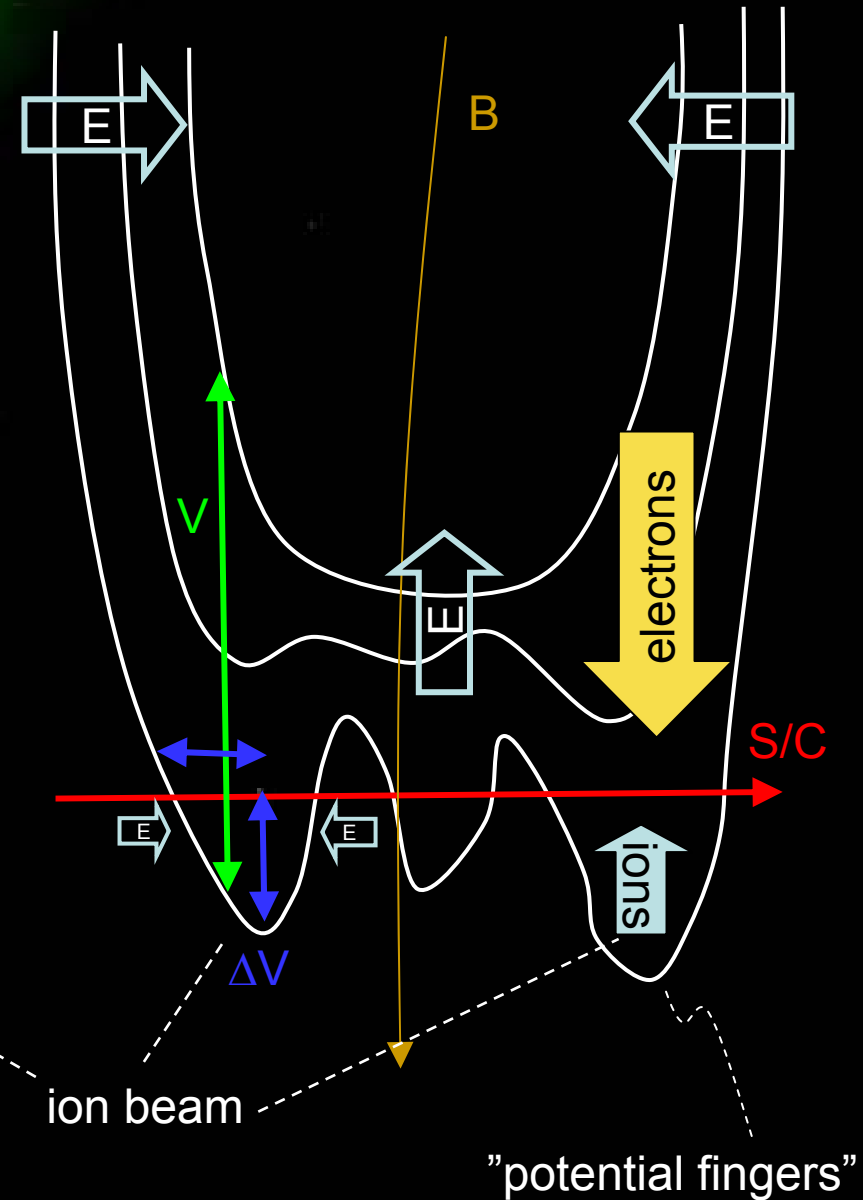
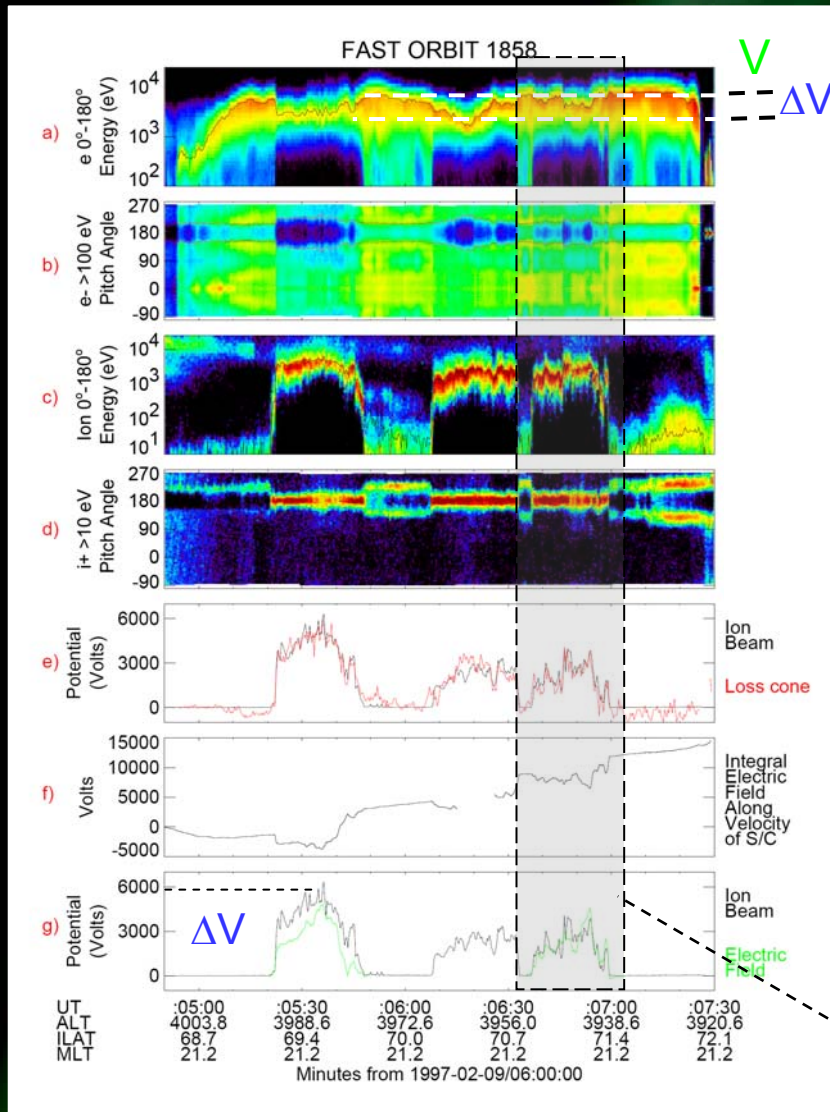
Bipolar electric field structures

Satellite signatures of U potential

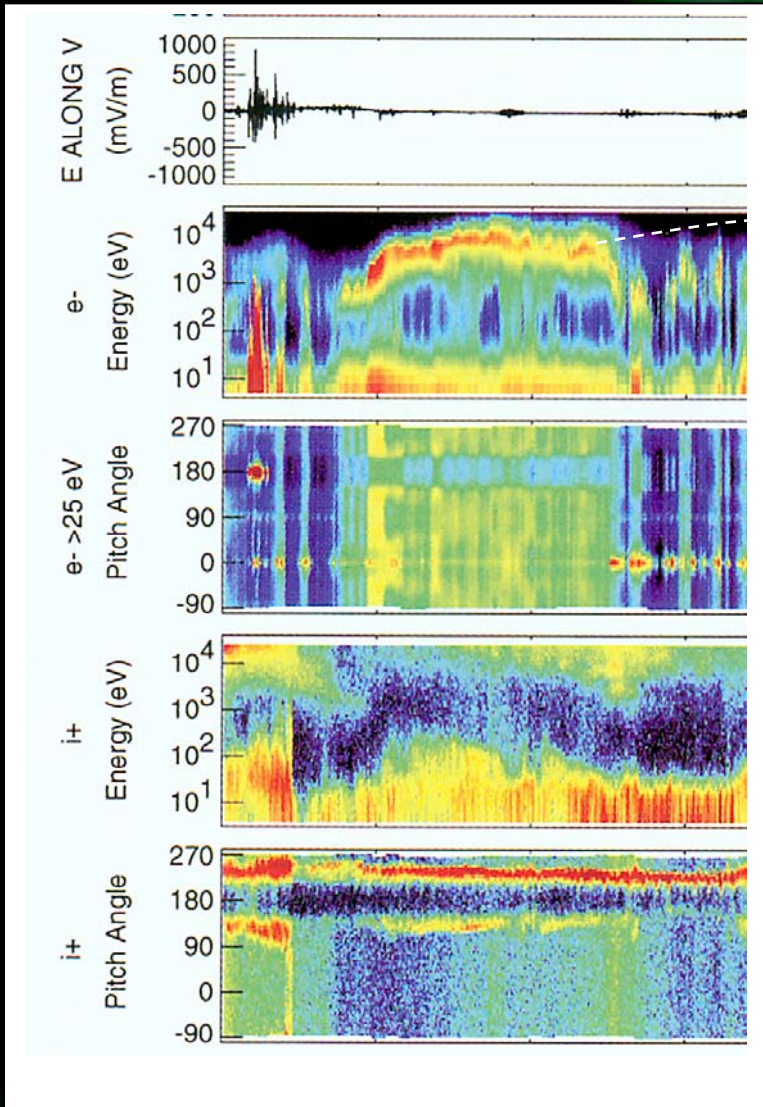


Measurements made by the ISEE satellite (Mozer et al., 1977)

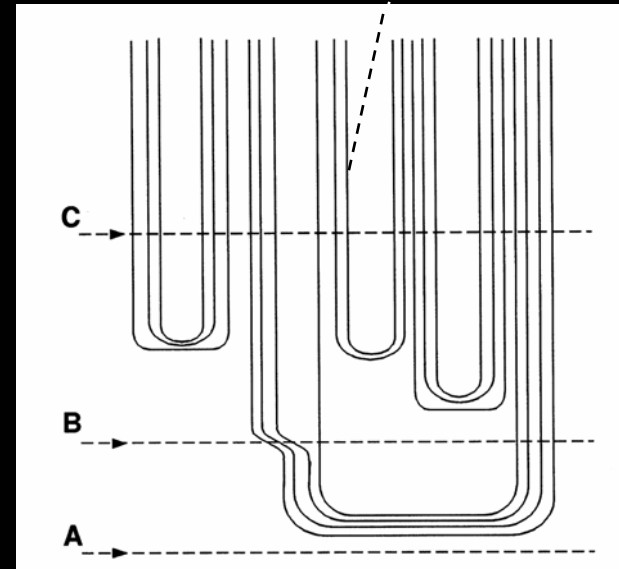
Acceleration potential structure I



Acceleration potential structure II

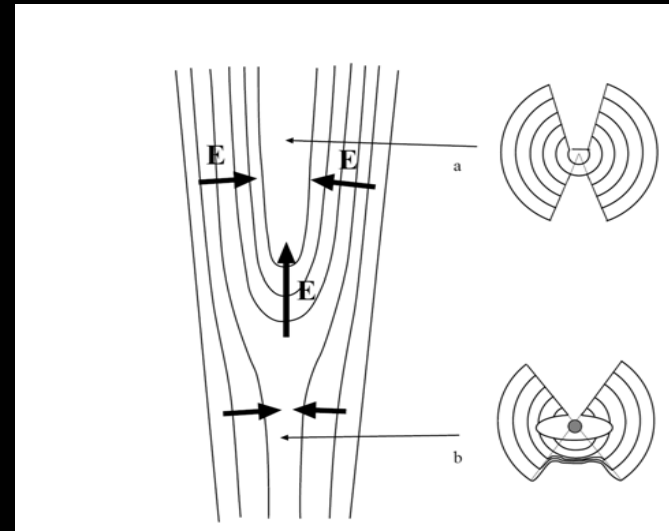
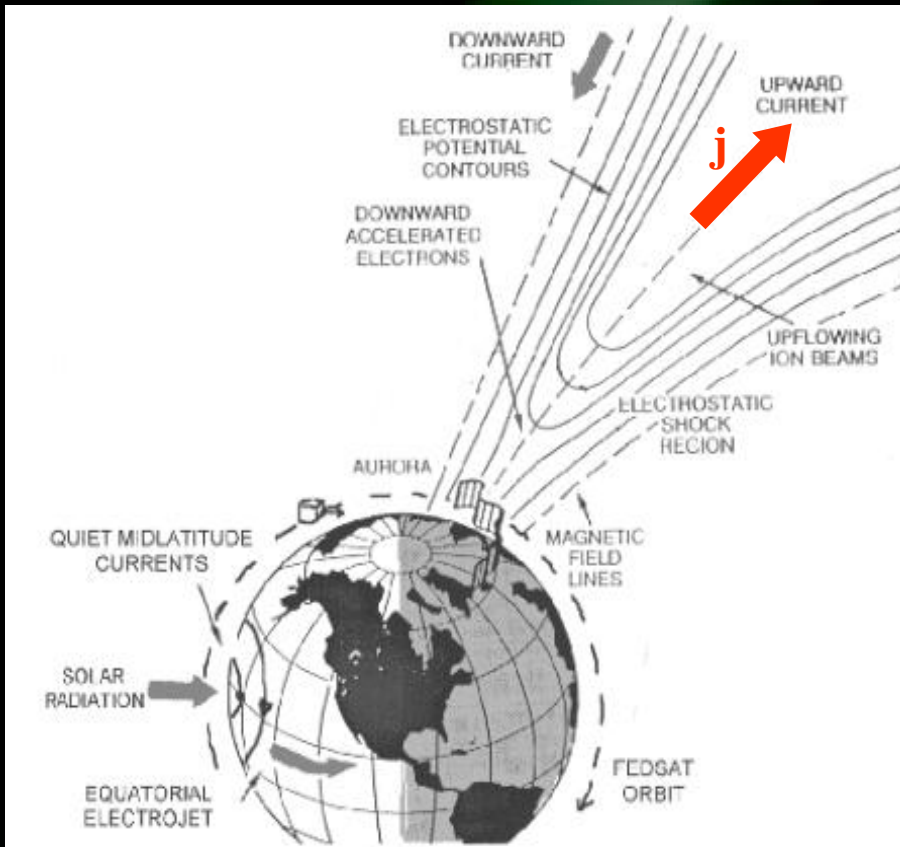


Localized regions of higher energy electrons without associated ion beams

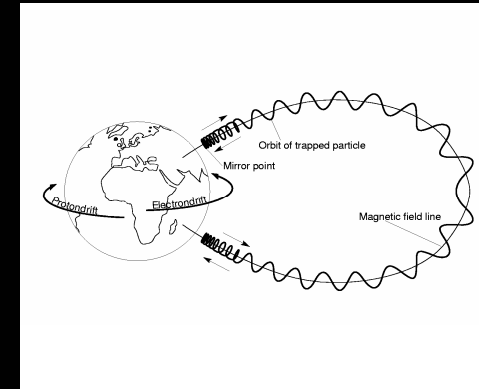
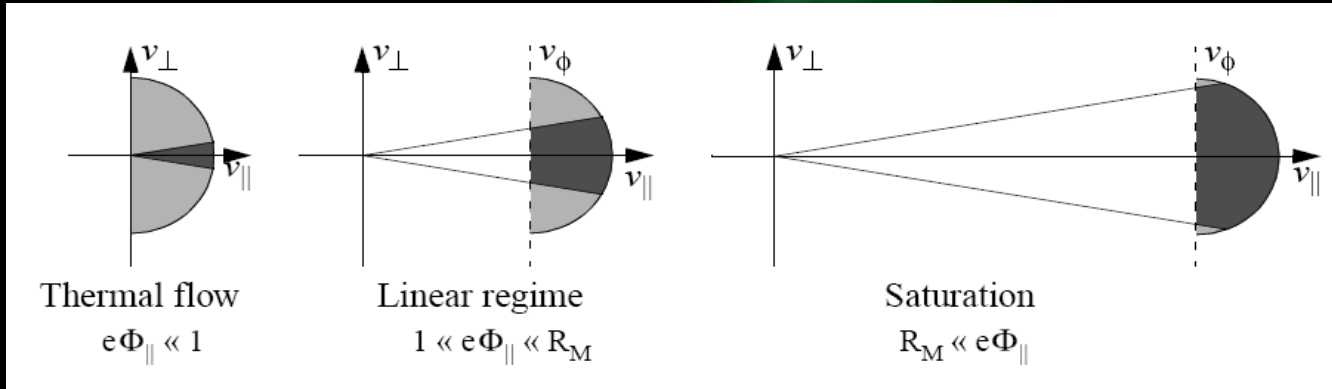


Why particle acceleration?

- The magnetosphere often acts as a current generator
- Electrons are accelerated downwards by upward E-field.
- This increases the pitch-angle of the electrons, and more electrons can reach the ionosphere, where the current can be closed.



Auroral currents – Knight relation



Fraction of particles in the loss cone:

$$f = \frac{\pi\theta_{lc}}{2\pi} \sim \frac{B_{ms}}{B_{ion}} \sim 2 \times 10^{-4}$$

Thermal current:

$$j_{\parallel,ms} = n_0 e v_{th} f$$

$$j_{\parallel,ion} = n_0 e v_{th} f \frac{B_{ms}}{B_{ion}} =$$

$$= n_0 e v_{th} \frac{B_{ms}}{B_{ion}} \frac{B_{ion}}{B_{ms}} = n_0 e v_{th} \approx$$

$$[n_e = 0.1 \text{ cm}^{-3}, T_e = 1 \text{ keV}] \approx$$

$$\sim 1 \mu\text{A/m}^2$$

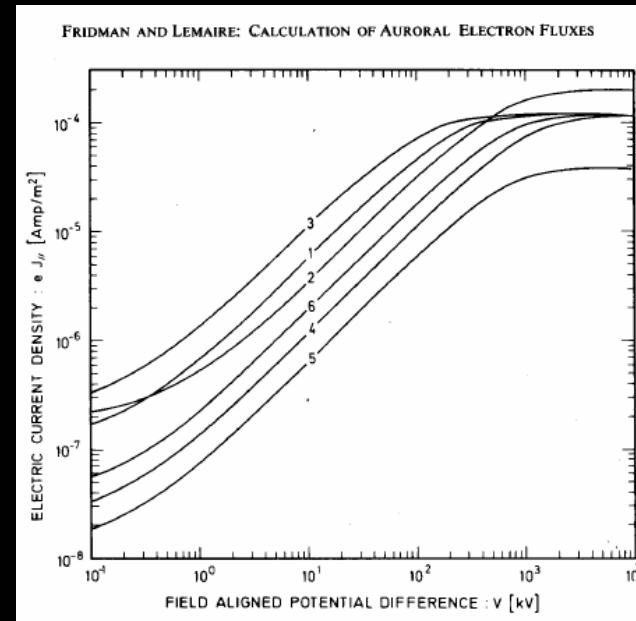
Apply a parallel potential drop:

$$j_{\parallel,ion} = n_0 e v_{th} \frac{B_{ion}}{B_{ms}} \left[1 - \frac{e^{-xe\Phi_{\parallel}/T_e}}{1+x} \right]$$

Linear regime :

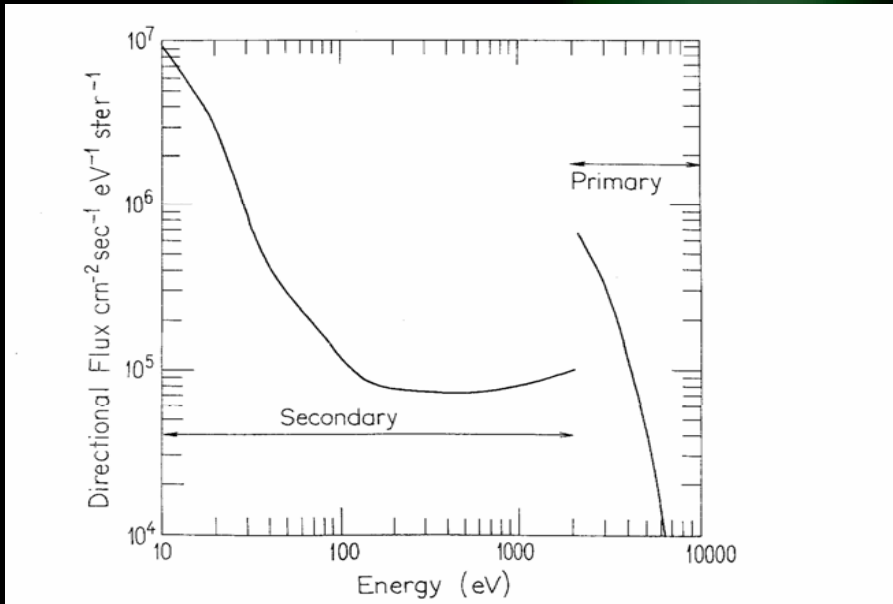
$$j_{\parallel,ion} \approx n_0 e v_{th} \frac{e\Phi_{\parallel}}{k_B T_e} = K \Phi_{\parallel}$$

$$K = \frac{e^2 n_0}{\sqrt{2\pi m_e k_B T_e}} \sim 10^{-9} \text{ S/m}^2$$

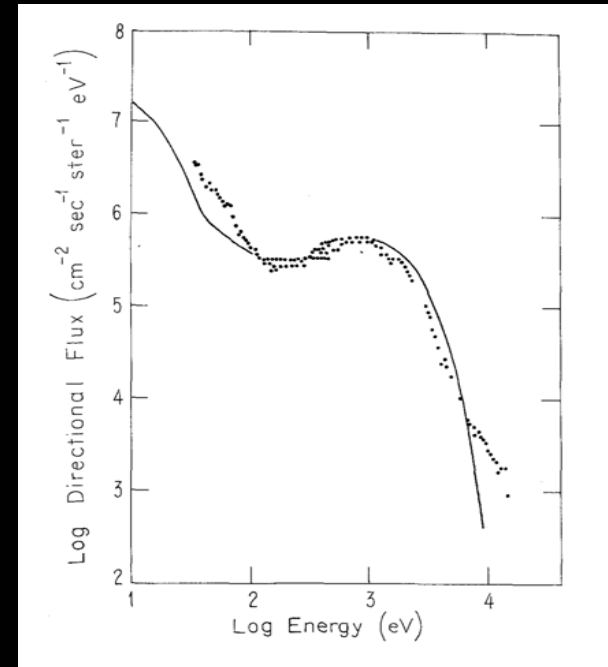


$$v_{th} = \sqrt{T_e / 2\pi m_e} \quad x = \frac{1}{B_i / B_0 - 1}$$

Particle distributions associated with inverted V's

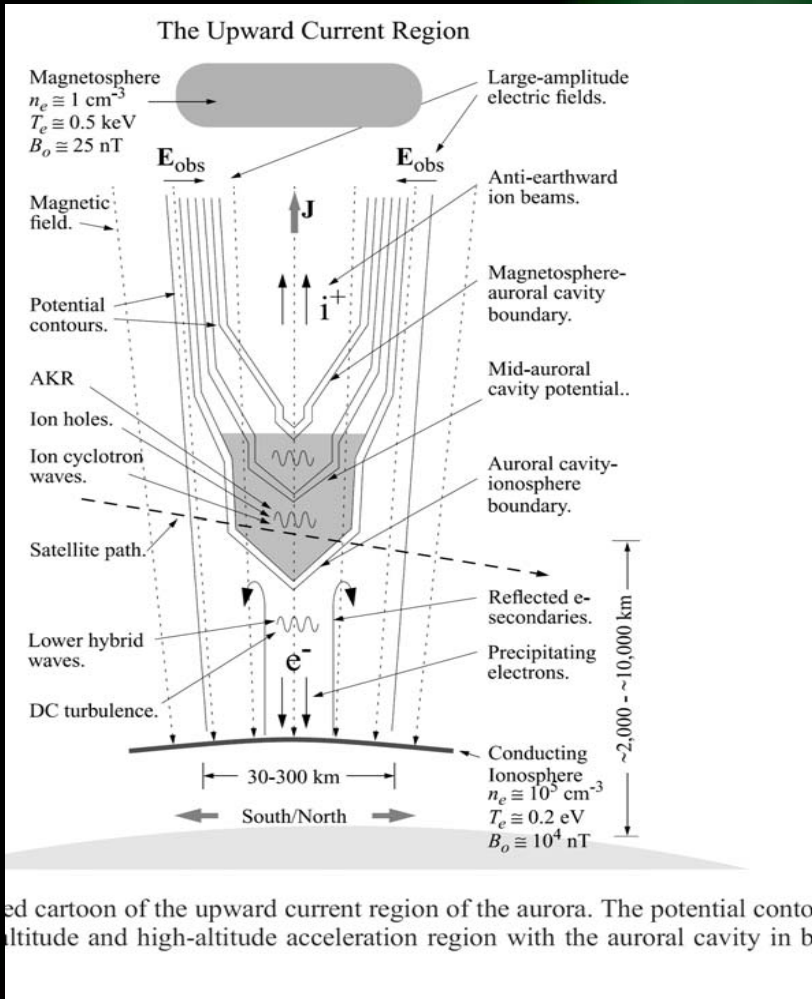


Model of cold beam producing secondaries (Evans, 1974)



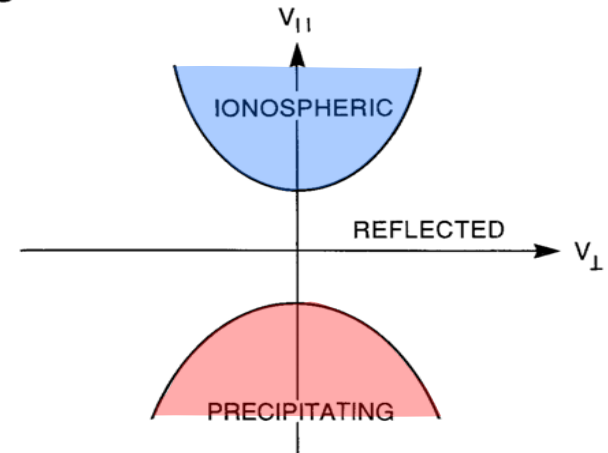
Model of hot electron beam and secondaries (Evans, 1974). Data from Franck and Ackerson, 1971

Auroral cavity and trapped populations

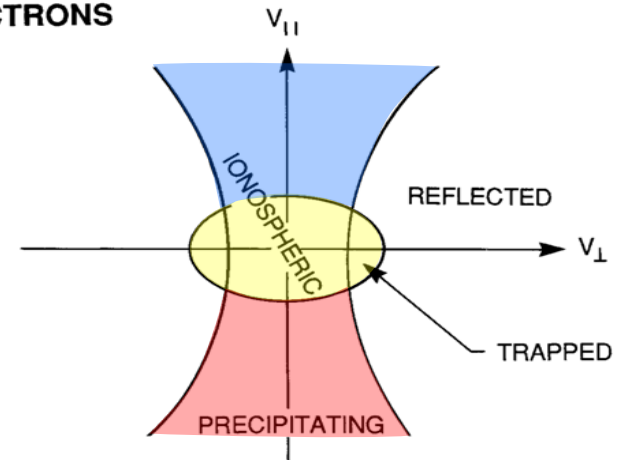


ed cartoon of the upward current region of the aurora. The potential conto
 altitude and high-altitude acceleration region with the auroral cavity in b

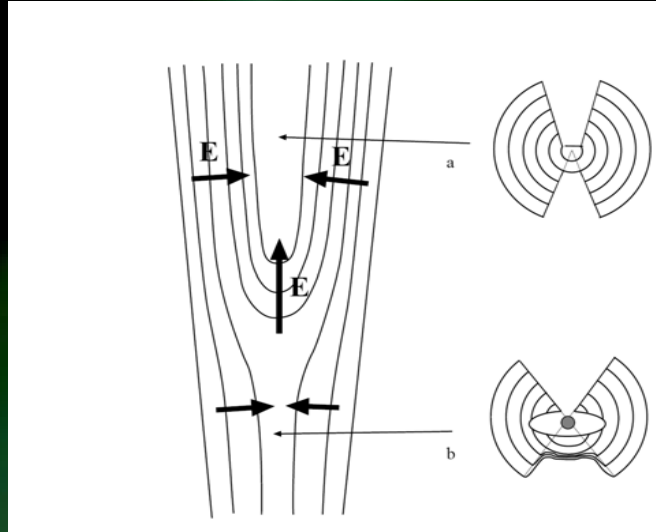
(a) IONS



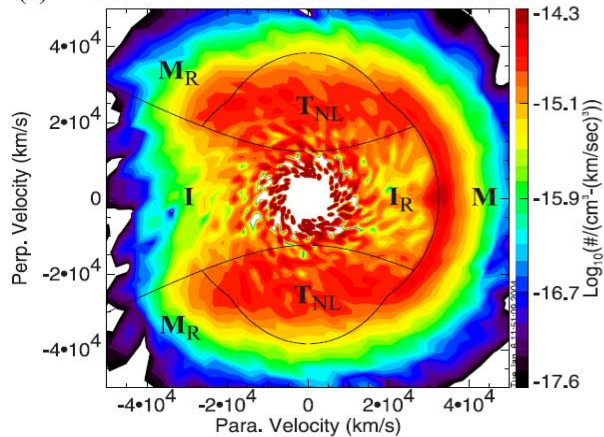
(b) ELECTRONS



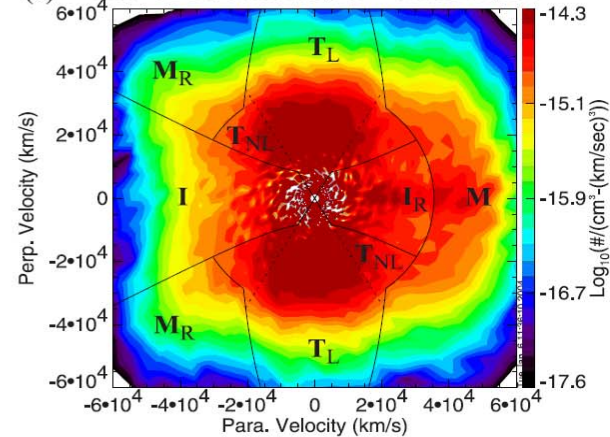
Accelerated Maxwellian



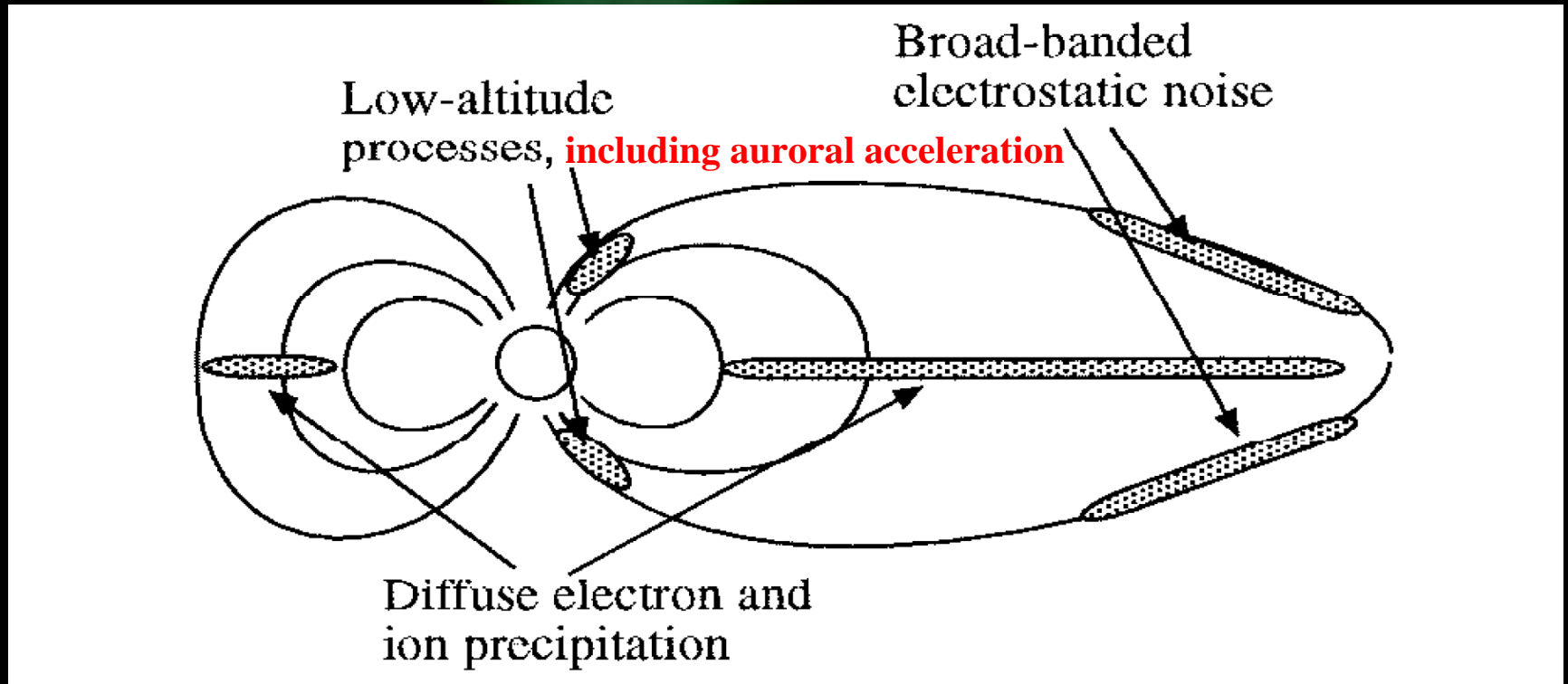
(a) FAST Eesa Survey DF
1997-02-17/05:00:09.126 - 05:00:09.443



(c) FAST Eesa Survey DF
1997-02-09/06:06:56.967 - 06:06:57.283



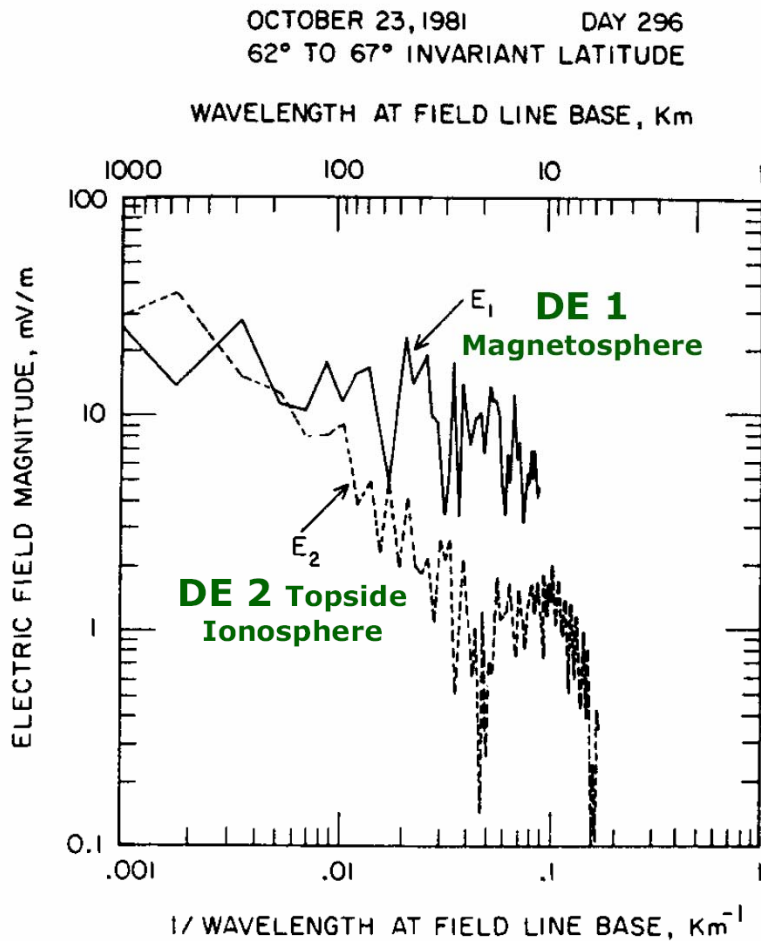
Acceleration regions



Koskinen

Auroral acceleration region typically situated at altitude of 1-3 R_E

Mapping of auroral electric fields



Large scales

Small scales

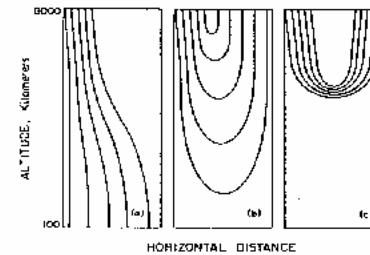


Figura 8. Contururi echipotentiale posibile in zona de accelerare (v. text). Din Hudson and Mozer (1978)

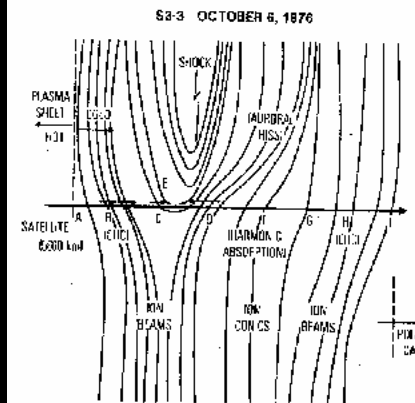
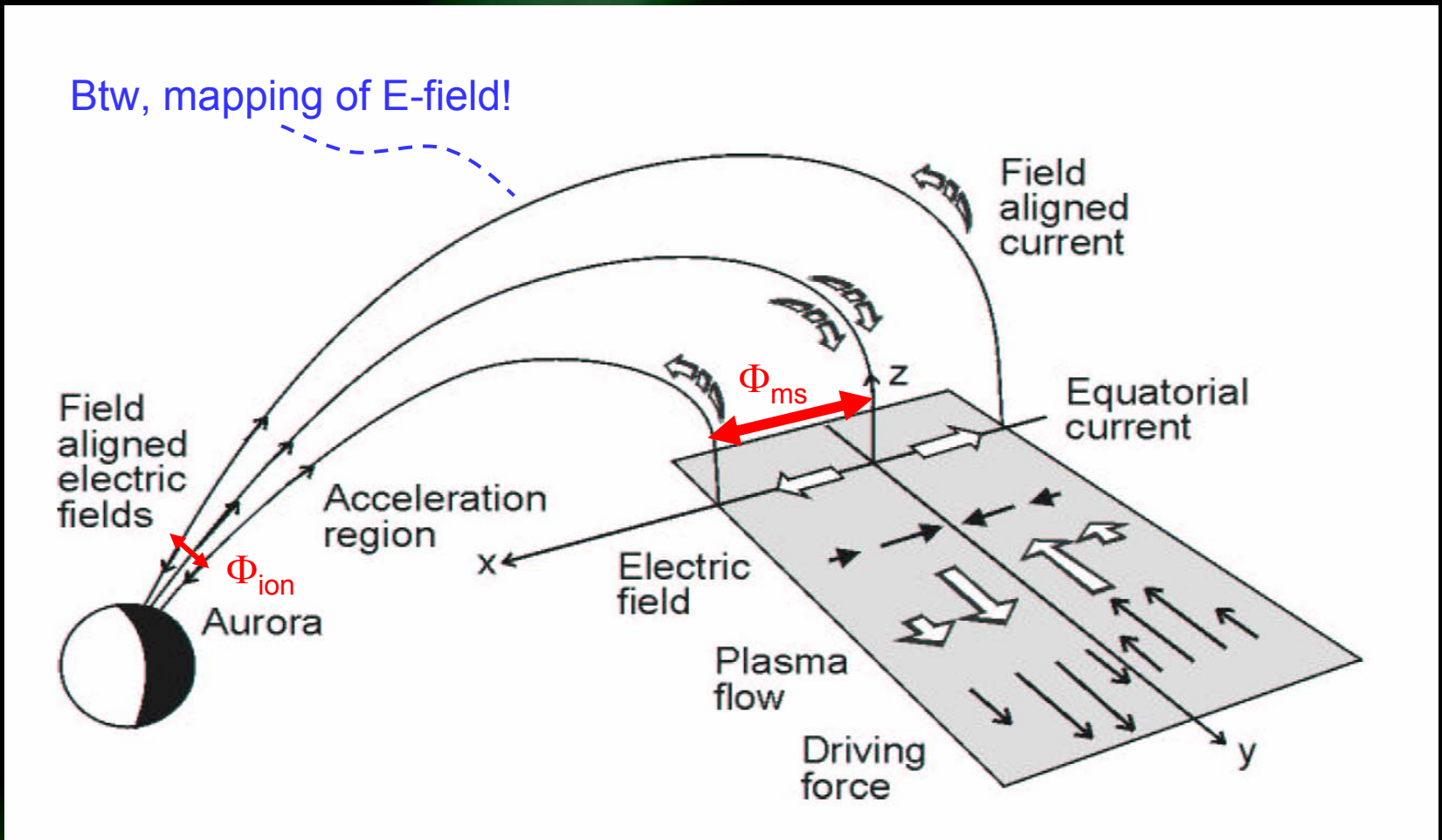


Figura 9. Reconstructia calitativa a contururilor echipotentiale din date 53-3. Figura ilustreaza 1 minut de observatii. Din Mizera et al. (1982)

Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (Weimer et al., 1985)

Static, medium-scale MI-coupling

MI-coupling critical scale size II



Static, medium-scale MI-coupling

MI-coupling critical scale size III

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp} + \mathbf{E}_{\perp} \cdot \nabla_{\perp} \Sigma_P + (\mathbf{E}_{\perp} \times \nabla_{\perp} \Sigma_H) \cdot \hat{\mathbf{b}}$$

$$K(\Phi_{ms} - \Phi_{ion}) \approx \Sigma_P \frac{\Phi_{ion}}{L^2}$$

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp}$$

$$\Phi_{ion} = \left(1 + \frac{\Sigma_P}{KL^2} \right)^{-1} \Phi_{ms}$$

$$j_{\parallel} = K \Delta \Phi_{\parallel} = K(\Phi_{ms} - \Phi_{ion})$$

$$\text{When } L \gg \sqrt{\frac{\Sigma_P}{K}} : \Phi_{ion} \approx \Phi_{ms}$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp ion}$$

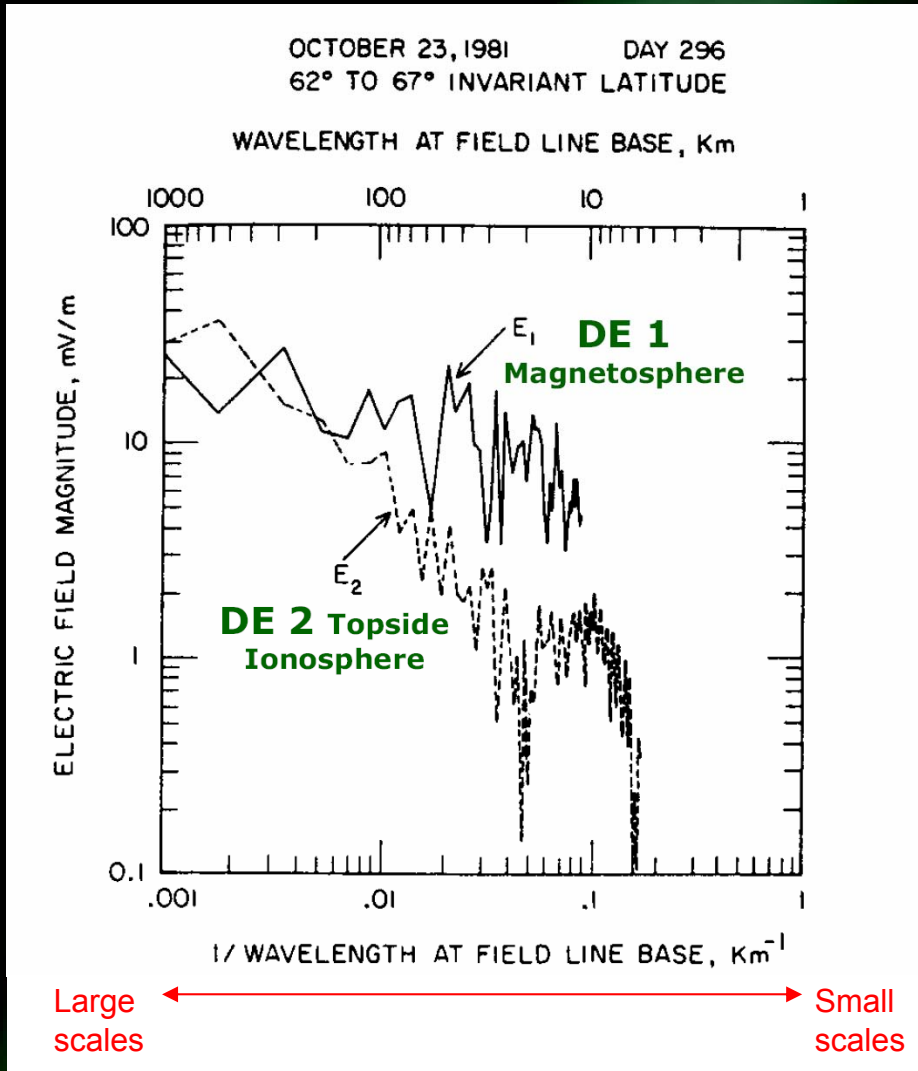
$$\text{When } L \ll \sqrt{\frac{\Sigma_P}{K}} : \Phi_{ion} \ll \Phi_{ms}$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp}^2 \Phi_{ion}$$

The last case means $|\Phi_{ms}| \approx |\Delta \Phi_{\parallel}|$

Static, medium-scale MI-coupling

MI-coupling critical scale size



Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (*Weimer et al., 1985*)

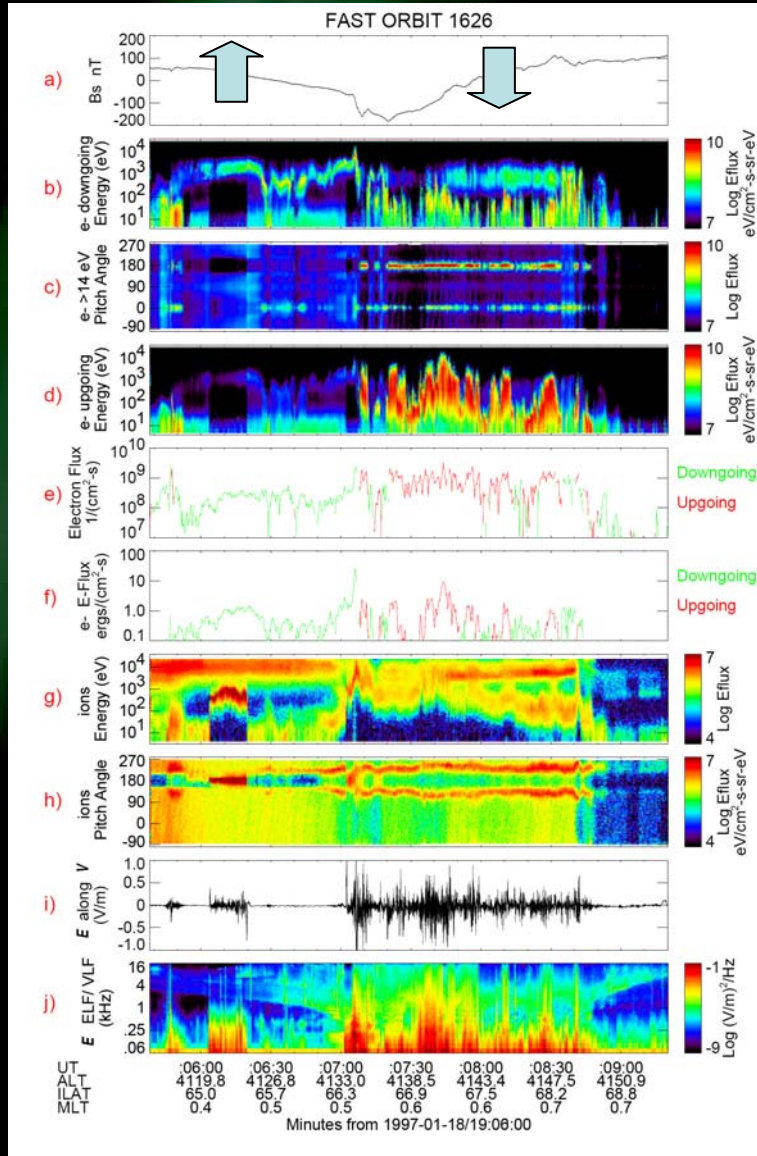
1/WAVELENGTH AT FIELD LINE BASE, Km⁻¹

Fig. 4. Electric field spectrums from day 296 (October 23) of 1981. The spectrums are obtained from a Fourier transform of the electric field data between 62° and 67° invariant latitude. The solid line shows the spectrum of the electric field measured by DE 1. The solid line shows the spectrum of the electric field measured by DE 2. The ordinate values are obtained from the square root of the "spectral power density." The actual units are mV m⁻¹ km^{1/2}. *Weimer et al, 1985*

Downward current region

Upward current region

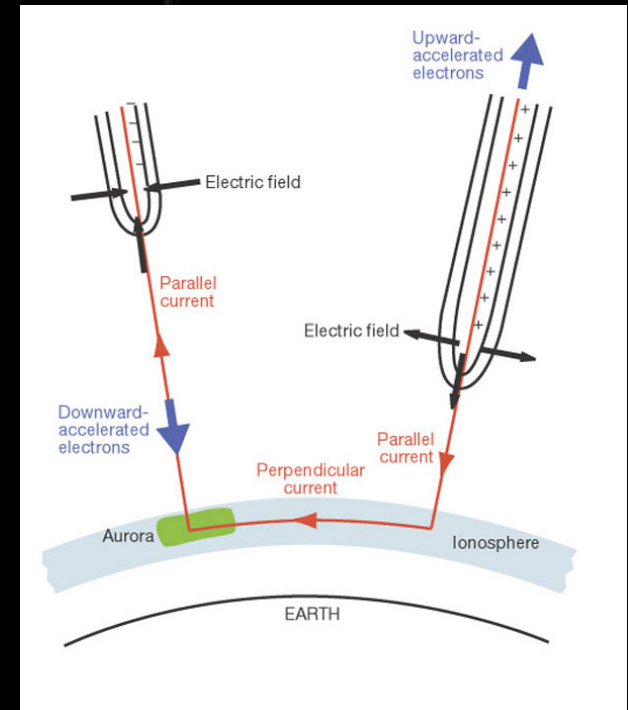
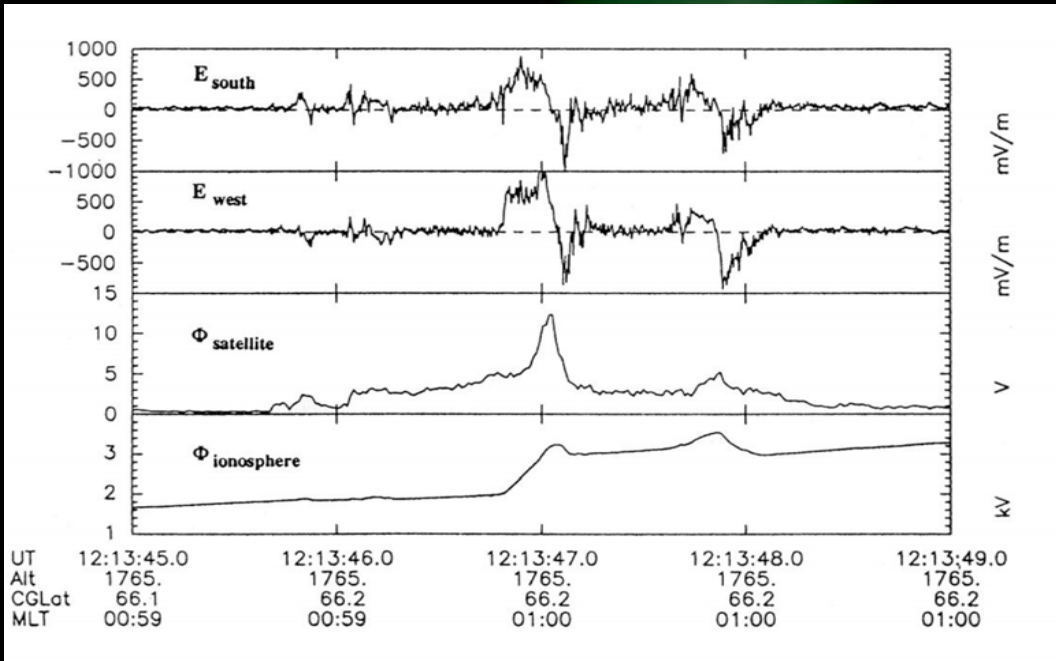
Downward electron beams:
Narrow in energy -
broad in pitch-angle



Downward current region

Upward electron beams:
Narrow in pitch angle - broad in energy

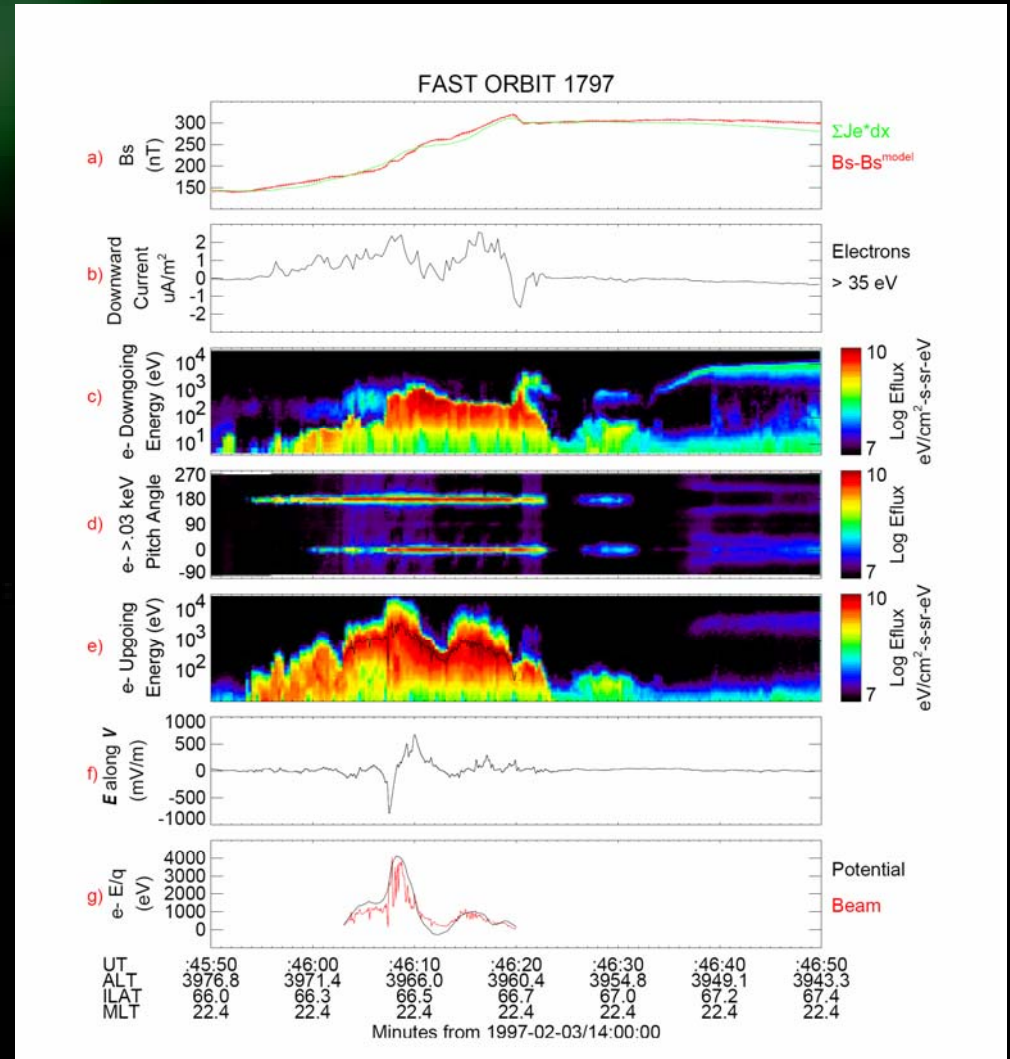
Potential structure in the downward current region



Freja electric field measurements,
(Marklund et al., 1994)

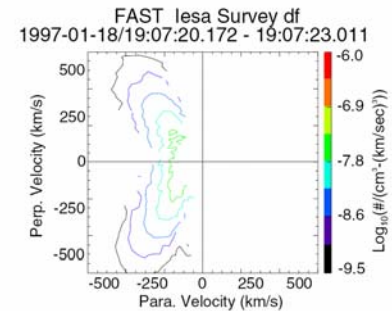
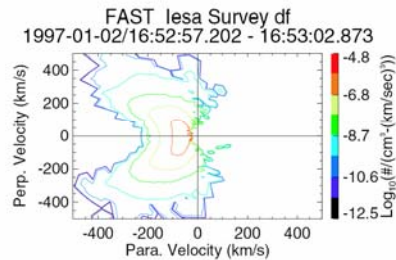
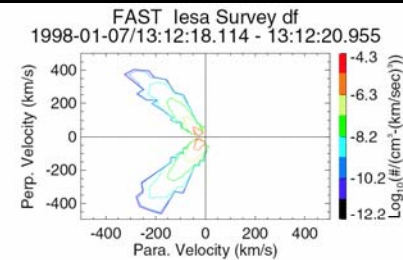
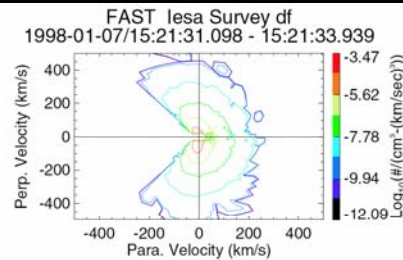
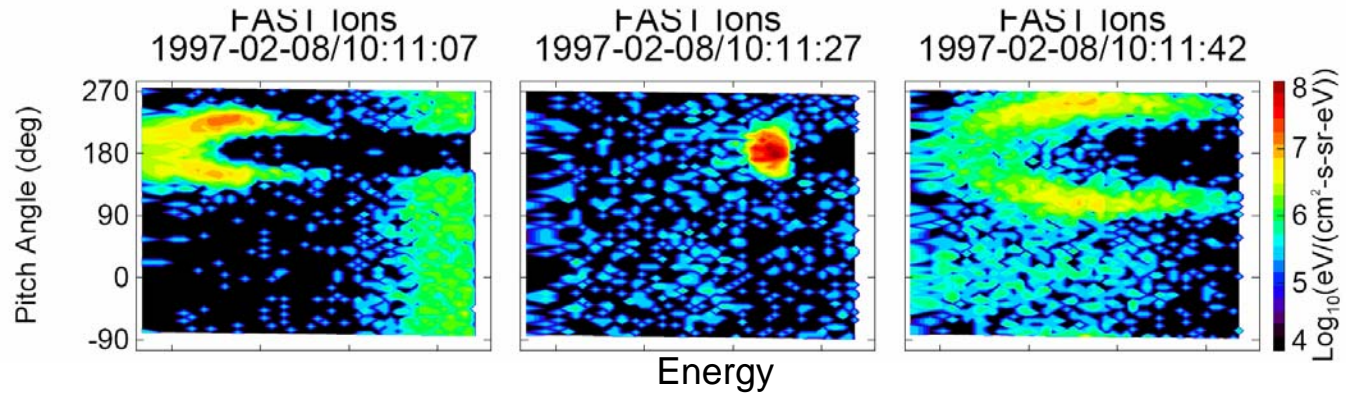
Upward electron beams

- Good agreement with integrated E-field
- Widening in energy is due to extensive wave-particle interaction.



Ion conics and beams

'Distribution functions'

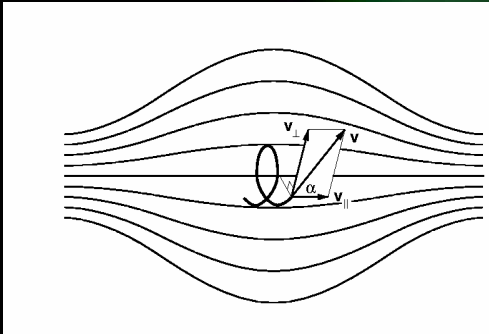


Ion conics – adiabatic motion

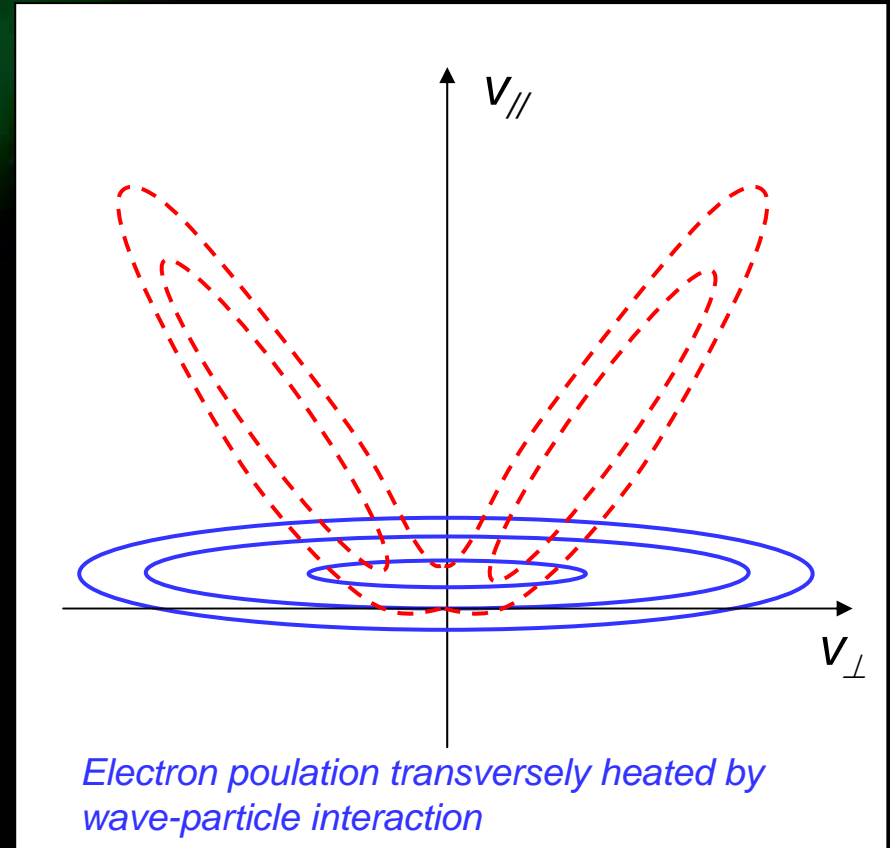
In a sense the opposite process to magnetic mirroring

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

Magnetic moment conserved



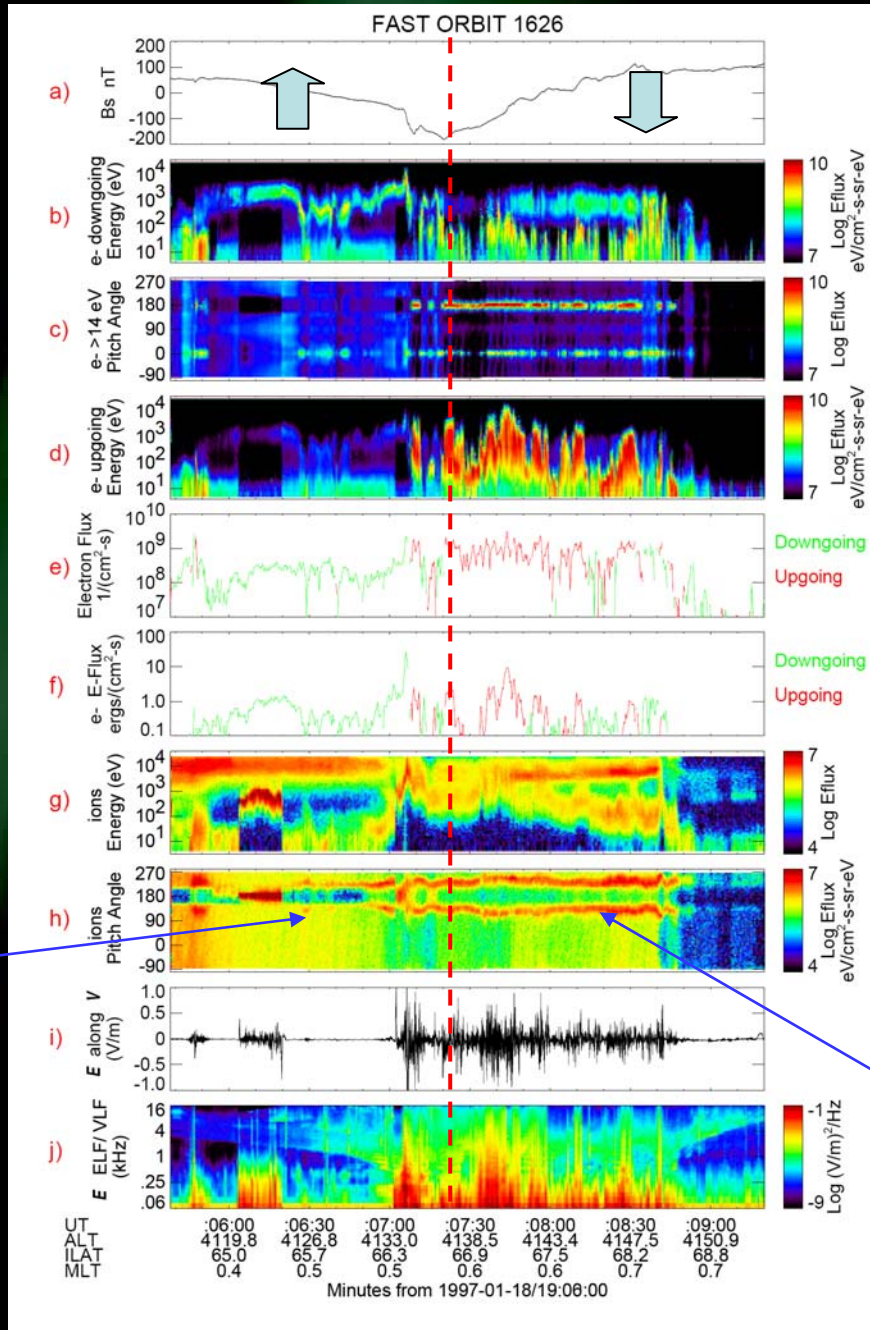
Weaker B means α decreases



An ion distribution originally heated in the direction perpendicular to B will fold up to a conic

$$\alpha = \sin^{-1} \sqrt{\frac{B}{B_0}}$$

Ion conics



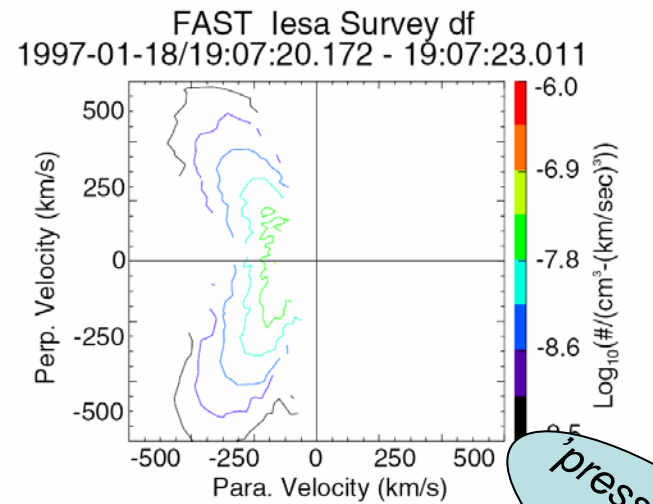
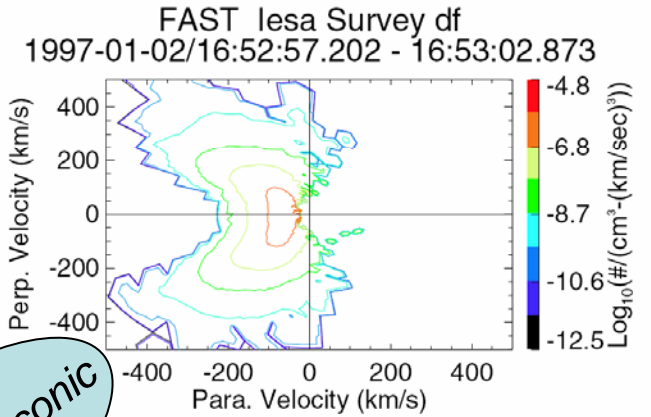
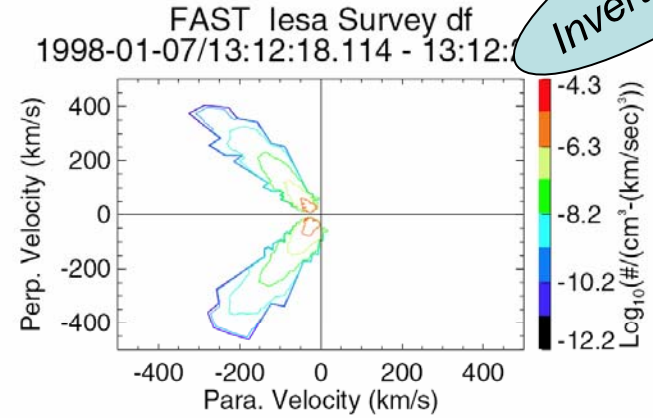
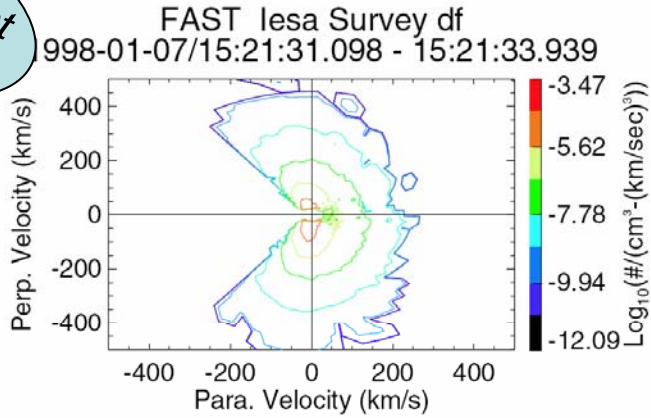
Ion conics – adiabatic motion

Downward current region conic

Inverted V conic

BBELF?

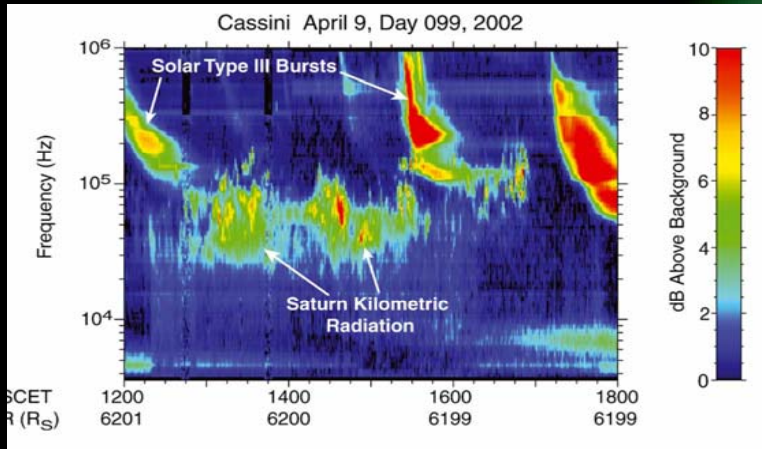
EMIC?



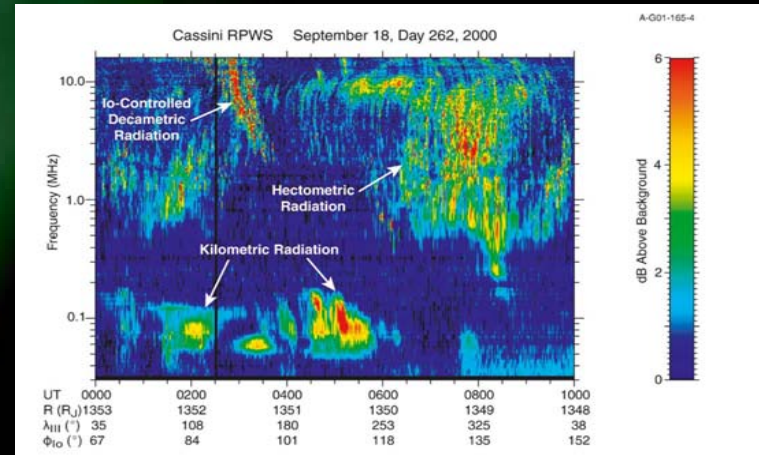
Alfvénic auroral conic

'pressure cooker' conic

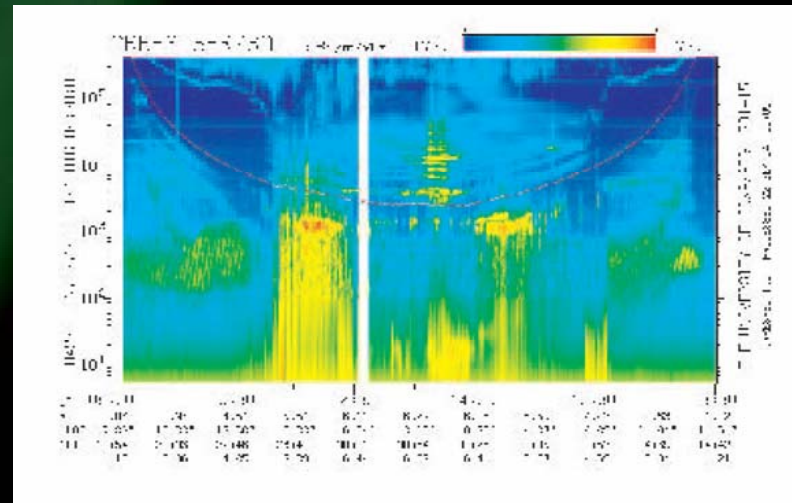
Waves in upward current region



Saturn kilometric radiation



Jupiter hectometric radiation

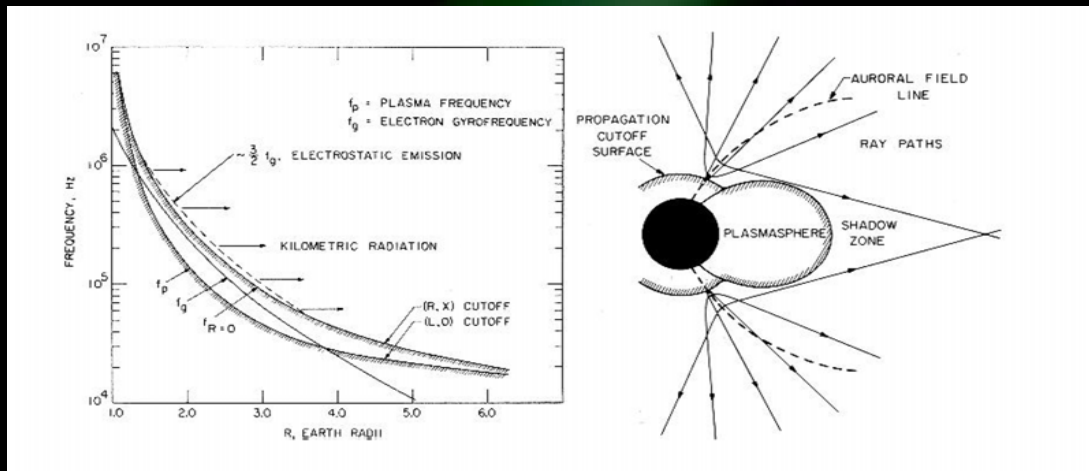


Auroral kilometric radiation

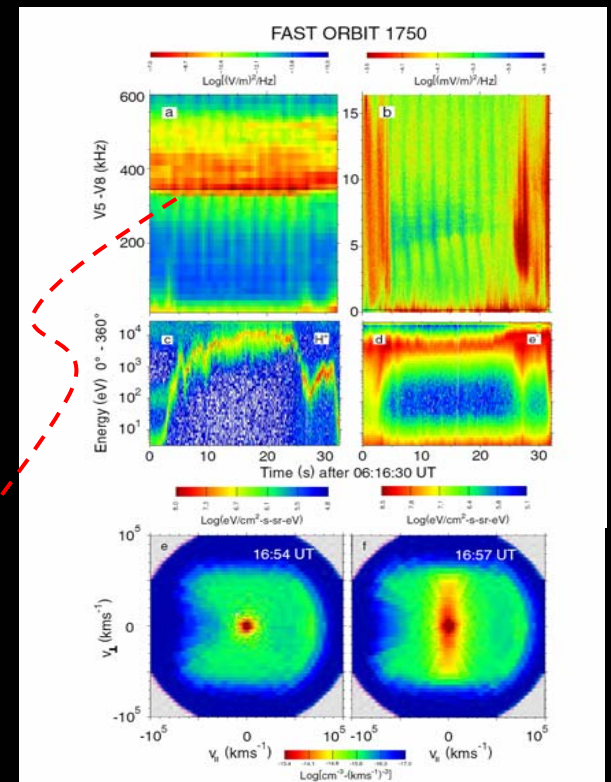
Auroral kilometric radiation

Dominating radiative feature of auroral zone

Generated by cyclotron-maser instability in auroral acceleration region



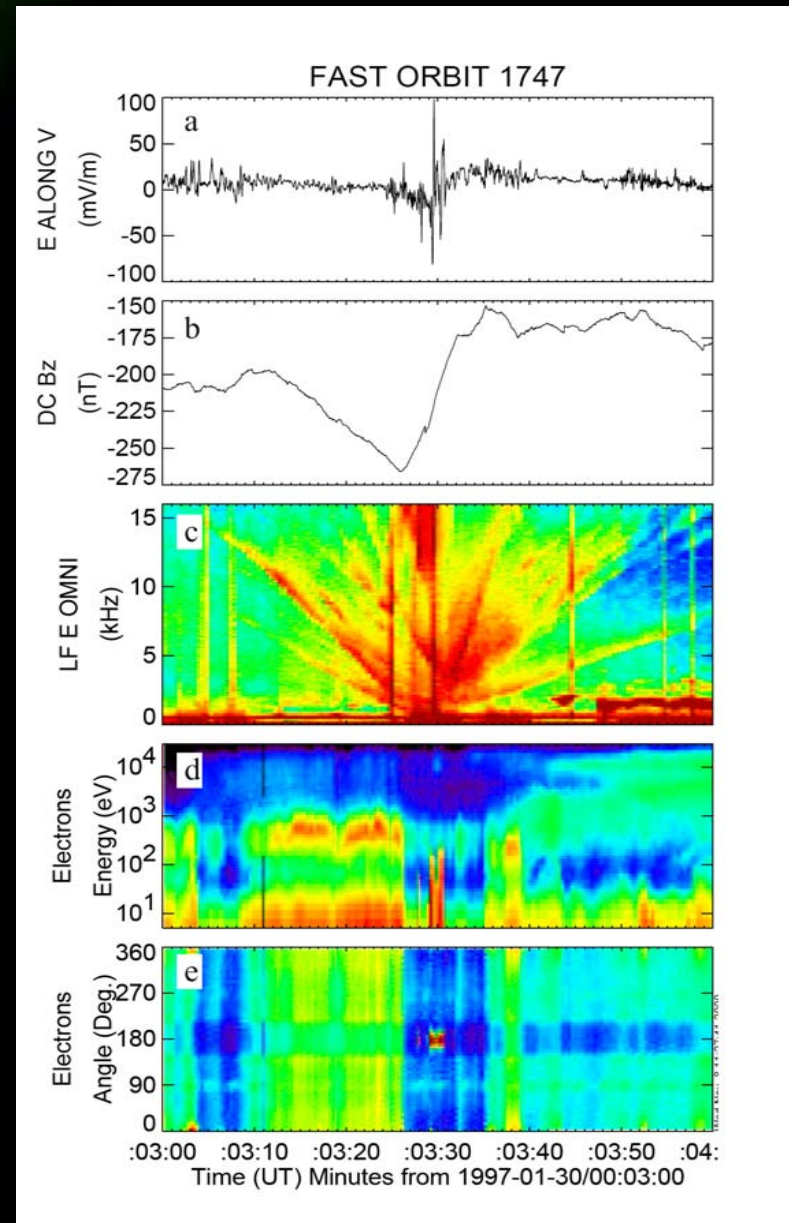
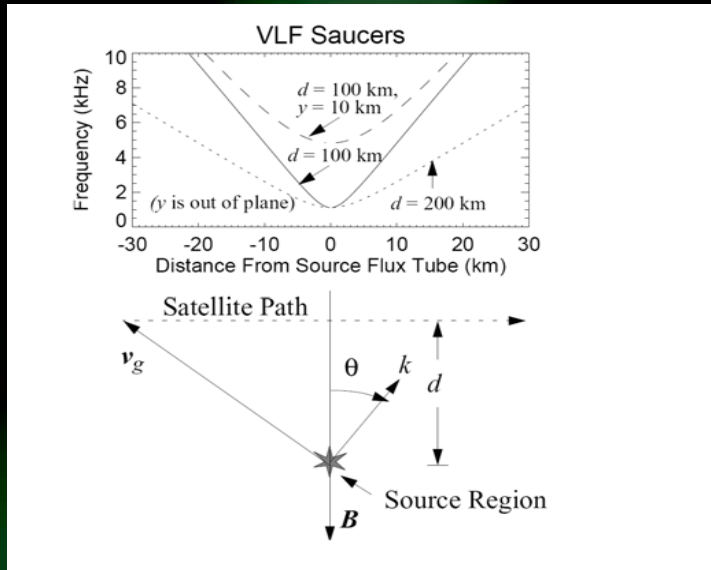
Lower cutoff at ω_{ce} of the source region.



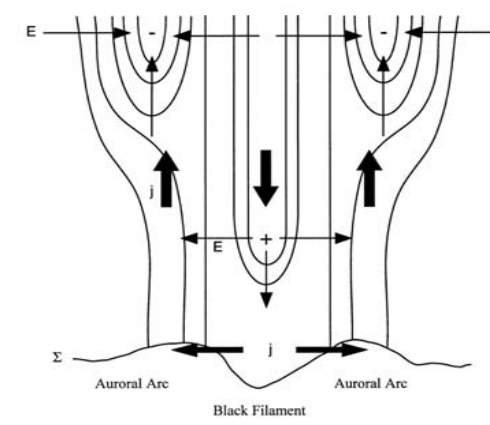
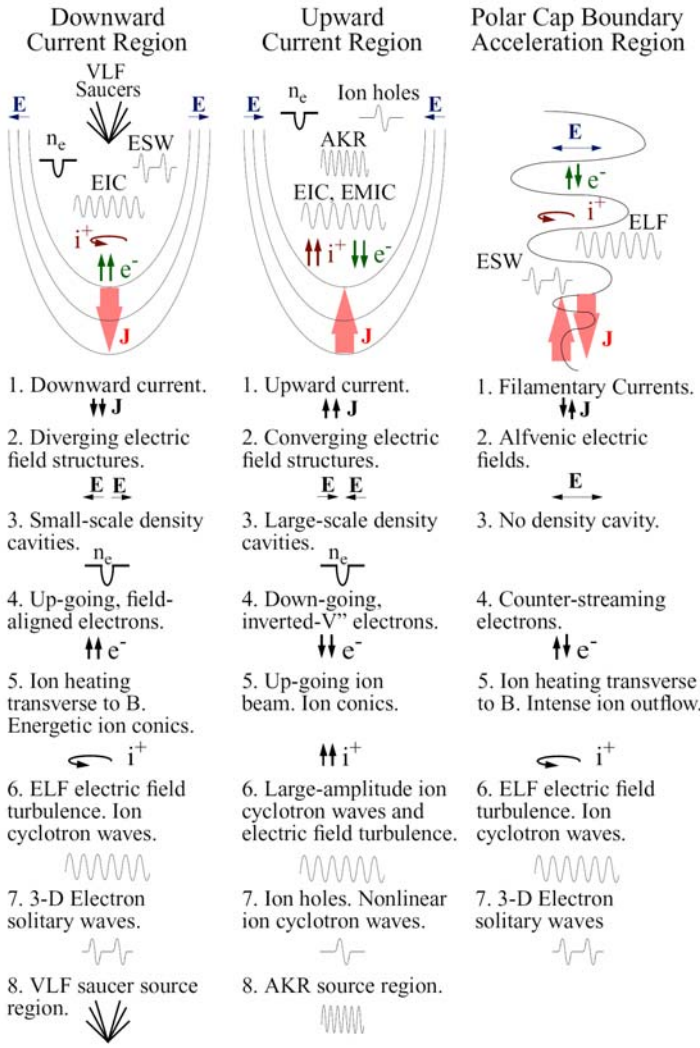
Waves in downward current region

VLF saucers

- Often most prominent wave feature of downward current region.
- k larger angle for higher frequencies
- Probably generated by upward ion beams



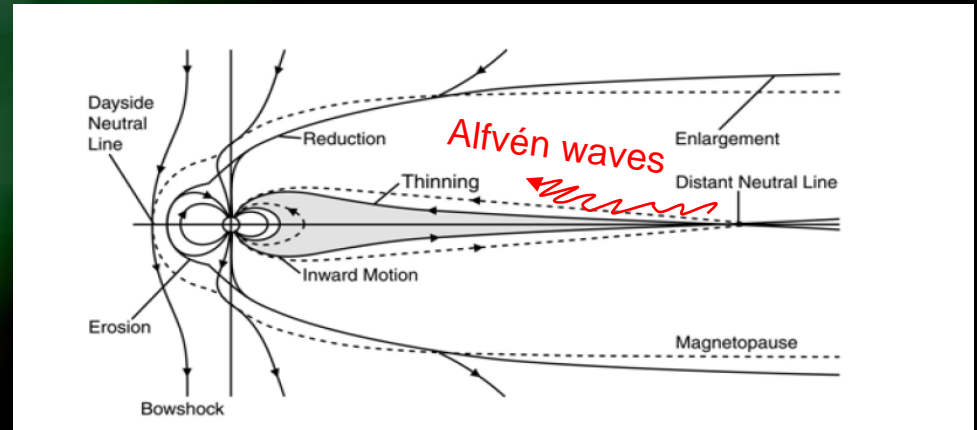
The symmetry between the upward and downward current regions



Dynamic MI-coupling

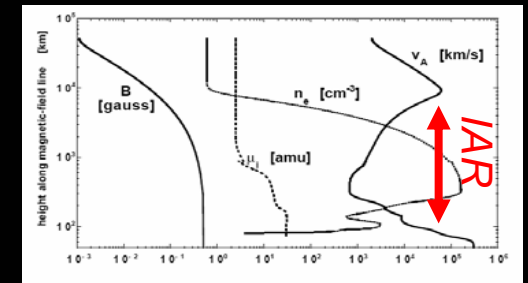
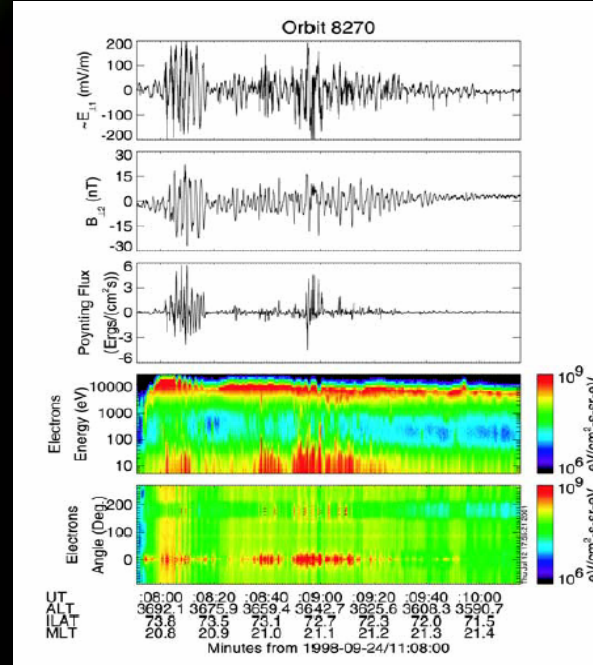
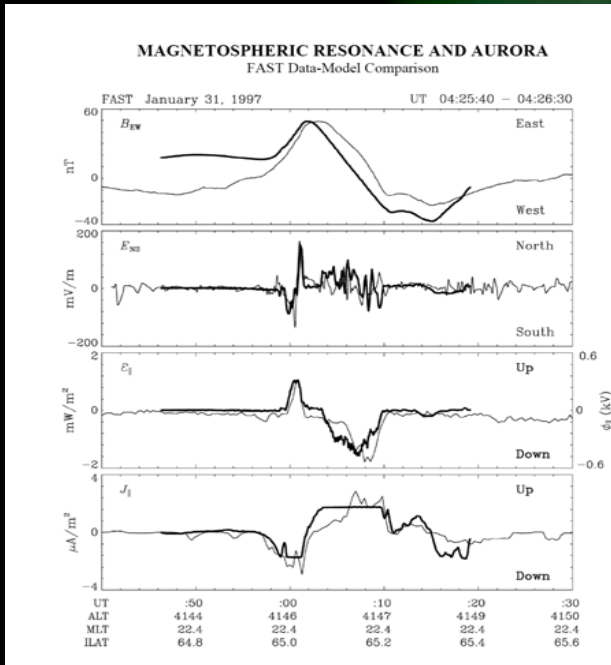
Alfvén wave driven aurora

X-line aurora

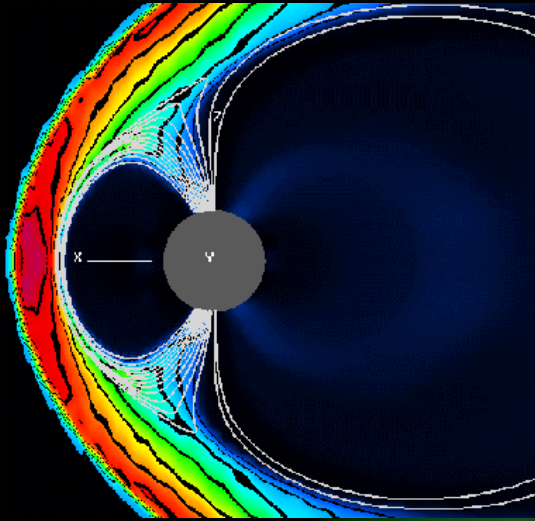


Field-line resonances

Ionospheric auroral resonator

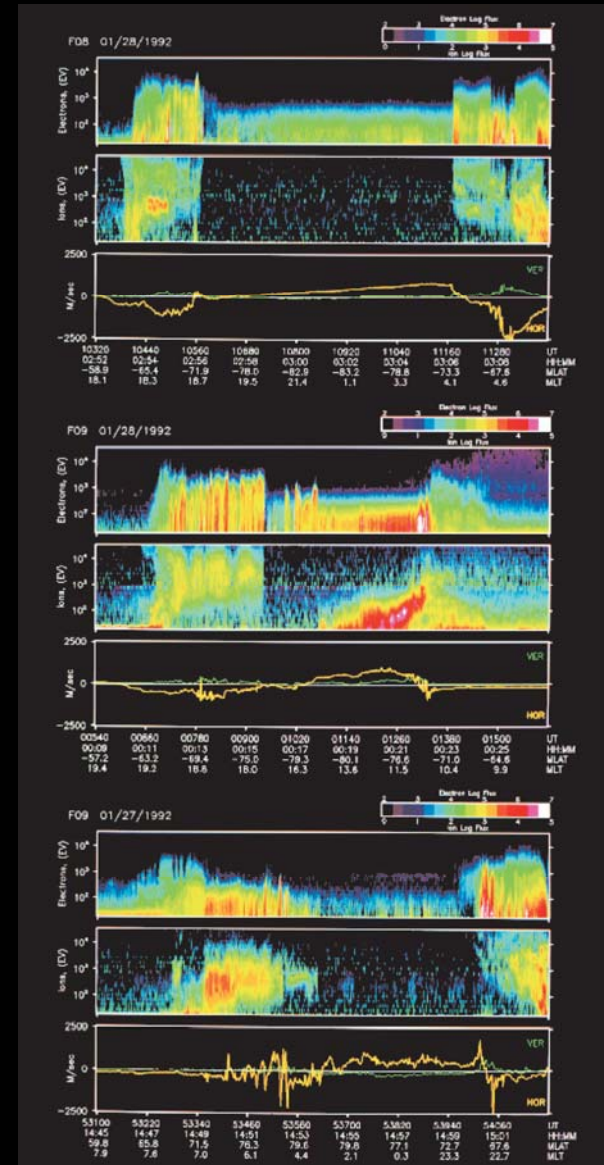
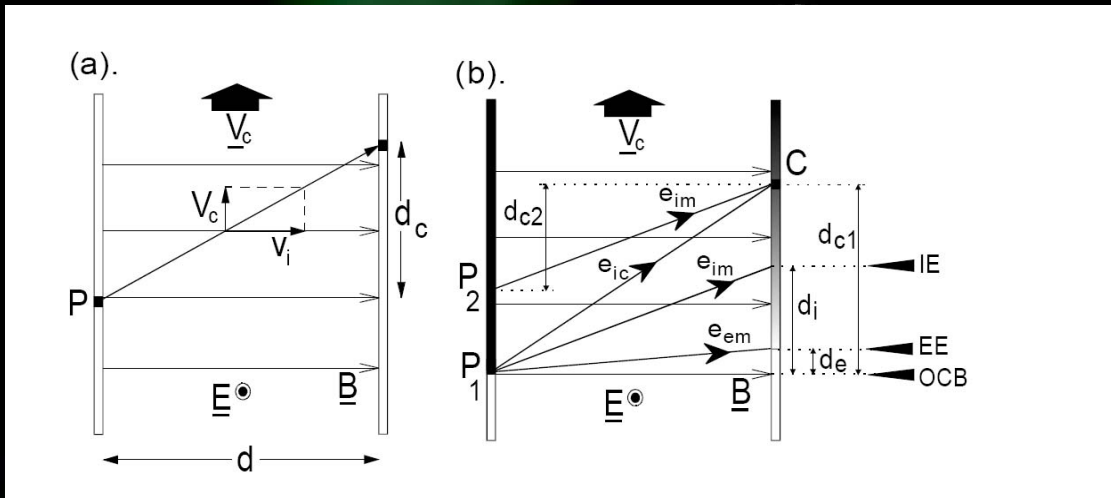


Cusp and dayside aurora



Direct entry of magnetosheath plasma
Fedder et al. (1997)

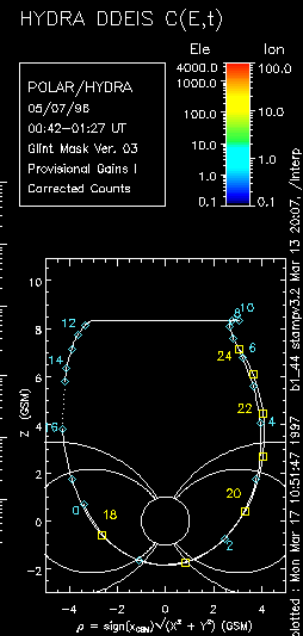
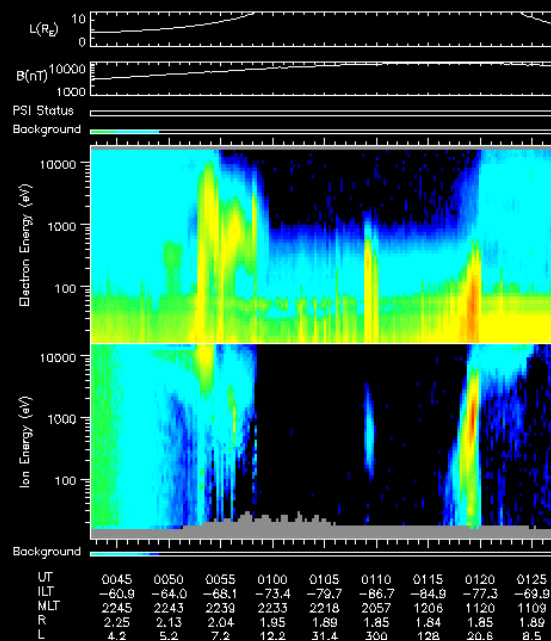
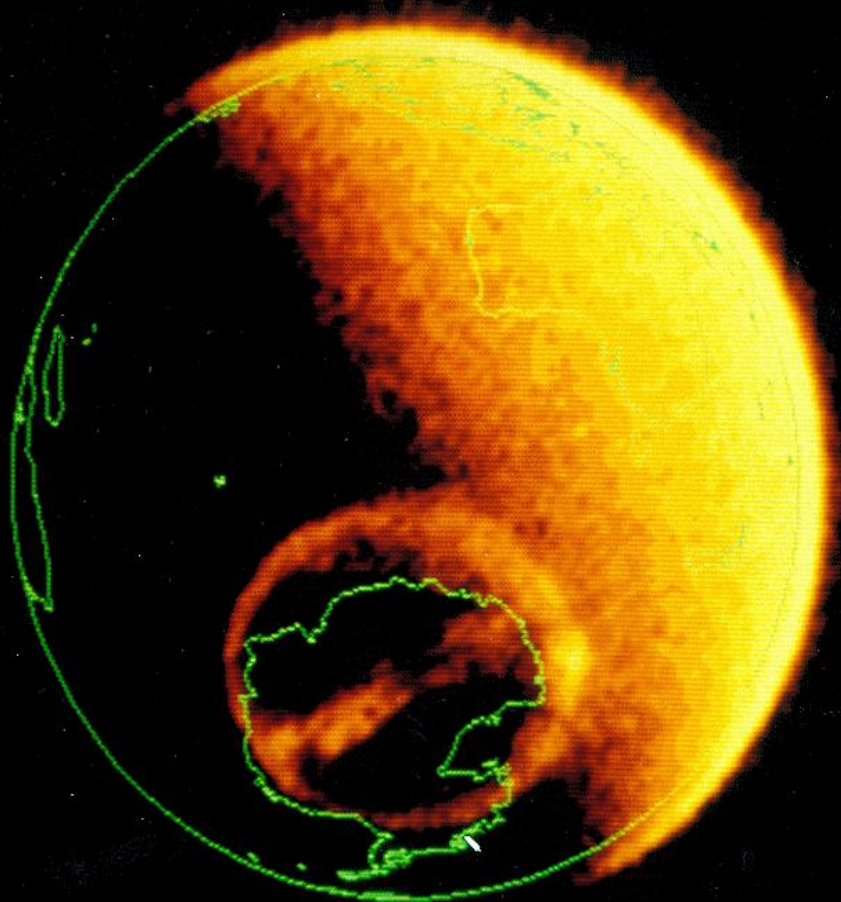
Velocity filter effect



DMSP 9 data (Lyons et al., 1999)

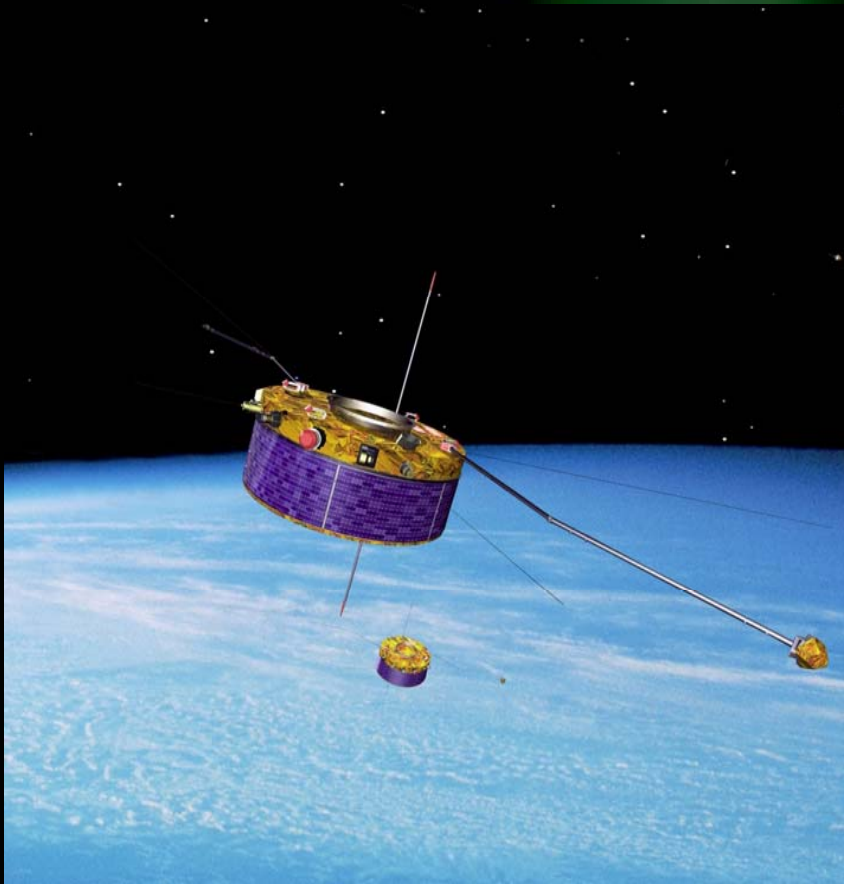
Theta aurora

What are they???



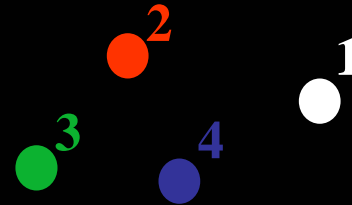
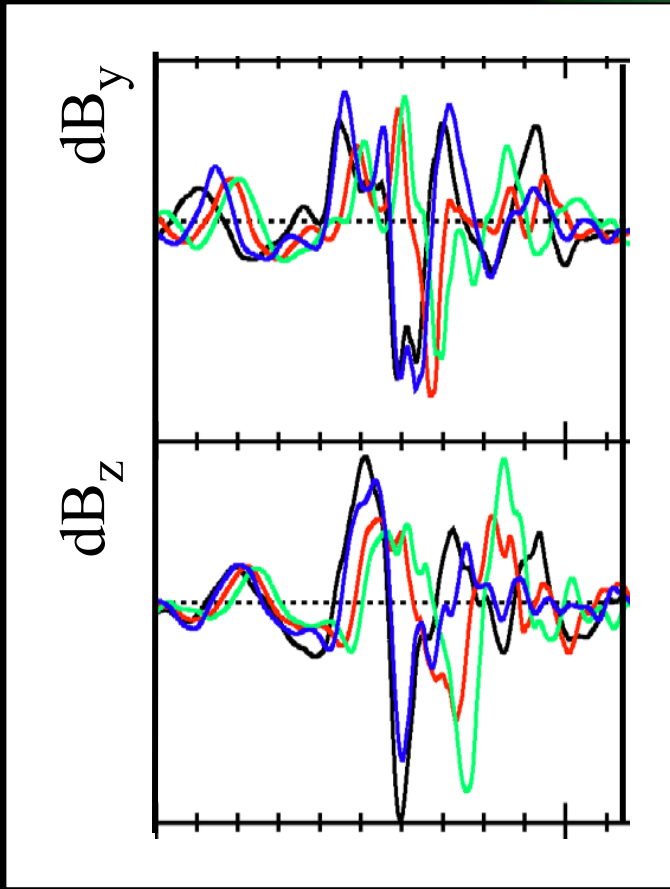
Cluster multi-point measurements

Seeing the temporal evolution

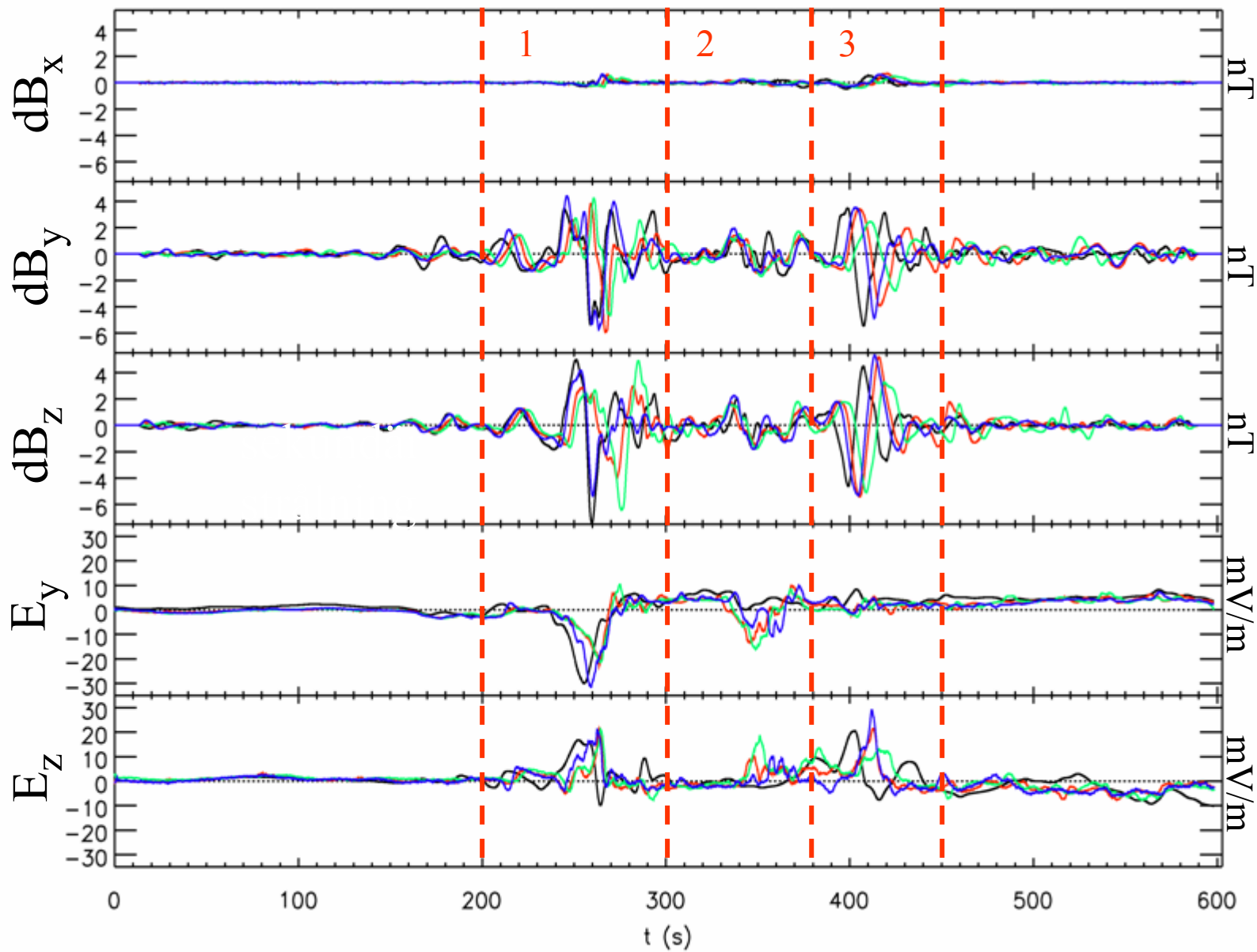


- Launched 2000
- Apogee: $20 R_E$
- Perigee: $4 R_E$
- Separations:
200-10000 km

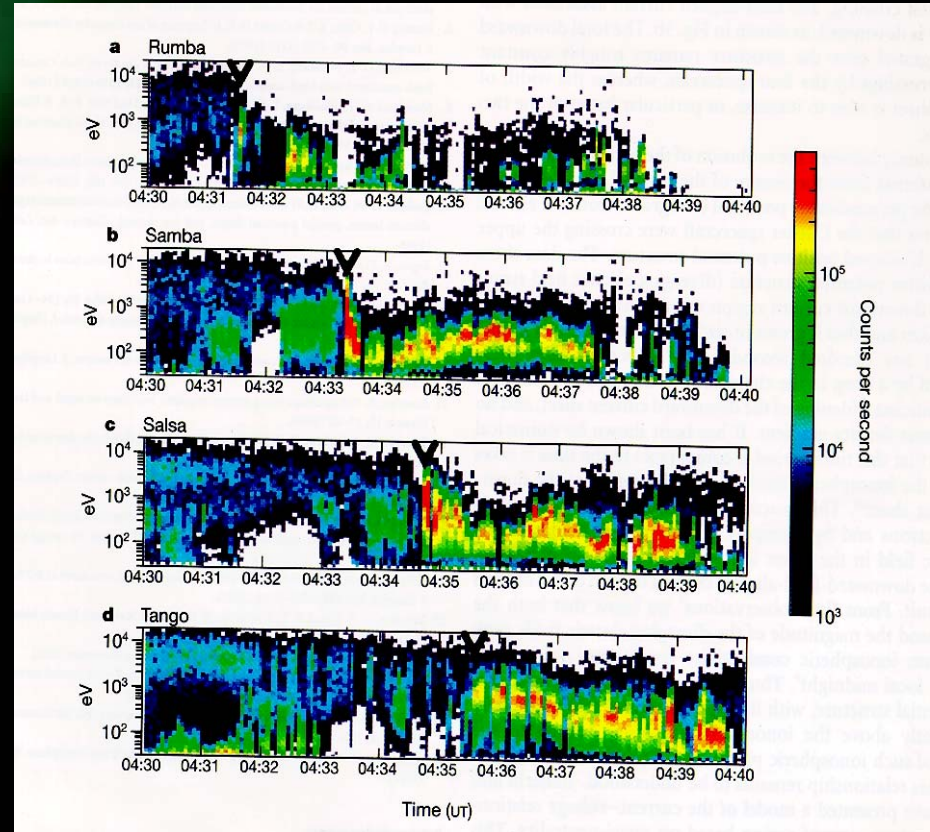
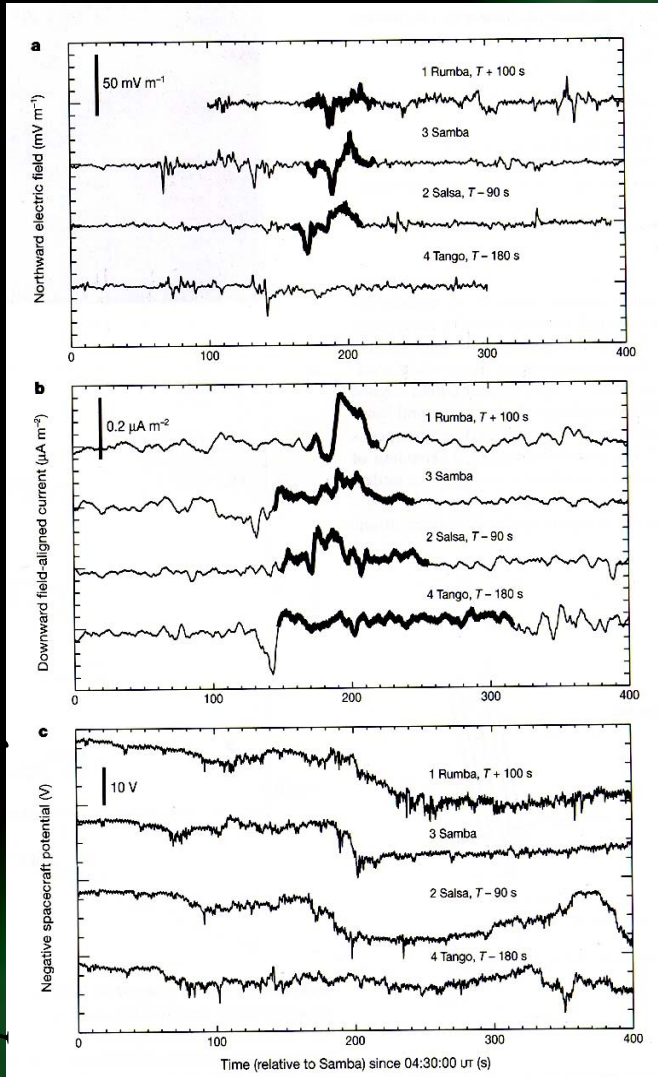
Interpreting Cluster multipoint measurements



Cluster field data



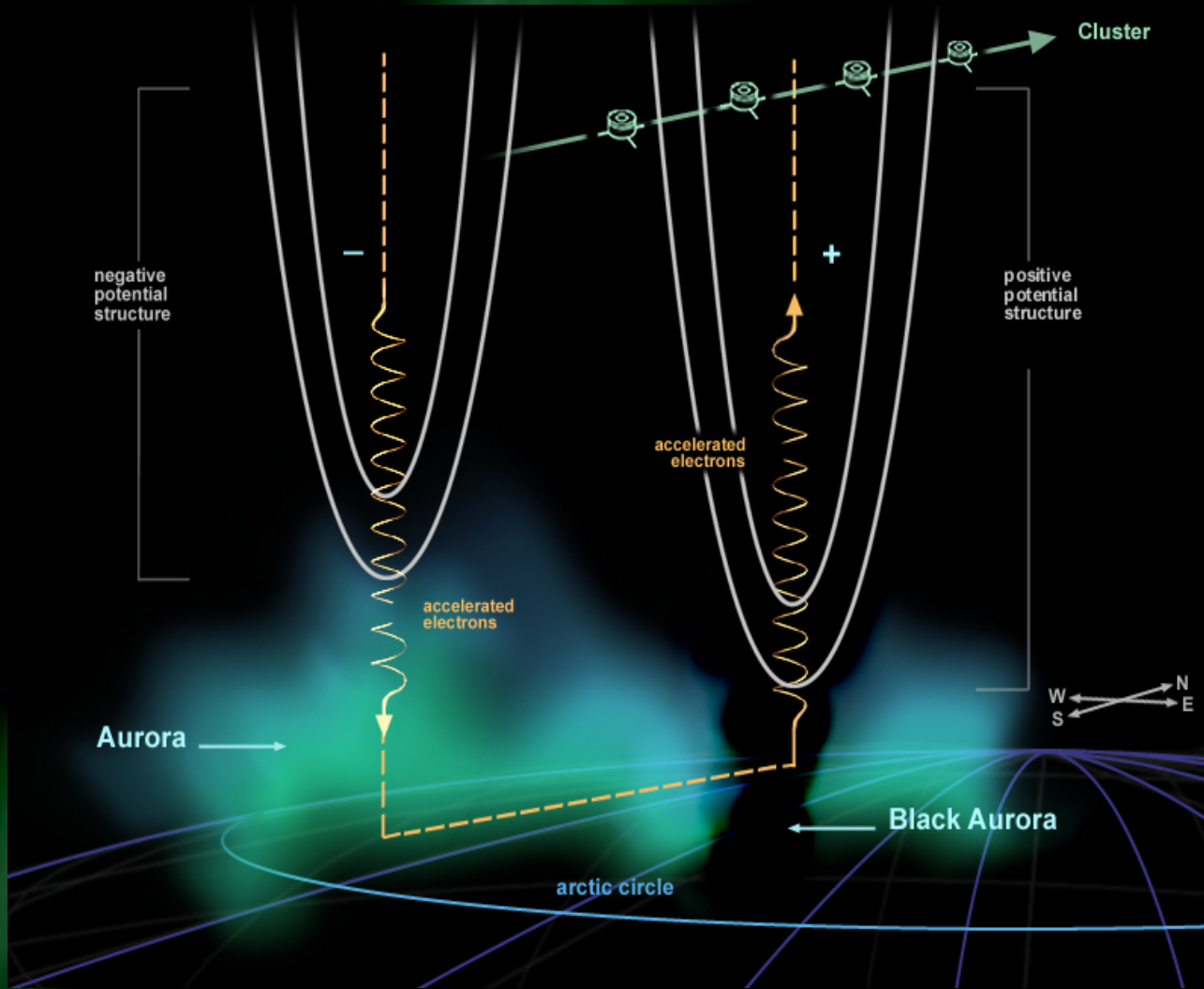
CLUSTER multipoint measurements (1)



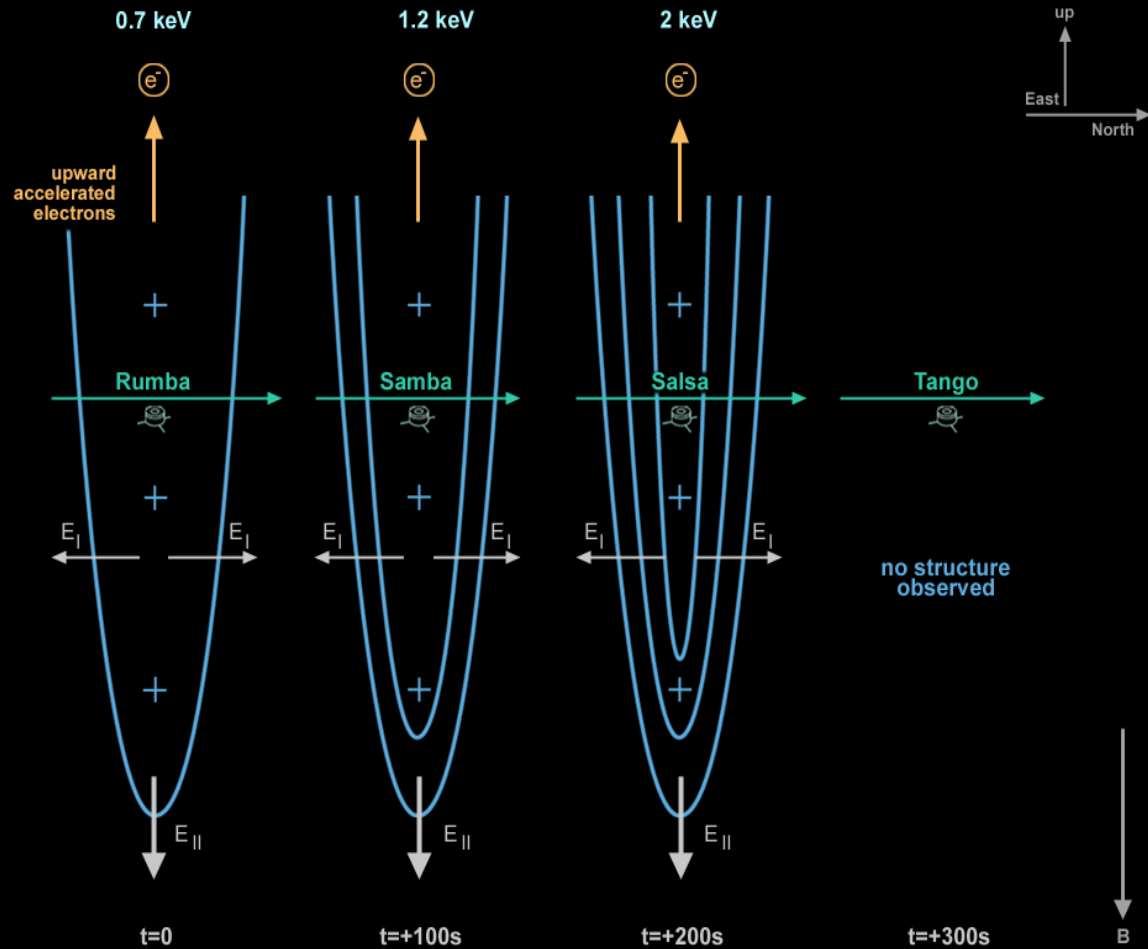
Marklund et al, 2001

plasma

Cluster passage through black aurora

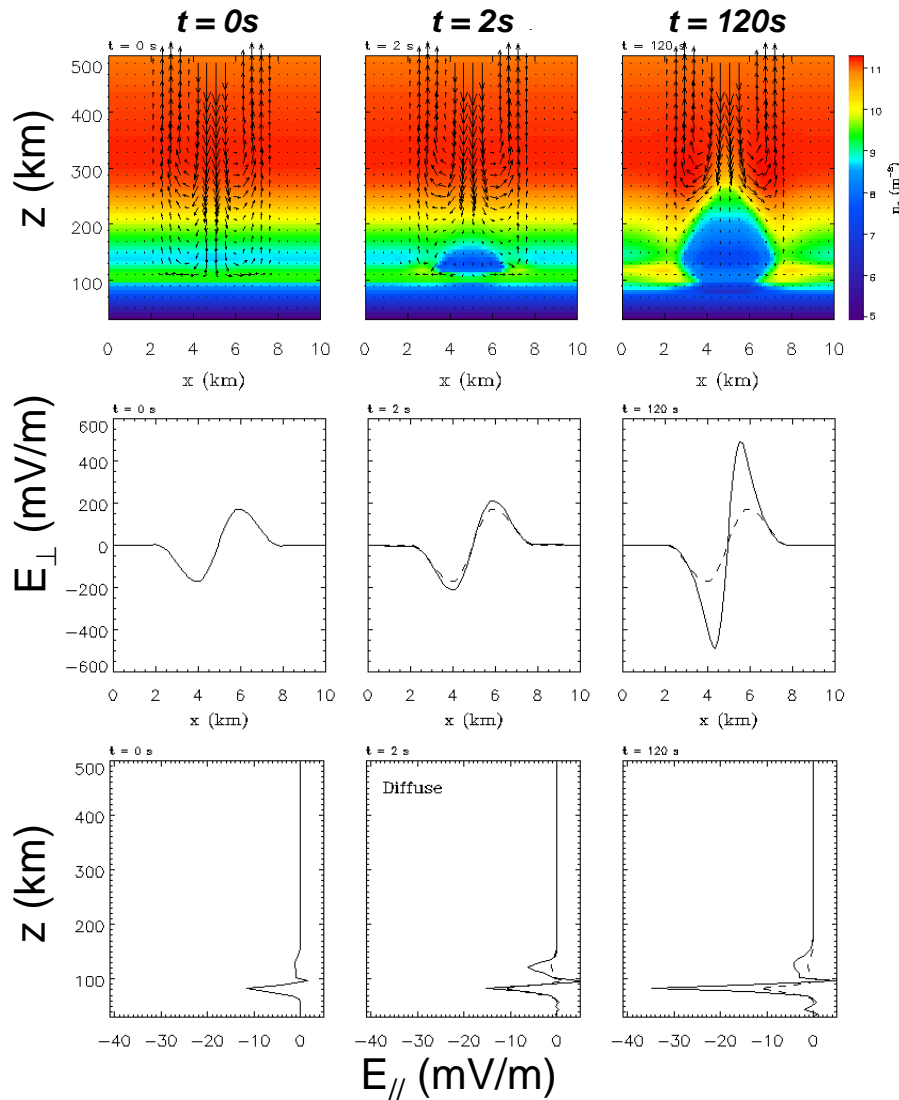


Temporal evolution of the acceleration potential above black aurora



The active role of the ionosphere

Density cavities



Simulations show deep density cavities formed by downward FAC and associated increased E -fields.

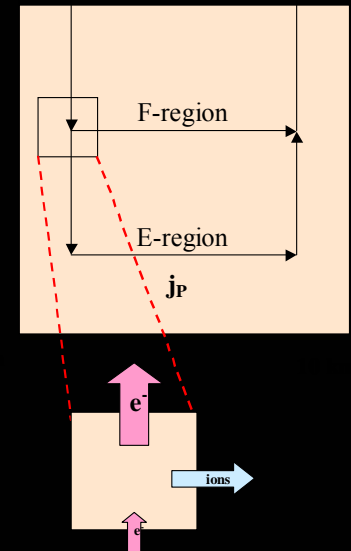
Important to take into consideration when mapping from high-altitude measurements.

Assumptions:

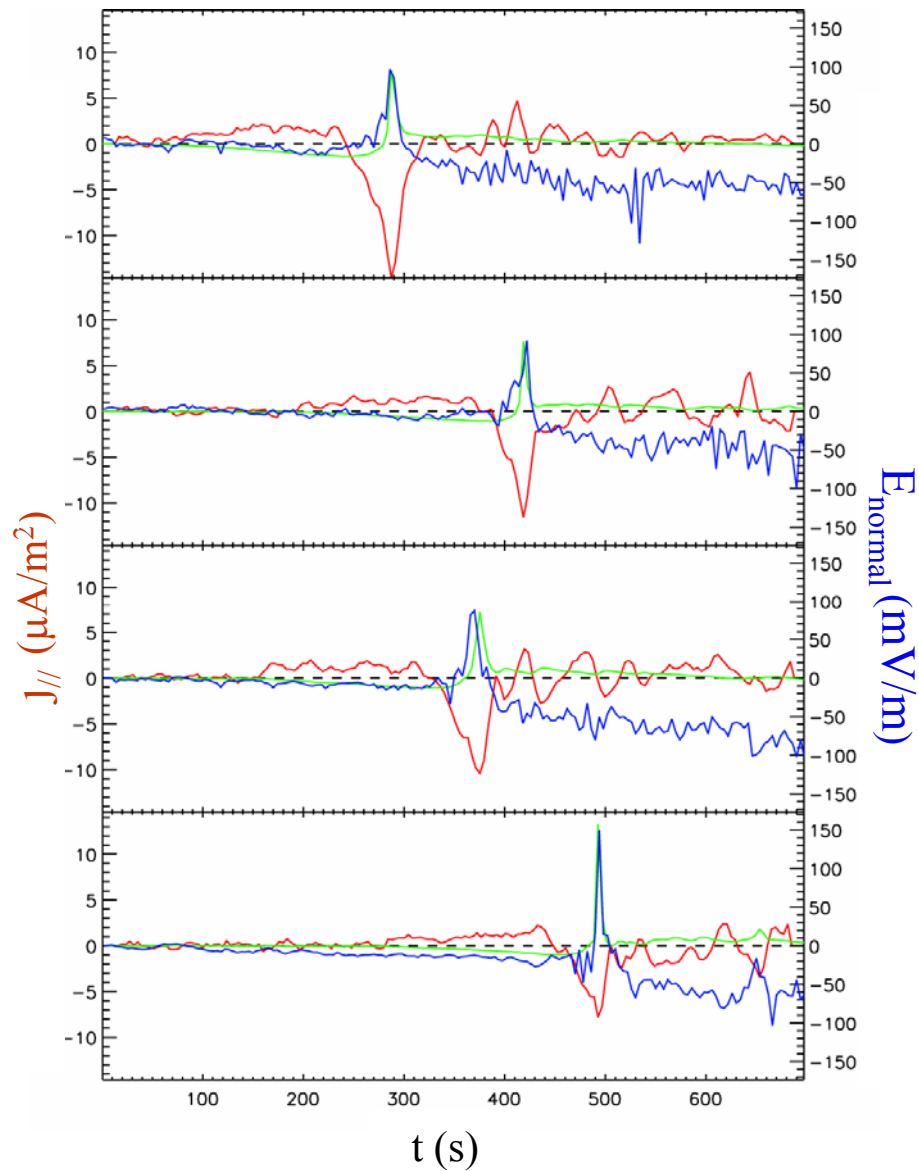
j_{\parallel} carried by e^{-}
 j_{\perp} carried by ions



$$\frac{\partial n_e}{\partial t} = -\frac{1}{q} \frac{\partial j_{\parallel}}{\partial z}$$



Cluster data, 2004-02-18, and model results



S/C 1-4

NH, $\text{MLT} \approx 4$, $\text{ILAT} \approx 66$

$E_n = 1$ mV/m

$E_t = -3$ mV/m

$\Sigma_{\text{P,BG}} = 5$ S

$k_{\text{down},1} = 0.31$ $\text{Sm}^2/\mu\text{A}$

$k_{\text{down},2} = 0.40$ $\text{Sm}^2/\mu\text{A}$

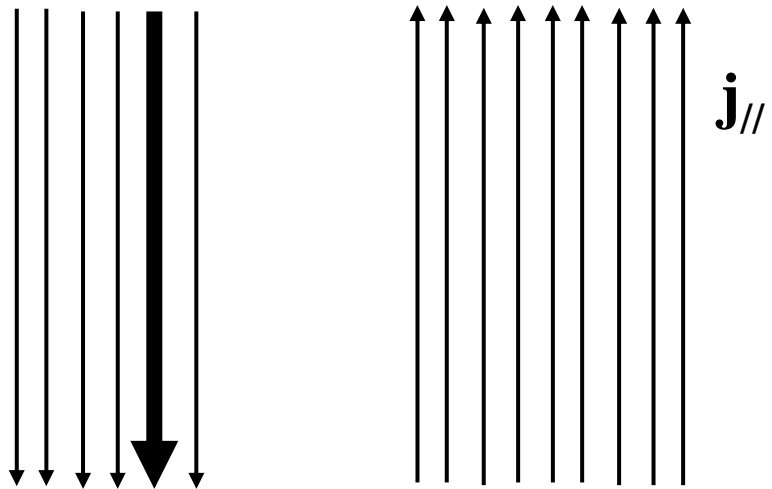
$k_{\text{down},3} = 0.42$ $\text{Sm}^2/\mu\text{A}$

$k_{\text{down},4} = 0.72$ $\text{Sm}^2/\mu\text{A}$

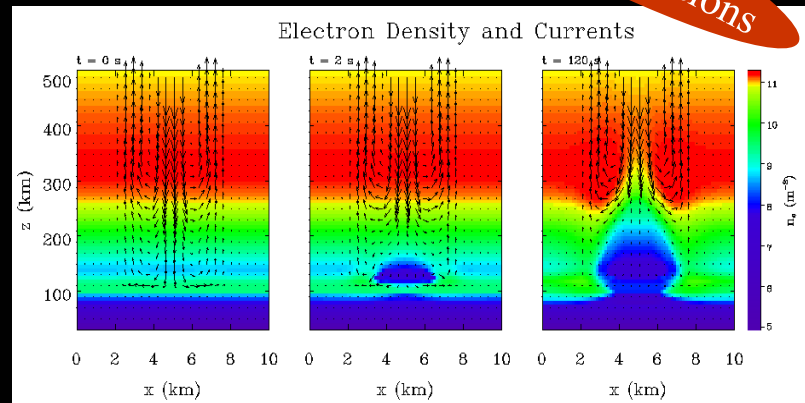
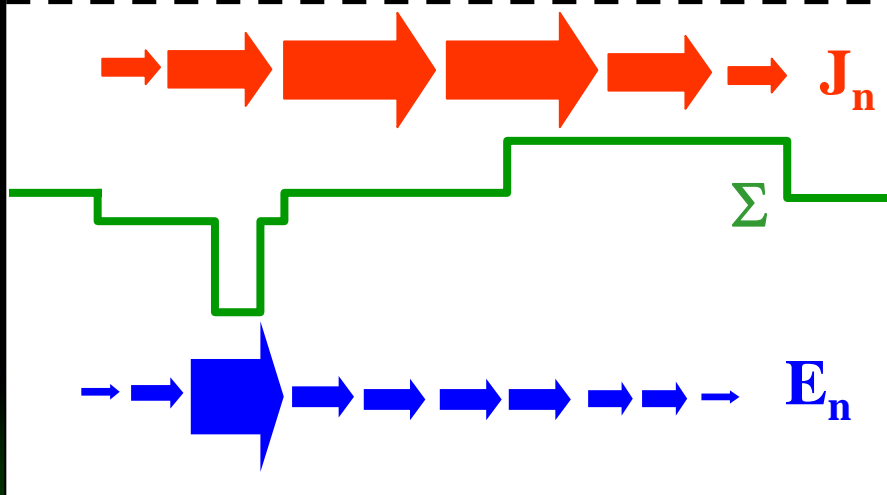
Model – ionospheric modification by downward FAC

From earlier simulations

Magnetosphere



Ionosphere



$$J_n = \int j_{||} dn + \Sigma_{P,0} E_{n,0} + \Sigma_{H,0} E_t$$

$$\Sigma_P = \Sigma_{P,0} + \begin{cases} k_{down} j_{||} & \text{downward } j_{||} \\ 0 & \text{upward } j_{||} \end{cases}$$

$$\Sigma_P \geq \Sigma_{P,min}, \quad \Sigma_H = 2\Sigma_P$$

$$E_n = \frac{\Sigma_{P,0} E_{n,0}}{\Sigma_P} + \frac{(\Sigma_H - \Sigma_{H,0}) E_t}{\Sigma_P} + \frac{1}{\Sigma_P} \int j_{||} dn$$

2004-02-18

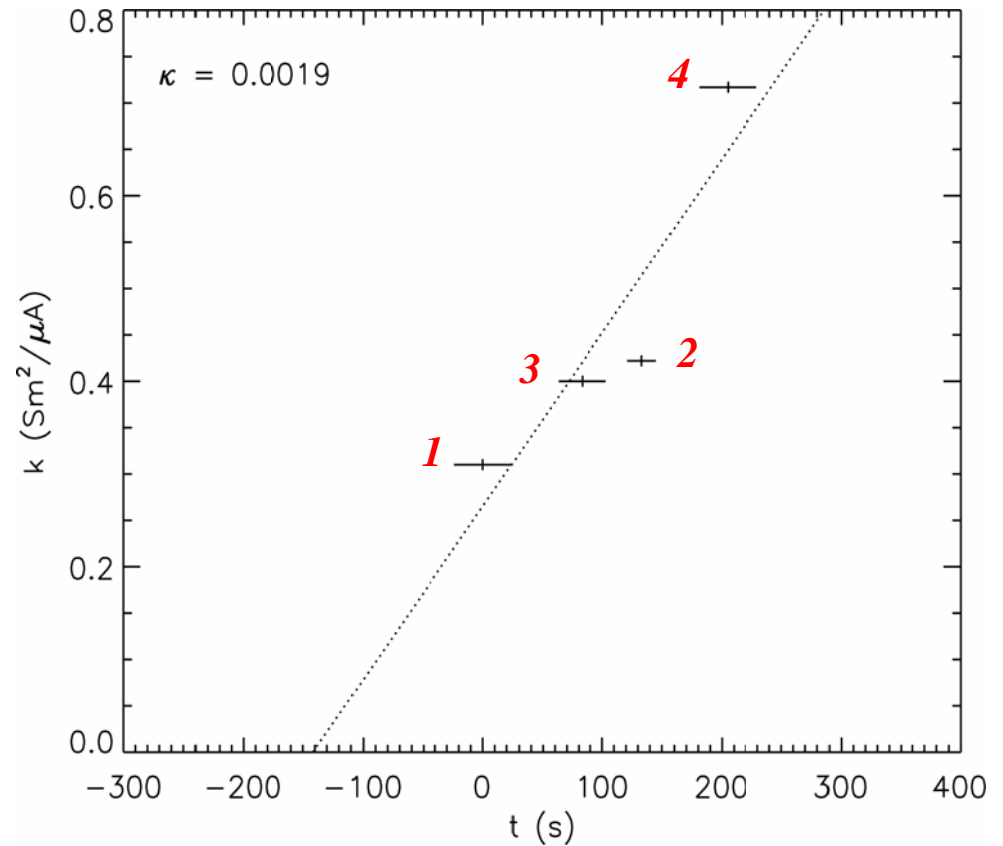
k as a function of time

Fit $k = \kappa t \Rightarrow$

$$\kappa = 1.9 \cdot 10^{-3} \text{ Sm}^2/\mu\text{As}$$

From simulations:

$$1 \cdot 10^{-5} \leq \kappa \leq 2 \cdot 10^{-3} \text{ Sm}^2/\mu\text{As}$$





Thank you for
your attention!